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Research Article

Keywords: Covid-19, epidemic dynamics, cascade of contacts, scale symmetry, growth and decline phases

Posted Date: November 12th, 2020

DOI: <https://doi.org/10.21203/rs.3.rs-106550/v1>

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COVID-19 scaling dynamics in growth and decline phases

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Abstract

The definition of optimal COVID-19 mitigation strategies remains worldwide on the top of public health agendas, particularly when facing a second wave. It requires a better understanding and a refined modelling of its dynamics. We emphasise the fact that epidemic models are phenomenologically based on the paradigm of a cascade of contacts that propagates infection. However, the introduction of ad-hoc characteristic times and corresponding rates spuriously break their scale symmetry.

Here we theoretically argue and empirically demonstrate that COVID-19 dynamics, during both growth and decline phases, is a cascade with a rather universal scale symmetry whose power-law statistics drastically differ from those of an exponential process. This involves slower but longer phases which are furthermore linked by a fairly simple symmetry. These results explain biases of epidemic models and help to improve them. Due to their generality, these results pave the way to a renewed approach to epidemics, and more generally to growth phenomena.

Introduction

The common root of most epidemics models is a cascade paradigm that can be traced back to their emergence with Bernoulli and d'Alembert¹, who preceded the celebrated quatrain of Richardson on the atmospheric dynamics cascade². Indeed, their basic assumption is that each individual of the infected fraction I of a given population N will “on average” contaminate R_0 - the mythical “basic reproduction number”- individuals of the “susceptible” fraction S of this population. The complement to the fraction S is the population fraction $R = N \setminus S$ of the “removed” individuals, either having become immune or deceased, hence the acronym SIR for this type of models. Originally³, these models were purely compartmental models, therefore with population fractions assumed to be homogeneous and the only variables are therefore their cumulative incidences $s(t)$, $i(t)$ and $r(t)$, usually driven by ordinary differential equations. Space inhomogeneity and resulting fluxes between different locations⁴, as well as age dependencies⁵ are now explicit in the current distributed epidemics models whose space time fields are driven by partial (integro-) differential equations.

At odd with cascade models developed in turbulence^{6,7,8} and other domains⁹, current epidemic models still rely on characteristic times that are assumed to be independent of the scale of observation, although there is no physical ground for this hypothesis. For instance, most epidemic models depend on or strive to estimate scale independent characteristic times T_c between contacts and T_r until recovery, whose ratio T_c/T_r , under strong homogeneity assumptions (e.g., uniform or Poisson distributions), yields a straightforward definition of the basic reproduction number R_0 . These characteristic times have been implemented in various epidemic models¹⁰ to obtain more elaborated estimates of an initial R_0 ranging from 2.06 – 2.52¹¹ to 3.6 – 4.0¹² and an “effective reproduction rate” $R \approx 4.08 \pm 0.36$ ¹³ for the COVID-19 pandemic. None of these studies tried to analyse the possible scale dependency of their estimates, whereas they admitted a sample dependency.

On the contrary, it is easy to point out that the time between contacts T_c sensitively depends on scale. Indeed, anyone would agree that the effective time between contacts t_c strongly depends on the mobility of the concerned agents, which has fluctuations over a wide range of space-time scales, e.g., from pedestrian walks to intercontinental flights. Not only recent epidemic models take a particular care to refine its representation, but the main goal and justification of the confinement strategies is to reduce the mobility nearly to zero^{15,16,17,18}. Indeed, the intensification of mobility is suspected to have

52 led to qualitatively change the type of spread: from normal diffusion (e.g. for the Middle Age
53 plague¹⁹) to anomalous diffusion. Normal diffusion does provide characteristic times due to the fact
54 that the generating Gaussian distribution has a fall-off faster than an exponential (e.g., the probability
55 to exceed 3 root mean squares (RMS) is only of the order of 10^{-5} , i.e., a Gaussian variable is never
56 too far from the mean). For instance, the spread radius $\rho(t)$ can be measured by the RMS $\sigma(t)$ of
57 particles emerging at the same time from the same source, which grows only like $t^{1/2}$, i.e. much
58 slower than a ballistic diffusion ($\rho(t) \approx t$). On the contrary, anomalous diffusion is generated by
59 Levy stable distributions that are devoid of characteristic times due to their power-law fall-off, i.e.,
60 the probability to exceed a distance $\rho(t)$ on a given time falls off as $\rho(t)^{-(1+\alpha)}$, with $0 < \alpha < 2$ ²⁰,
61 ²¹. As a consequence, the RMS $\sigma(t)$ does not exist. Furthermore, the probability to exceed c times a
62 given distance $\rho(t)$ decreases only by the factor $c^{-\alpha}$, which shows that we can easily be far away
63 from the mean... which furthermore exists only for $\alpha > 1$! Such anomalous diffusion due to human
64 mobility has received some empirical support from a large scale experiment ironically nicknamed
65 “where is George”. This experiment corresponds to first tracking via an internet site, over ten millions
66 of displacements of about half a million of one dollar bills (bearing the image of George Washington)
67 then to analyse them²². The important result obtained was that this diffusion was anomalous and
68 nearly isotropic in space and time with a common exponent $\alpha \approx 0.6$. It therefore corresponds to a
69 strong super diffusivity. It was also shown that a random walker in a Chernobyl contaminated field
70 accumulates radiation doses that also anomalously fluctuate, contrasting sharply with the usual
71 average estimates²³. Refined small scales statistics of time between contacts could be obtained with
72 the deployment of mobile tracking applications, which nevertheless should protect privacy²⁴. Another
73 important aspect to be taken into account is the fractality of measuring networks²⁵.
74

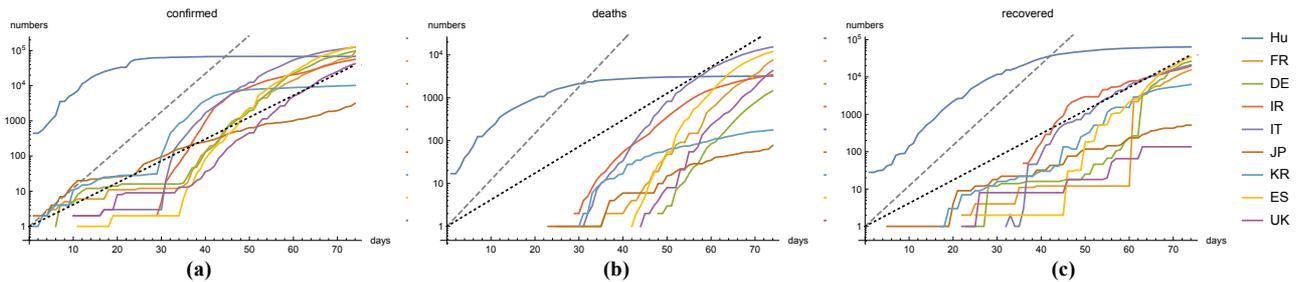


Fig.1: Sub-exponential growth phase: log-linear plots of (a) confirmed, (b) death, (c) recovered cumulative incidences $X(t)$ of the following entities: Hubei, France, Germany, Iran, Italy, Japan, Korea, Spain and United Kingdom, over the period 22 January - 04 April. An exponential behaviour would correspond to straight lines like those drawn for characteristic times $T = 4; 7$ days, and therefore to doubling times $T_2 = T \ln(2) \approx 3; 5$ days (respectively dashed grey and dotted black straight lines). Due to their (strong) concavity, all the trajectories are sub-exponentials, i.e., grow slower than any of their local approximations by an exponential that corresponds to a tangent to these log-linear graphs.

75
76 Characteristic times T and the corresponding “doubling times” T_2 are related to differential equations
77 of the type:

$$78 \quad dX/dt \approx X(t)/T; T_2 = T \ln(2) \quad (1)$$

79 and their exponential solution $X(t) \approx X(t_0)\exp((t - t_0)/T)$, where the symbol \approx denotes an
80 equivalence that can be more general than a strict deterministic equality, e.g., asymptotic equivalence
81 on a given range of times or/and up to a slowly varying factor, or/and an equality in probability
82 distribution, etc.. The SIR model provides again a rather simple example of this behaviour. For a
83 constant population ($s(t) + i(t) + r(t) = n(t) = n = const$) it yields the following evolution
84 equation for the infected population:

$$85 \quad di(t)/dt \approx (R_0s(t)/n - 1)i(t)/T_c \approx (R_0s(0)/n - 1)i(t)/T_c \quad (2)$$

86 The second r.h.s. approximation is valid during the initial outbreak ($s(t) \approx s(0)$), whereas the
 87 awaited “flattening of the epidemic wave” ($\frac{di(t)}{dt} = 0$) occurs in the SIR model only when the
 88 susceptible population $s(t)$ has been significantly reduced with respect to the total population n ,
 89 hence it occurs at the time t_f such that: $s(t_f)/n = R_0^{-1}$ or $(i(t_f) + r(t_f))/n = 1 - R_0^{-1}$.

90
 91 **Database**

92 We will not discuss either the interests (e.g., analytically solvable³) or limitations (e.g., non-accounted
 93 spatial and intra-compartmental heterogeneities) of the SIR model that we only used to introduce
 94 basic concepts needed to explore the database of the Center for Systems Science and Engineering at
 95 Johns Hopkins University²⁶, more precisely its convenient github web site
 96 https://github.com/CSSEGISandData/COVID-19/tree/master/csse_covid_19_data. This database
 97 provides daily estimates of the “confirmed”, “deaths” and “recovered” cumulative incidences for 187
 98 countries. These data are refined for some regions, e.g.: the data on China are distributed according
 99 to all the provinces, metropolitan France is distinct from the overseas departments and the American
 100 data are provided according to 3261 urban agglomerations and dependencies.

101 For presentation clarity, results are only displayed for the following entities: Hubei, France, Germany,
 102 Iran, Italy, Japan, Korea, Spain and United Kingdom. The choice of these time series is based on
 103 their length, as well as their relative high counts. Both factors could presumably reduce artefacts such
 104 as incomplete counts or even possible offsets for various reasons, including political ones. However,
 105 important changes in tests or/and count methodologies or simply data transmission delays did
 106 introduce non negligible, artificial fluctuations in these series. For instance, important changes on
 107 COVID-19 infection evaluations such as +15100 in the Hubei province occurred on 13 February
 108 and +26800 in France on 12 April. An additional cumulation of 325 cases in the Hubei province
 109 was added to the datasets on ad-hoc dates, a month later the official ending of the epidemic spread in
 110 this area. Often, the deaths in rest homes began to be included only at a later stage of the epidemic.
 111 Furthermore, data were not always daily provided, therefore introducing a zero increment between
 112 two consecutive days, followed by an artificially enlarged increment for the second day. However,
 113 our data analysis is rather robust with respect to this sort of time translation, contrary to the direct
 114 analysis of the time evolution of the cumulative incidence $X(t)$. There are other disrupting examples
 115 of data adjustments, such as the cumulation of the “deaths” incidences was decreased by 1918 cases
 116 in Spain on 25 May.

117

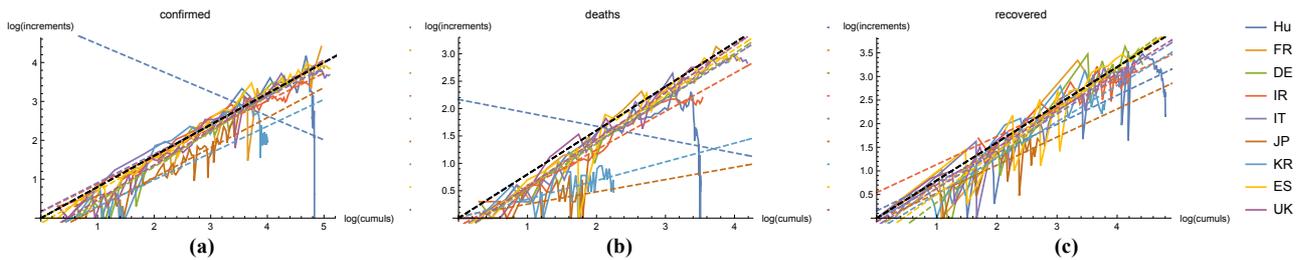


Fig.2: Cumulative-incremental analysis of the growth phase: log-log plots of the couples $(X(t), \Delta X(t))$, where $\Delta X(t) = X(t) - X(t - \Delta t)$ is the increment (the time increment Δt being one day) of (a) “confirmed”, (b) “deaths”, (c) “recovered” cumulative incidence $X(t)$ of the following entities: Hubei, France, Germany, Iran, Italy, Japan, Korea, Spain and United Kingdom, over the period 22 January- 04 April. Graphs should be read from left to right to follow the time arrow ($X(t)$ being non decreasing with time) and the almost vertical parts on the right-hand side corresponds to compressed views of decline phases to be expanded in Fig.3.

These graphs display a rather universal behaviour, especially a common scaling/power-law behaviour $\Delta X(t_i) \approx \alpha X(t_i)^b$, a discrete version of the differential equation Eq. 4, over a non negligible range of couples $(X(t), \Delta X(t))$ and therefore of time, with only a few exceptions. The best linear fits are represented by the dashed lines (of the same colour as the graph of this entity), with high determination coefficients R^2 and limited dispersion of the estimates of the feedback exponent b (see Supplementary Information). The black dashed line indicates the average estimates over the period of 22 January - 04 April.

119 However other entities, and other epidemic data as well, display similar graphs that have called
120 attention to power-law growth, including with the help of log-log curve fittings^{27, 28, 29}. But, these fits
121 are sensitive to the choice of the initial time and data artefacts. We therefore developed a specific data
122 analysis presented below.

123 We analysed data from 22 January till 07 June hence containing both the end of the crisis in Hubei
124 and its beginning in Europe to enable to analyse both stages of the epidemic. Two intermediate dates
125 were used for a comparative analysis: 4 April as a tentative date for the well-advanced growth phase
126 in Europe (about 3 weeks after the pandemic announcement by WHO), and 16 April as the ending
127 date of the epidemic for the Hubei province). Nevertheless, tentative explorations until the beginning
128 of November have confirmed the interest of a follow-up paper focused on the beginning of the second
129 wave.

130

131 **Growth phase analysis**

132 Figure 1 shows that the time evolution of the cumulative incidence $X(t)$ of the aforementioned
133 categories and entities are sub-exponential. Hence, we introduce a scaling dependency (i.e. a power-
134 law) of their characteristic times $T(t)$ (and doubling time $T_2(t)$) with respect to the current scale
135 $X(t)$:

$$136 \quad T(t) \approx X(t)^{1-b}/a; \quad T_2(t) \approx \ln(2)X(t)^{1-b}/a \quad (3)$$

137 where the exponent b can be close to unity, but nevertheless distinct of it, and the prefactor a is no
138 longer the inverse of a (characteristic) time except for $b = 1$. The presence of the prefactor $X(t)^{1-b}$
139 in the local time $T(t)$ (for $b \neq 1$ in Eq. 3) makes all the difference (except at the initial time t_0
140 of the contamination $X(t_0) = 1$) with respect to the exceptional case $b = 1$ (Eq.1): the local times $T(t)$
141 and $T_2(t)$ are no longer invariants of the contamination dynamics, but increase with t (for $b < 1$).
142 This merely corresponds to the fact that the system became less and less efficient due to its larger and
143 larger outer scale, which rather makes sense. As expected, this increase of $T(t)$ yields a slower
144 growth of $X(t)$ with respect to the exponential model.

145

146 Injecting the time $T(t)$ in Eq.1 yields its scaling generalisation:

$$147 \quad dX/dt \approx aX(t)^b \quad (4)$$

148 where b can be now be understood as the feedback exponent from the cumulative incidence $X(t)$
149 onto the elementary increment $dX(t) = X(t) - X(t - dt)$, both are thus non stationary. This is a
150 common feature for growth processes²⁹, but not for cascades, which are usually supposed to have a
151 stationary outer scale. We therefore need to proceed to a “cumulative-incremental analysis” to
152 quantify this feedback, mainly to determine its exponent b . The simplicity of this analysis might help
153 to change the human behaviour with respect to epidemic risks, whose importance has been often
154 underlined¹⁴.

155 The very first step (see Fig.2) corresponds to simply graph $\Delta X(t_i)$ vs. $X(t_i)$ in a log-log plot, where
156 $\Delta X(t_i) = X(t_i) - X(t_{i-1})$. It is rather surprising to see how much the curves (corresponding to the
157 various entities recalled above) collapse together providing a first indication of a given universality
158 of the virus dynamics over a range of time scales (see Supplementary Information). Furthermore, the
159 range over which this universality is better observed corresponds to power-laws $\Delta X(t_i) \approx aX(t_i)^b$,
160 i.e., a discrete versions of Eq.3. This scaling behaviour is supported by least square fits of the
161 cumulative-incremental incidence curves (dashed straight lines in the log-log plot of Fig.2) with high
162 determination coefficients R^2 and limited dispersion of the estimates throughout the time scales (see
163 Supplementary Information).

164

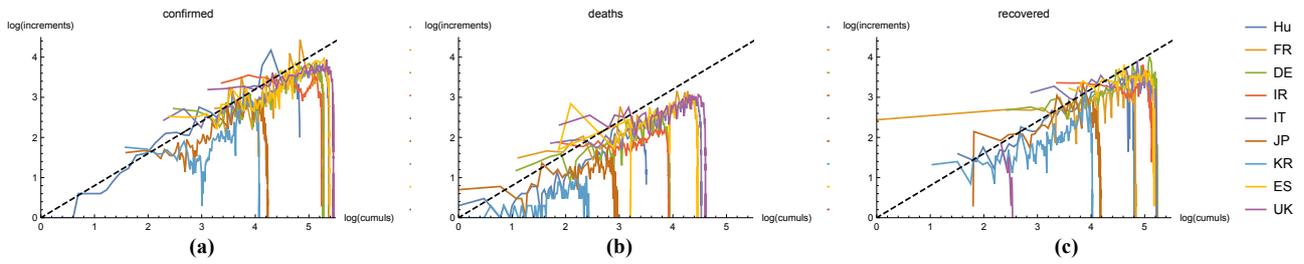


Fig.3: Cumulative-incremental analysis of the decline phase: log-log plots of the couples $(X_{max} - X(t), \Delta X(t))$ to analyse the decline phase of (a) confirmed, (b) deaths, (c) recovered incidences cumulative $X(t)$ of the aforementioned entities (Hubei, France, Germany, Iran, Italy, Japan, Korea, Spain and United Kingdom), over the period 22 January – 07 June, with an exception of shorter period for Hubei till 16 April. Contrary to those of Fig.2, these graphs should be read from right to left to follow the time arrow ($X_{max} - X(t)$ being non increasing with time) and the almost vertical parts on the right hand side corresponds to compressed views of growth phases, which were expanded in Fig.2.

The reference curve (dashed black straight line) corresponds to the equation $\Delta X(t) = (X_{max} - X(t))^{\bar{b}}$, where \bar{b} denotes the average value of the feedback exponents b 's of the growth phase analysis (Fig.2). Hubei (dark blue) and Korea (light blue) rather clearly confirm a feedback exponent $b' \approx \bar{b}$, while other time series were beginning their decline phase. The loop exhibited by the time series of Spain in the category “deaths” is due to a sharp decrease (–1918 cases) of this cumulation on 25 May. All the posterior daily cumulation remain inferior to the cumulation of 24 May. There are similar behaviours, although less obvious due to smaller adjustments, for some other countries.

165

166 The most important result is that the estimates of the feedback exponent b are clearly below 1 for all
 167 types of time series over the period 22 January to 04 April: 0.81 ± 0.09 (confirmed cases), $0.68 \pm$
 168 0.25 (deaths), 0.78 ± 0.13 (recovered). These estimates are mainly obtained on growth phases,
 169 however as displayed by Fig.3, their mean values fit rather well the corresponding decline phases (see
 170 Supplementary Information).

171

172 The local time $T(t)$ (Eq.3) increases linearly with respect to time, with the prefactor $(1 - b)$. Indeed,
 173 a straightforward integration of Eq.4 yields:

$$174 \forall t_1, t_2: t_2 - t_1 = (X(t_2)^{1-b} - X(t_1)^{1-b}) / (a(1 - b)) = (T(t_2) - T(t_1)) / (1 - b) \quad (5)$$

175 The empirical, local times $T(t)$ (defined by Eq.3 and normalised by $(1 - b)$) are in agreement with
 176 this relation in the growth phase (see Supplementary Information).

177

178 Decline phase analysis

179 The second stage of the cumulative-incremental analysis is to explore the decline phase. By reversing
 180 the process $X(t)$ from its maximal value $X_{max} = X(t_{max})$, we obtain a process $X'(t')$:

$$181 X'(t') = X_{max} - X(t); t' = t_{max} - t \quad (6)$$

182 whose growth phase corresponds to the decline phase of $X(t)$ and vice-versa. As $dX/dt = dX'/dt'$,
 183 the scaling decline phase is ruled by a slight modification of Eq.4:

$$184 dX/dt \approx a'(X_{max} - X(t))^{b'} \quad (7)$$

185 Thanks to this symmetry, the cumulative-incremental analysis remains rather the same: it analyses
 186 the (discrete) graph $\Delta X(t_i)$ vs. $(X_{max} - X(t_i))$, seeking a power law relation for large times t_i 's.
 187 Figure 3 shows an overall agreement with $b' \approx b$, particularly for the Hubei province that has a fully
 188 developed decline phase. For other countries the total cumulation is weakly approached by the
 189 empirical X_{max} , resulting in a temporary flattening of last points (see Supplementary Information).

190

191 Discussion and conclusions

192 So far, we have been using only deterministic calculus, whereas Figs. 2, 3 display non negligible
 193 fluctuations with respect to a power-law. Not all of these fluctuations result either from data

194 uncertainty, or from quasi-Gaussian perturbations. These deviations can be easily accounted for in
195 the framework of stochastic multiplicative cascades³⁰. These processes are exponentials of a
196 stochastic generators $\Gamma(t)$ that broadly generalise $\ln(t - t_0)/(b - 1)$. However, their stochastic
197 differentiation remains rather close to the deterministic one, mostly adding a supplementary term σdt
198 to the differential of the generator:

$$199 \quad dX(t) \approx X(t) d\Gamma(t) + \sigma dt \quad (8)$$

200 When the generator $\Gamma(t)$ is Gaussian, σ is the “quadratic variation”, extensions to Lévy stable
201 generators, which are strongly non-Gaussian, have been also considered⁸.

202
203 Therefore, the next stage of the cumulative-incremental analysis is to investigate the stochastic nature
204 of the cascade generator, as well as how it respects the aforementioned symmetry between the growth
205 and decline phases (Eq. 6). Universal multifractals³⁰ can provide some preliminary, universal scaling
206 characterisation, despite the present sample limitations (see Supplementary Information). The present
207 results paved the way for joint scaling analysis of the vector-valued time series $(X_1(t), X_2(t), X_3(t))$,
208 instead of separately analysing the time series corresponding to the three categories, and to introduce
209 location and other dependencies. This would correspond to enlarge the domain, on which the process
210 and its generator are defined, as well as their co-domain, on which they are valued⁸.

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270

271 **Supplementary Information**

272 is linked to the online version of the paper at <https://www.nature.com/>.

273

274 **Acknowledgments**

275 This work is supported by the Academic Chair “Hydrology for Resilient Cities” of Ecole des Ponts
276 ParisTech, endowed by Veolia.

277

278 **Author contributions**

279 Both authors equally contributed to the paper and therefore the alphabetical order is used.

280

281 **Conflict of interest**

282 There is no conflict of interest

283

284 **Data and results duplication**

285 As explained in the text, only the open data of the Center for Systems Science and Engineering at
286 Johns Hopkins University have been used (https://github.com/CSSEGISandData/COVID-19/tree/master/csse_covid_19_data). Furthermore, authors are willing to provide on request the
287 standalone Mathematica program that was used to perform these analyses and to generate a large
288 number of figures, from which were selected the few presently presented in this paper and its
289 Supplementary Information.

290

Figures

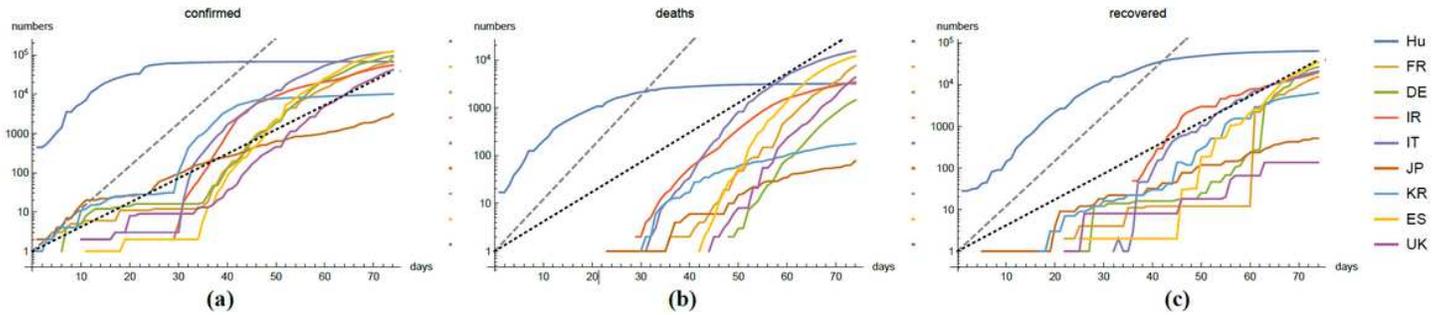


Figure 1

Sub-exponential growth phase: log-linear plots of (a) confirmed, (b) death, (c) recovered cumulative incidences $\mathbb{I}(t)$ of the following entities: Hubei, France, Germany, Iran, Italy, Japan, Korea, Spain and United Kingdom, over the period 22 January - 04 April. An exponential behaviour would correspond to straight lines like those drawn for characteristic times $\tau = 4; 7$ days, and therefore to doubling times $\tau = \tau \ln(2) \approx 3; 5$ days (respectively dashed grey and dotted black straight lines). Due to their (strong) concavity, all the trajectories are sub-exponentials, i.e., grow slower than any of their local approximations by an exponential that corresponds to a tangent to these log-linear graphs.

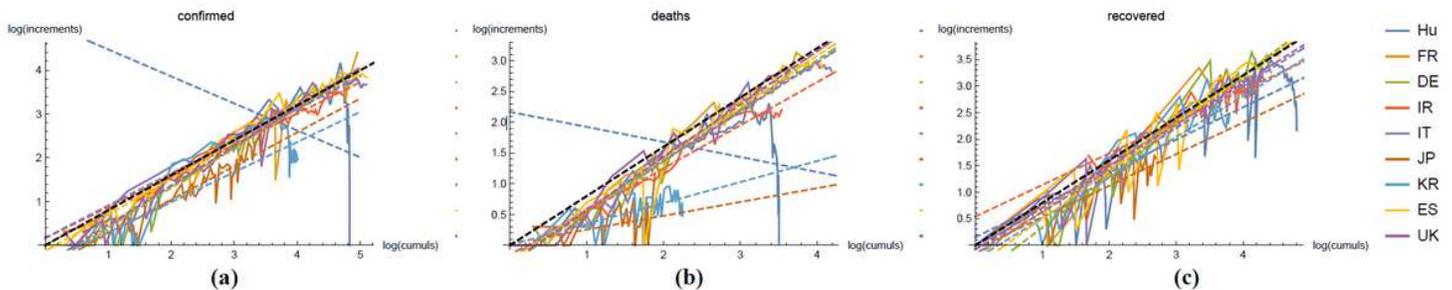


Figure 2

Cumulative-incremental analysis of the growth phase: log-log plots of the couples $(\mathbb{I}(t), \Delta\mathbb{I}(t))$, where $\Delta\mathbb{I}(t) = \mathbb{I}(t) - \mathbb{I}(t - \Delta t)$ is the increment (the time increment Δt being one day) of (a) “confirmed”, (b) “deaths”, (c) “recovered” cumulative incidence $\mathbb{I}(t)$ of the following entities: Hubei, France, Germany, Iran, Italy, Japan, Korea, Spain and United Kingdom, over the period 22 January- 04 April. Graphs should be read from left to right to follow the time arrow ($\mathbb{I}(t)$ being non decreasing with time) and the almost vertical parts on the righthand side corresponds to compressed views of decline phases to be expanded in Fig.3. These graphs display a rather universal behaviour, especially a common scaling/power-law behaviour $\Delta\mathbb{I}(t) \approx \mathbb{I}(t)^\alpha$, a discrete version of the differential equation Eq. 4, over a non negligible range of couples $(\mathbb{I}(t), \Delta\mathbb{I}(t))$ and therefore of time, with only a few exceptions. The best linear fits are represented by the dashed lines (of the same colour as the graph of this entity), with high determination coefficients R^2 and limited dispersion of the estimates of the feedback exponent α (see Supplementary

Information). The black dashed line indicates the average estimates over the period of 22 January - 04 April.

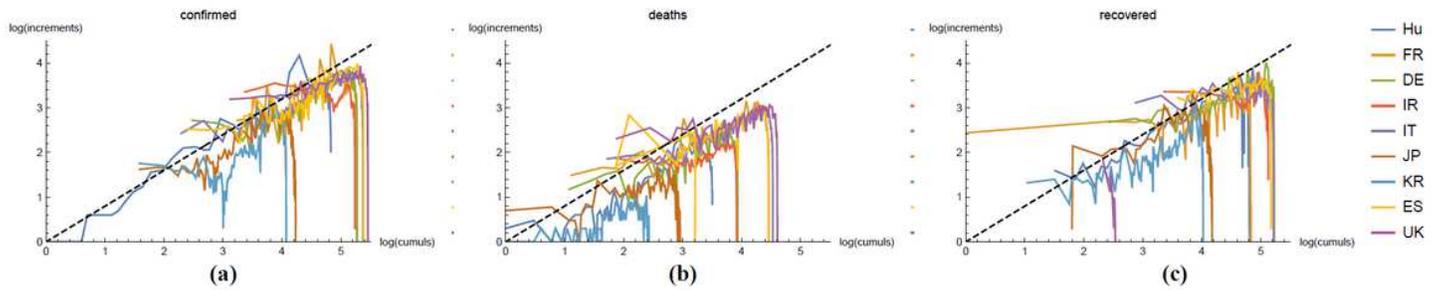


Figure 3

Cumulative-incremental analysis of the decline phase: log-log plots of the couples $(\log(\max - X(t)), \Delta \log(X))$ to analyse the decline phase of (a) confirmed, (b) deaths, (c) recovered incidences cumulative $X(t)$ of the aforementioned entities (Hubei, France, Germany, Iran, Italy, Japan, Korea, Spain and United Kingdom), over the period 22 January – 07 June, with an exception of shorter period for Hubei till 16 April. Contrary to those of Fig.2, these graphs should be read from right to left to follow the time arrow ($\log(\max - X(t))$ being non increasing with time) and the almost vertical parts on the right hand side corresponds to compressed views of growth phases, which were expanded in Fig.2. The reference curve (dashed black straight line) corresponds to the equation $\Delta \log(X) = (\log(\max - X(t)))^\alpha$, where α denotes the average value of the feedback exponents α 's of the growth phase analysis (Fig.2). Hubei (dark blue) and Korea (light blue) rather clearly confirm a feedback exponent $\alpha \approx 1$, while other time series were beginning their decline phase. The loop exhibited by the time series of Spain in the category "deaths" is due to a sharp decrease (-1918 cases) of this cumulation on 25 May. All the posterior daily cumulation remain inferior to the cumulation of 24 May. There are similar behaviours, although less obvious due to smaller adjustments, for some other countries.

Supplementary Files

This is a list of supplementary files associated with this preprint. Click to download.

- [COVID19cascadeSupplementaryInformation20201106.pdf](#)