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Floppy hydrated salt microcrystals

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Summary

The crystalline structure of minerals due to the highly ordered assembly of its constituent atoms, ions or molecules confers a considerable hardness and brittleness to the materials^{1,2}. As a result, they are generally subject to fracture. Here we report that microcrystals of natural inorganic salt hydrates such as sodium sulfate decahydrate ($\text{Na}_2\text{SO}_4 \cdot 10\text{H}_2\text{O}$) and magnesium sulfate hexahydrate ($\text{MgSO}_4 \cdot 6\text{H}_2\text{O}$) can behave remarkably differently: instead of having a defined faceted geometrical shape and being hard or brittle, they lose their facets and become soft and deformable when in contact with their saturated salt solution at their deliquescence point. As a result, the hydrated crystals simultaneously behaves as crystalline and liquid-like. We show that the observed elasticity is a consequence of a trade-off between the excess capillary energy introduced by the deformation from the equilibrium shape and viscous flow over the surface of the microcrystals. This surprising, unusual mechanical properties reported here reveals the role of the water molecules present in the crystalline structure. Although many compounds can incorporate water molecules in their crystalline frameworks, the relationship between the different hydrates and anhydrate crystal forms are still poorly investigated. Our results on the floppy behaviour of such crystalline structure reveals some unexplored properties of water molecules entrapped in the crystalline structure and can open novel routes for their application in various fields such as pharmaceutical sciences³, thermal energy storage⁴⁻⁶ and even the traceability of water on Mars⁷⁻⁹

Main text

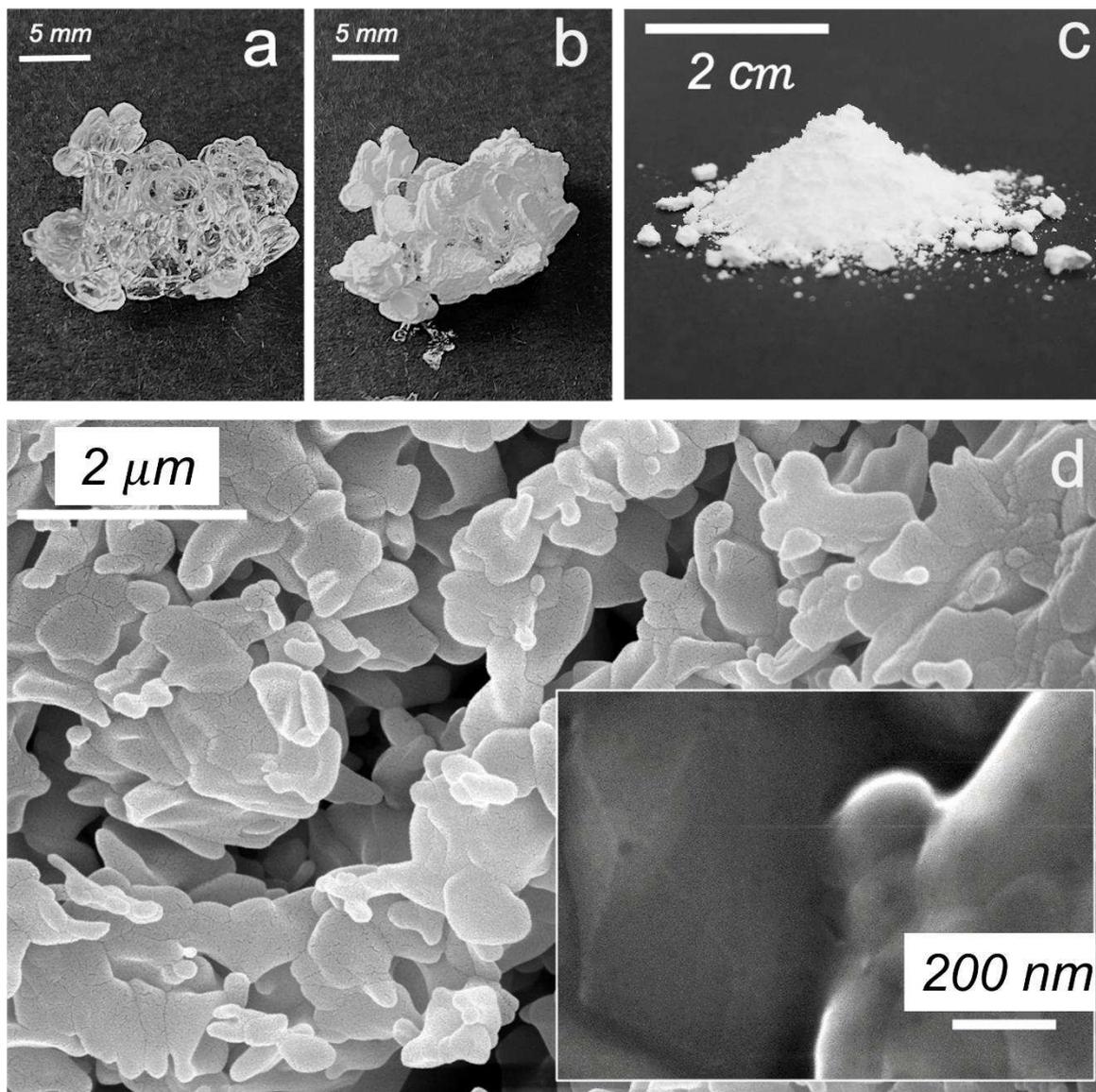
Minerals are natural materials that make up a large part of the earth and other planets⁷⁻⁹. Their chemical formula giving the types and proportions of the chemical elements and their crystal lattice defining the arrangement of atoms and the nature of bonding determine generally their physical properties. These physical properties and some other exceptional discovered properties of minerals such as fluorescence, magnetism or radioactivity, etc are of great technological significance for scientists and engineers for the design of novel materials.

For a typical crystalline solid, such as a salt, the spacing is of the order of a few ångström and the interaction energy several $k_B T$, which immediately leads to the conclusion that its elastic modulus is many GPa¹. The consequence of this large modulus is that, when large stresses are applied to them, the crystal will break, preferentially along directions/planes corresponding to the lowest bond energy or along random directions^{1,2}. It is well known that, due to the presence

48 or absence of defects, the mechanical properties at small scales can differ from the macroscopic
49 ones. However these effects are usually only visible at the nanoscale ¹⁰.

50
51 Here we show that microcrystals of natural inorganic salt hydrates such as sodium sulfate
52 decahydrate ($\text{Na}_2\text{SO}_4 \cdot 10\text{H}_2\text{O}$) and magnesium sulfate hexahydrate ($\text{MgSO}_4 \cdot 6\text{H}_2\text{O}$) can
53 behave remarkably differently instead of being hard, brittle and faceted, they become soft and
54 deformable and lose their facets at their deliquescence point. In addition, we report 'self-
55 healing' properties as indentations on the crystal surface spontaneously and rapidly close
56 themselves under the action of surface tension. This unusual behavior appears to be related to
57 the number of water molecules present in the crystalline structure and their mobility; in some
58 hydrated salts the water is stoichiometrically bound into a crystal but not directly bonded to the
59 central metal cation.

60
61 We study the formation of sodium sulfate decahydrate ($\text{Na}_2\text{SO}_4 \cdot 10\text{H}_2\text{O}$) and magnesium sulfate
62 hexahydrate ($\text{MgSO}_4 \cdot 6\text{H}_2\text{O}$) by the deliquescence of the anhydrous form of these salts using
63 electron and optical microscopy as well as confocal Raman micro spectroscopy. Deliquescence
64 is the absorption of water vapor by the crystals from the atmosphere until it dissolves and forms
65 a solution. First, the anhydrous form of sodium sulfate is obtained by drying millimetric
66 mirabilite crystals precipitated in salt solution at room temperature. Mirabilite or the sodium
67 sulfate decahydrate also referred to as Glauber's salt ¹¹, is the thermodynamically stable state
68 in contact with a saturated sodium sulfate solution at room temperature (**Extended Data Fig.**
69 **1**) ¹². However, once the crystals are taken out of solution, the water in the crystalline structure
70 evaporates easily at room temperature, and subsequently the initially transparent hydrated
71 crystal transforms into a white porous structure made up of micro crystallites, while retaining
72 the overall shape of the original millimeter-sized mirabilite crystal (**Fig. 1a,b**). Scanning
73 Electron microscopy imaging shows that this dehydrated white porous structure is made of an
74 assembly of microcrystals of thenardite (**Fig. 1d**), as identified by confocal Raman microscopy
75 with an SO_4^{2-} -symmetric stretch peak at 994 cm^{-1} (**Extended Data Fig. 2**) in agreement with
76 previous work ^{13,14}. It is this porous structure that is subsequently rehydrated by deliquescence.
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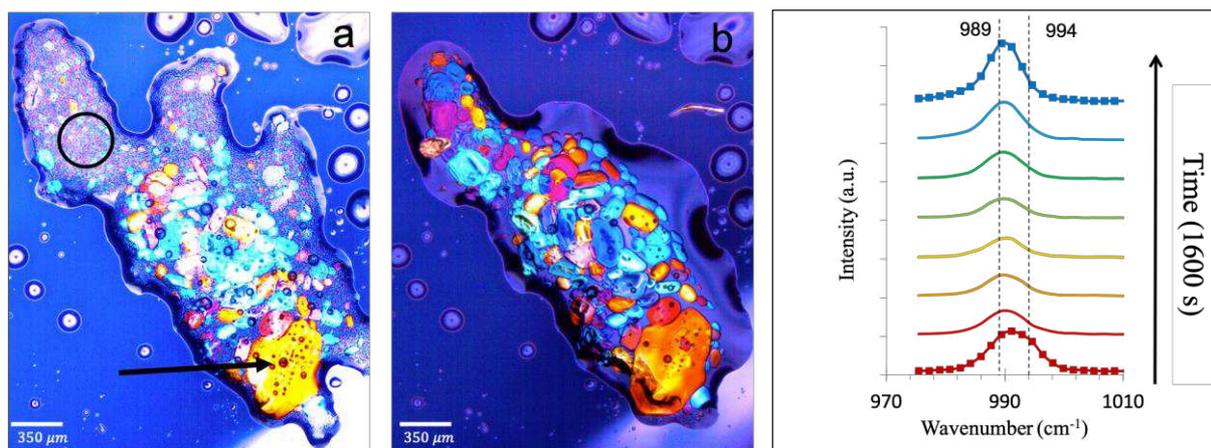


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79 **Fig. 1. Mirabilite ($\text{Na}_2\text{SO}_4 \cdot 10\text{H}_2\text{O}$) and thenardite (Na_2SO_4) crystals.** (a) typical mirabilite
 80 crystals precipitated out of a sodium sulfate solution, and (b) mirabilite crystals dried
 81 (dehydrated) at $T=21^\circ\text{C}$ (c) Powder of the dried crystals (d) Electron microscopy images of the
 82 dried crystal consisting of an assembly of nanocrystals of thenardite.

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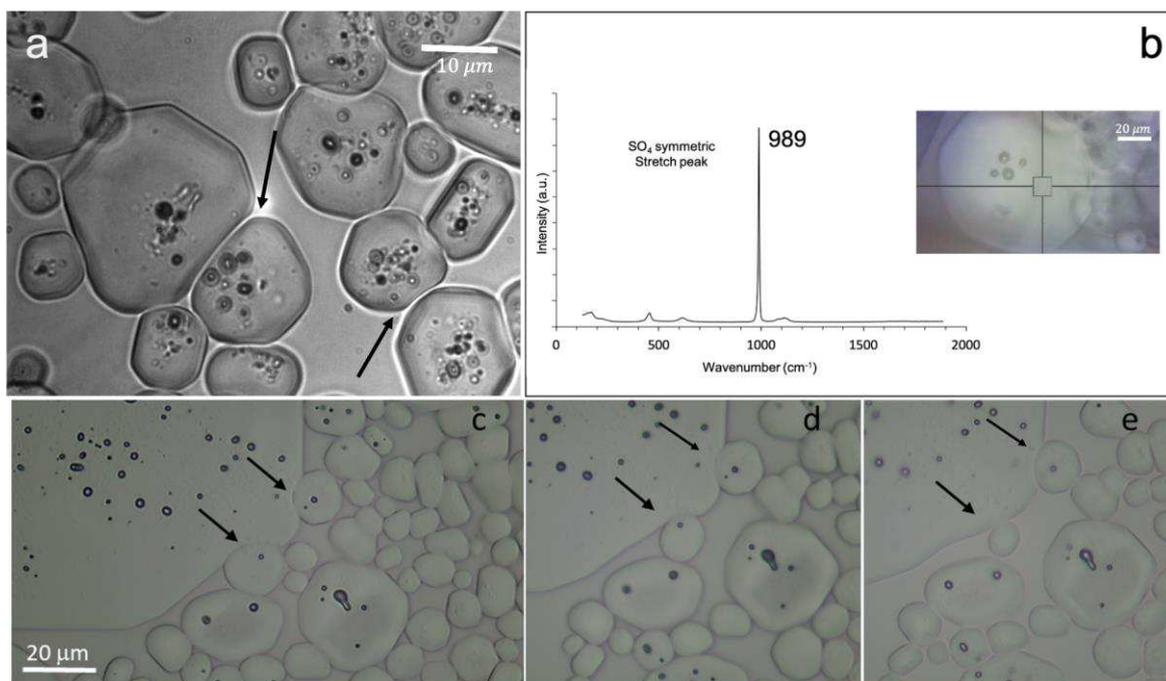
84 To do so, a small amount of the anhydrous powder is placed under a polarizing light microscope
 85 in a climate chamber at a relative humidity above the equilibrium relative humidity $\text{RH}_{\text{eq}} \sim 96\%$
 86 of this salt (**Extended Data Fig. 1**)⁴. The progressive moisture sorption and dissolution of the
 87 thenardite nanocrystals results in the gradual formation of a salt solution in which subsequently
 88 the growth of new crystals of mirabilite can be observed (**Fig. 2a,b**). This is confirmed by
 89 analysis of sequential Raman spectra of the process, that reveals an overlay of the thenardite
 90 (994 cm^{-1}) and mirabilite (989 cm^{-1}) spectra; the intensity of the anhydrous polymorph
 91 (thenardite) decreases in time due to the deliquescence, whereas that of the mirabilite spectrum
 92 increases with the formation and growth of more mirabilite microcrystals (**Fig. 2c**); the latter
 93 can also entrap undissolved thenardite as impurities during their growth (seen as black dots in
 94 **Fig. 2**) in their structures during the growth. The process thus leads to have large amount of
 95 microscale size mirabilite crystals instead of a few large millimetric ones as seen in **Fig1a**.



97
98 **Fig. 2. Gradual deliquescence of thenardite (Na_2SO_4) polycrystals and rapid precipitation**
99 **of mirabilite ($\text{Na}_2\text{SO}_4 \cdot 10\text{H}_2\text{O}$) micro crystals. (a-b)** Microscopy images made with cross
100 cross polarizers of the deliquescence process of the dehydrated powder. The colors that are seen with
101 cross polarized lights highlights the fact that what we are seeing is not a liquid structure.
102 Initially, both polymorphs (mirabilite and thenardite) coexist as shown by the Raman spectra
103 of the region circled in black (b) After 660 s, most of the thenardite nanocrystals are dissolved;
104 further water vapour absorption will cause the mirabilite crystals to deliquesce in turn. The
105 arrow shows entrapped remaining thenardite nanocrystals (black dots) in the mirabilite crystals.
106 (c) Raman spectra of the evolution of the SO_4 symmetric stretch peak during deliquescence of
107 thenardite powder; the dotted lines indicate the peak maxima of thenardite (994 cm^{-1}) and
108 mirabilite (989 cm^{-1}) (see **Extended Data Fig. 2**).
109

110 Generally, the higher the number of water molecules in the crystalline structure of a given
111 inorganic hydrated crystal, the lower its solubility at a given temperature (phase diagrams from
112 the literature are shown in **Extended Data Fig. 1**). 10,13,14,15 In our experiment, the salt
113 solution formed by deliquescence of thenardite nanocrystals is therefore highly concentrated
114 with respect to the formation of the decahydrate form (**Extended Data Fig. 1**); the resulting
115 supersaturation ($S \sim m_{\text{thenardite}} / m_{\text{mirabilite}} \sim 2.5$ at 21°C) induces the rapid precipitation of the
116 large amount of mirabilite microcrystals. Their spectacular-colored appearance in polarizing
117 light microscopy (**Fig. 2**) and Raman identification of these crystals as mirabilite (**Fig. 3b**) rules
118 out the formation of an intermediate metastable amorphous phase or dense liquid droplets on
119 the time scale of our experiments. The mirabilite microcrystals are observed to grow only until
120 a certain size, until the surrounding salt solution concentration drops to the saturation
121 concentration for mirabilite (**Extended Data Fig. 1**). Beyond this threshold, the precipitated
122 mirabilite crystals will then dissolve in their turn, as the deliquescence process (i.e., the water
123 vapour uptake) continues.
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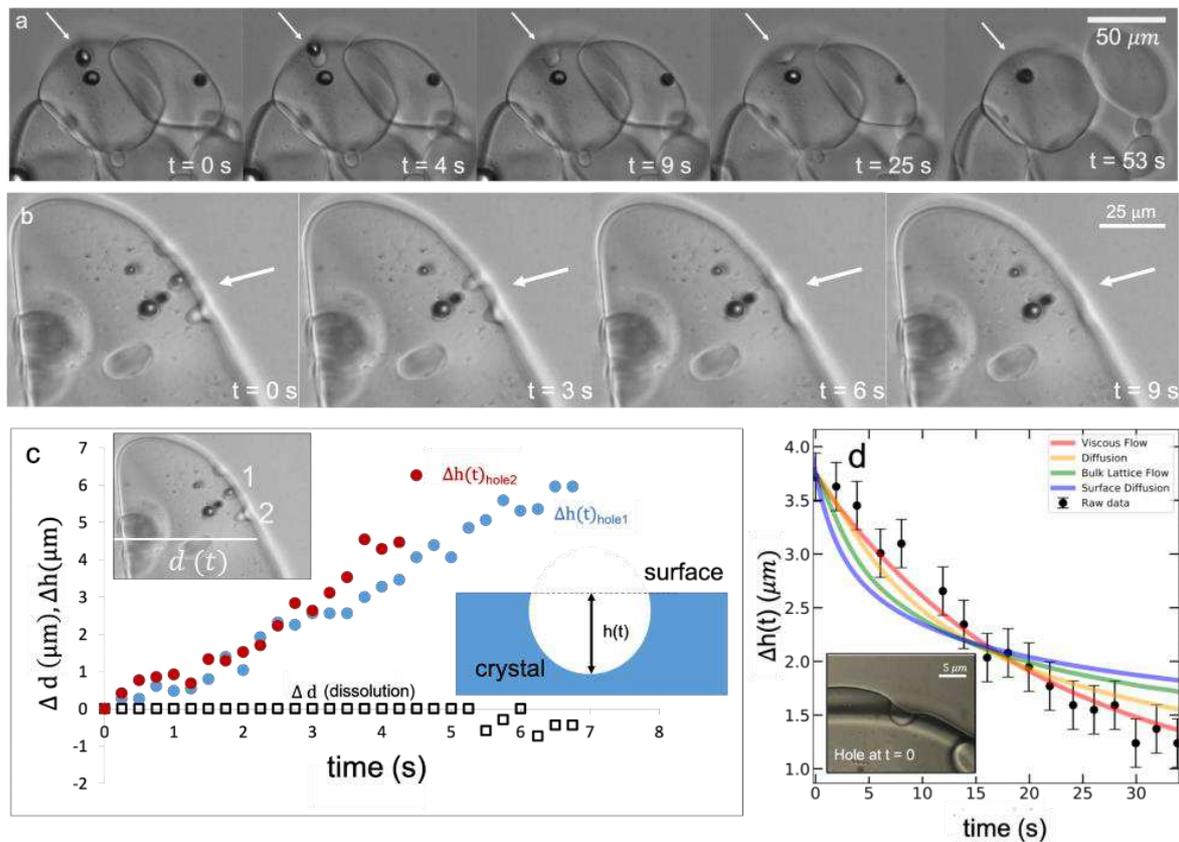
125 The remarkable observation is that, although mirabilite has a clear monoclinic crystal structure,
126 the precipitated small mirabilite crystals close to their deliquescence point exhibit no facets at
127 all and appear very rounded. In addition, they are observed to be soft; this can be most clearly
128 observed in regions with a high local density of crystals: the microcrystals are observed to
129 deform each other and adapt to the shape of the other neighboring crystals, much like liquid
130 droplets in an emulsion (**Fig. 3a**). Once the strain imposed by the presence of other crystals is
131 removed, the deformed crystals regain their non-faceted spherical shape (**Fig. 3c-e**).



132
 133 **Fig. 3. Floppiness of mirabilite micro crystals surrounded by sodium sulfate solution at**
 134 **their deliquescence point. (a)** At the deliquescence point, non-faceted mirabilite micro crystals
 135 surrounded by sodium sulfate solution are deformable and can adapt their shape with respect to
 136 the neighboring crystals. **(b)** Raman spectrum of the crystals identified as mirabilite. **(c)** the
 137 arrows show the local deformation of a large mirabilite crystal by two neighboring small
 138 crystals like pushing on a soft ball; **(d-e)** The crystals show an elastic response once the strain
 139 is removed indicated with the black arrows.

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 141 One way to interpret these results would be to assume that the mobility of the molecules at the
 142 crystalline surface is very high. In agreement with this hypothesis, we observe in a separate set
 143 of experiments that any indentations at the crystal surface can heal spontaneously. Such
 144 indentations appear during the deliquescence of the mirabilite microcrystals when entrapped
 145 thenardite nanocrystals (seen as black dots in **Fig. 2 and 3**) are released into the solution and
 146 quickly dissolve in the surrounding salt solution, which is not fully saturated with respect to the
 147 thenardite polymorph (**Fig. 4a,b**). We observe that the holes thus formed close up
 148 spontaneously.

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Fig. 4. Self-healing behavior at the surface of mirabilite crystals. (a-b) evolution of surface defects on two different crystals, indicated by arrows. During the deliquescence, while the entrapped thenardite dissolves, the indentation at the surface of the hydrated crystal is seen to level out rapidly. (c) The temporal evolution of the levelling out $\Delta h(t) = h(t) - h_0$ of two holes (filled circles) and the dissolution $\Delta d(t) = d(t) - d_0$ of the same crystal (open squares) (d) Comparison between the measured depth of a hole $h(t)$ levelling towards a flat surface (symbols) and the Mullins model (solid lines) describing the four possible mechanisms explaining the levelling. The model describing a viscous flow (red curve) shows the best agreement with the data.

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One can imagine that the levelling of the holes could simply be due to the dissolution of the mirabilite crystal because of the deliquescence. However, the comparison of the dissolution speed of the crystals $d(t)$ with the temporal evolution of the depth of several holes $h(t)$ of different sizes on various mirabilite crystals shows that the dissolution is much too slow to account for the disappearance of the holes (Fig. 4c). The self-healing behaviour of the crystal surface has thus another physical origin. To determine what mechanism is responsible for the hole closing, we use Mullins model¹⁷⁻¹⁹ which solves the levelling of a hole due to four different possible mechanisms: (1) viscous-capillary flow of the crystal, (2) diffusion of the salt from the solution into the holes, (3) a bulk lattice rearrangement, and (4) a rearrangement of surface molecules by surface diffusion. We compare this model to the full experimental hole profiles measured during the levelling in (Extended Data Fig. 3) and to the hole depth (Fig. 4d) and find that the viscous-capillary flow mechanism is the only one able to describe the data (for details on the model, equation and the comparison with experiments, see Materials and Methods in the Extended Data). The constant in the equation represents a capillary velocity; its value is approximately $0.1 \mu\text{m s}^{-1}$ for the hole shown in Figure 4d and does not vary much between different experiments. Therefore, the self-healing, i.e., the relaxation of the holes is driven by

179 the gradients in Laplace pressure along the crystal-solution interface; the excess capillary
180 energy introduced by the curvature of the hole with respect to the flat equilibrium state is
181 dissipated by viscous flow into the hole that leads to the observed self-healing. The driving
182 force for the healing is thus clearly the surface tension that requires the crystals to be spherical.
183

184 This unusual soft, non-faceted and self-healing behavior at the deliquescence point is not
185 specific to sodium sulfate. Repeating the same experiments for micro crystals of magnesium
186 sulfate hexahydrate ($\text{MgSO}_4 \cdot 6\text{H}_2\text{O}$), we observed a similar behavior under the polarizing and
187 Raman confocal microscope (**Extended Data Fig. 4**). The process is, however, much faster
188 since the deliquescence of magnesium sulfate starts at a lower equilibrium relative humidity
189 $\text{RH}_{\text{eq}} = 85\%$ ¹⁶, which induces a faster water vapor sorption rate in comparison with sodium
190 sulfate in our climate chamber (where the relative humidity of the chamber was fixed at
191 $\text{RH} \sim 100\%$). In addition, the magnesium sulfate hexahydrate microcrystals were smaller in size,
192 which prevented us from measuring the self-healing process quantitatively with our
193 experimental setup as the indentations were too small and closed too rapidly. Nonetheless, the
194 magnesium sulfate salt behaved qualitatively like the sodium salt described above.
195

196 Our observations of the self-healing and softness of hydrated crystals are reminiscent of the
197 case of Helium-4 crystals close to zero temperature, which under the application of extremely
198 small stresses present a giant plasticity due to the motion of crystalline defects. This leads,
199 amongst other things, to the ‘dripping’ of such crystals reported to behave liquid-like rather
200 than solid-like^{20–22} due to the bulk lattice motion (‘model 3’ in our terminology)^{20–22}

201 In our case of hydrated salts, the deformability of the microcrystals at the deliquescence point
202 is mainly caused by capillary-driven viscous flow at the surface of the crystal. This governs the
203 self-healing dynamics of the crystals and shows why the molecules in the crystalline lattice
204 appear to be mobile, leading to spherical and deformable crystals on our experimental time
205 scale. The time scale for the viscous flow limits the size range over which the floppiness can
206 be observed. The comparison of the behavior between the sodium and magnesium salt is
207 interesting; one hypothesis could be that the higher numbers of water molecules in the hydrated
208 crystal, water can be more easily exchanged between the crystal and the layer adjacent to it.
209 Second, since the magnesium is bivalent, the water is likely to be more strongly bound, again
210 decreasing the exchange rate with the exterior water. The ease of exchange of the water
211 molecule in the crystalline structure with the surrounding is also indicated by the fact that this
212 water of hydration evaporates more readily from hydrated crystals with a relatively high number
213 of water molecules^{4,5,23}. For the sodium sulfate decahydrate this readily happens at room
214 temperature, whereas for hydrated salts with fewer water molecules in the lattice the water
215 evaporates at higher temperatures (for magnesium sulfate hexahydrate $T > 60^\circ\text{C}$, for magnesium
216 sulfate monohydrate and calcium sulfate dihydrate $T > 150^\circ\text{C}$). The same reasoning then also
217 explains why the softness is not observable for the anhydrous form of the salt.

218 Another important observation in this respect is that when the relative humidity is decreased
219 below the equilibrium relative humidity RH_{eq} (i.e., below the deliquescence point), the
220 unfaceted hydrated crystals will lose their softness, facets appear again and increase in size over
221 time. Thus, the surface mobility only appears at the deliquescence point (i.e., at a certain RH
222 for a given salt); perhaps in analogy to the roughening transition that appears at a fixed
223 temperature for certain crystals, and above which the facets also disappear.
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225
226 The hydrated salts that we study here are not only abundant and ubiquitous on the surface of
227 the Earth but are also present on Mars^{8,9,15}. In fact, their discovery on the sediments of Mars
228 provided vital information on the hydrogeologic history of other planets. In addition, they are

229 an important class of natural materials that are currently investigated both as thermochemical
230 materials and phase change materials (PCM) and for energy storage purposes because of their
231 large potential storage capacity ($\sim 1-3 \text{ GJ/m}^3$) and the easily accessible temperature range that
232 they cover. Our results reveal an unexplored fundamental property of inorganic hydrated salts,
233 which can open new routes for optimizing their performance with improved control over their
234 geometry and size. On the other hand, such materials have tremendous potential for
235 applications where the key parameter is relative humidity rather than temperature. More
236 generally, the conception of novel *soft microcrystalline materials* with peculiar macroscopic
237 properties with respect to their shape, colour, and luminescence as ‘flexible responsive
238 systems’ remains an underdeveloped area that deserves further exploration.²⁴⁻²⁶
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Main references

1. Kittel, C. *Introduction to Solid State Physics*. (2005).
2. Ashcroft, N. W. & Mermin, N. David. *Solid state physics*. (Holt, Rinehart and Winston, 1976).
3. Healy, A. M., Worku, Z. A., Kumar, D. & Madi, A. M. Pharmaceutical solvates, hydrates and amorphous forms: A special emphasis on cocrystals. *Advanced Drug Delivery Reviews* vol. 117 25–46 (2017).
4. Donkers, P. A. J., Linnow, K., Pel, L., Steiger, M. & Adan, O. C. G. Na₂SO₄·10H₂O dehydration in view of thermal storage. *Chemical Engineering Science* **134**, 360–366 (2015).
5. Donkers, P. A. J. *et al.* Water Transport in MgSO₄·7H₂O during Dehydration in View of Thermal Storage. *Journal of Physical Chemistry C* **119**, 28711–28720 (2015).
6. Donkers, P. A. J., Sögütöglu, L. C., Huinink, H. P., Fischer, H. R. & Adan, O. C. G. A review of salt hydrates for seasonal heat storage in domestic applications. *Applied Energy* vol. 199 45–68 (2017).
7. Ojha, L. *et al.* Spectral evidence for hydrated salts in recurring slope lineae on Mars. *Nature Geoscience* **8**, 829–832 (2015).
8. Vaniman, D. T. *et al.* Magnesium sulphate salts and the history of water on Mars. *Nature* **431**, 663–665 (2004).
9. Wang, A., Freeman, J. J., Jolliff, B. L. & Chou, I. M. Sulfates on Mars: A systematic Raman spectroscopic study of hydration states of magnesium sulfates. *Geochimica et Cosmochimica Acta* **70**, 6118–6135 (2006).
10. Meyers, M. A., Mishra, A. & Benson, D. J. Mechanical properties of nanocrystalline materials. *Progress in Materials Science* vol. 51 427–556 (2006).
11. Levy, H. A. & Lisenskyt, G. C. *Crystal Structures of Sodium Sulfate Decahydrate (Glauber's Salt) and Sodium Tetraborate Decahydrate (Borax). Redetermination by Neutron Diffraction**. *Acta Cryst* vol. 34 (1978).
12. Steiger, M. & Asmussen, S. Crystallization of sodium sulfate phases in porous materials: The phase diagram Na₂SO₄-H₂O and the generation of stress. *Geochimica et Cosmochimica Acta* **72**, 4291–4306 (2008).
13. Hamilton, A. & Menzies, R. I. Raman spectra of mirabilite, Na₂SO₄·10H₂O and the rediscovered metastable heptahydrate, Na₂SO₄·7H₂O. *Journal of Raman Spectroscopy* **41**, 1014–1020 (2010).
14. Lindström, N., Talreja, T., Linnow, K., Stahlbuhk, A. & Steiger, M. Crystallization behavior of Na₂SO₄-MgSO₄ salt mixtures in sandstone and comparison to single salt behavior. *Applied Geochemistry* **69**, 50–70 (2016).
15. Steiger, M., Linnow, K., Ehrhardt, D. & Rohde, M. Decomposition reactions of magnesium sulfate hydrates and phase equilibria in the MgSO₄-H₂O and Na⁺-Mg²⁺-Cl⁻-SO₄²⁻-H₂O systems with implications for Mars. *Geochimica et Cosmochimica Acta* **75**, 3600–3626 (2011).
16. López-Arce, P., Fort, R., Gómez-Heras, M., Pérez-Monserrat, E. & Varas-Muriel, M. J. Preservation strategies for avoidance of salt crystallisation in El Paular Monastery cloister, Madrid, Spain. *Environmental Earth Sciences* **63**, 1487–1509 (2011).
17. Mullins, W. W. Theory of Thermal Grooving. *Journal of Applied Physics* **28**, 333–339 (1957).
18. Mullins, W. W. Flattening of a nearly plane solid surface due to capillarity. *Journal of Applied Physics* **30**, 77–83 (1959).
19. Backholm, M., Benzaquen, M., Salez, T., Raphaël, E. & Dalnoki-Veress, K. Capillary levelling of a cylindrical hole in a viscous film. *Soft Matter* **10**, 2550–2558 (2014).

- 292 20. Balibar, S., Alles, H., Parshin, A. Y. & Kapitza, P. L. *The surface of helium crystals*.
 293 (2005).
- 294 21. Ishiguro, R., Graner, F., Rolley, E., Balibar, S. & Eggers, J. Dripping of a crystal.
 295 *Physical Review E - Statistical, Nonlinear, and Soft Matter Physics* **75**, (2007).
- 296 22. Balibar, S. *et al.* Dislocations in a quantum crystal Solid helium: A model and an
 297 exception. *Comptes Rendus Physique* vol. 17 264–275 (2016).
- 298 23. Ruben, W. *et al.* Crystal Structure and Entropy of Sodium Sulfate Decahydrate.
 299 *Journal of the American Chemical Society* **83**, 820–824 (1960).
- 300 24. Kato, M., Ito, H., Hasegawa, M. & Ishii, K. Soft Crystals: Flexible Response Systems
 301 with High Structural Order. *Chemistry - A European Journal* **25**, 5105–5112 (2019).
- 302 25. Horike, S., Shimomura, S. & Kitagawa, S. Soft porous crystals. *Nature Chemistry* vol.
 303 1 695–704 (2009).
- 304 26. Gupta, P., Karothu, D. P., Ahmed, E., Naumov, P. & Nath, N. K. Thermally Twistable,
 305 Photobendable, Elastically Deformable, and Self-Healable Soft Crystals. *Angewandte*
 306 *Chemie - International Edition* **57**, 8498–8502 (2018).

307 Figure legends

308 **Fig. 1. Mirabilite ($\text{Na}_2\text{SO}_4 \cdot 10\text{H}_2\text{O}$) and thenardite (Na_2SO_4) crystals.** (a) typical mirabilite
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 310 (dehydrated) at $T=21^\circ\text{C}$ (c) Powder of the dried crystals (d) Electron microscopy images of
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 317 Initially, both polymorphs (mirabilite and thenardite) coexist as shown by the Raman spectra
 318 of the region circled in black (b) After 660 s, most of the thenardite nanocrystals are dissolved;
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 323 mirabilite (989 cm^{-1}) (see **Extended Data Fig. 2**).
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325 **Fig. 3. Floppiness of mirabilite micro crystals surrounded by sodium sulfate solution at
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333 **Fig. 4. Self-healing behavior at the surface of mirabilite crystals.** (a-b) evolution of
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341 explaining the levelling. The model describing a viscous flow (red curve) shows the best
342 agreement with the data.
343

344 **Methods**

345 **Recrystallization of the anhydrous Na₂SO₄ nanocrystals in the form of powder**

346 First, sodium sulfate solutions were prepared by dissolving anhydrous thenardite Na₂SO₄
347 (Sigma Aldrich; 99.8% purity) little by little while stirring in Millipore water ($\rho \approx 18.2$
348 M Ω .cm) till a concentration of 18.6 wt % or 1.6 molal; beyond this concentration the solution
349 becomes turbid due to the precipitation of mirabilite crystals (Na₂SO₄ • 10 H₂O) from the
350 supersaturated solution (concentration at saturation being $m_s = 1.4$ molal). The stirring is then
351 stopped and the bottle is kept at rest for one day at room temperature. Some precipitated
352 mirabilite crystals were subsequently removed from the saturated salt solution and dried at
353 ambient temperature and relative humidity RH~45% of the lab. The stability of the
354 hydrated/anhydrous phase in open air depends on the temperature and relative humidity.
355 Upon drying and dehydration, the transparent mirabilite crystals become white and crumbly
356 while keeping the shape of the initial hydrated crystals (see **Fig. 1ab**). The composition and
357 the morphology of the white structure has been identified by Raman spectroscopy and high-
358 resolution scanning electron microscopy as being an assembly of thenardite nanocrystals.

359

360 **Deliquescence of anhydrous Na₂SO₄ nanocrystals**

361 The anhydrous Na₂SO₄ powder was placed on a clean glass slide and compressed to minimize
362 spaces between the powder particles. The glass slide is then placed into a controlled mini
363 climatic chamber under the microscope and subjected to an increasing relative humidity. The
364 temporal physical changes of the powder due to moisture adsorption was followed under the
365 Leica DM IRM inverted microscope equipped with a 800 x 600 pixels and 8-bit sensitivity
366 CCD camera. The images of the ensuing process were taken at different time intervals at 10x,
367 20x, 40x, and 60x magnifications. The images were subsequently analyzed using the Fiji
368 ImageJ software.

369

370 **Recrystallization of the anhydrous MgSO₄ nanocrystals in the form of powder**

371 A similar procedure as for the recrystallization of the Na₂SO₄ was followed.
372 First, magnesium sulfate solutions were prepared by dissolving anhydrous magnesium sulfate
373 (Sigma Aldrich; 99.5% purity) little by little while stirring in Millipore water ($\rho \approx 18.2$
374 M Ω .cm) to obtain a supersaturated solution (concentration at saturation being $m_s = 3.8$ molal).
375 The stirring was then stopped and the bottle was kept at rest for one day at room temperature.
376 Droplets of the solution were put on a Corning glass slide (borosilicate) and were dried at T =
377 60°C. After ~24 hours the dried droplets were removed from the oven and crushed to get a
378 white powder which was stored at RH = 20% to ensure the MgSO₄ stayed anhydrous. The
379 composition and the morphology of the powder has been identified by Raman spectroscopy
380 and high-resolution scanning electron microscopy as being anhydrous MgSO₄ crystals.

381

382 **Deliquescence of anhydrous MgSO₄ nanocrystals**

383 The anhydrous MgSO₄ powder was placed on a clean glass slide. The glass slide was then
384 placed into a controlled mini climatic chamber under the microscope and subjected to an
385 increasing relative humidity. The temporal physical changes of the powder due to moisture
386 adsorption were followed under the Leica DM IRM inverted microscope equipped with a 800
387 x 600 pixels and 8-bit sensitivity CCD camera. The images of the ensuing process were taken
388 at different time intervals at 10x, 20x, 40x, and 60x magnifications. The images were
389 subsequently analyzed using ImageJ software. The floppy crystals appeared at RH ~ 85%.

390

391 **Confocal Raman spectroscopy**

392 For chemical identification of elements and various polymorphs of Na₂SO₄ and MgSO₄ during
393 the deliquescence, we applied Raman microscopy using a Renishaw InVia spectrometer with

394 a 532 nm laser calibrated with a silicon mono crystal. First, we took reference spectra in the
395 fingerprint area of each compound: the mirabilite $\text{Na}_2\text{SO}_4 \cdot 10\text{H}_2\text{O}$, the anhydrous Na_2SO_4
396 powder obtained after dehydration and 1.4 molal saturated solution, using 20x and 50x
397 magnification objectives, backscattering geometry and 100% laser power. These spectra help
398 to distinguish different phases as the wavenumber position and number of peaks differ for
399 each phase. The peak values of the reference spectra were compared to the literature values
400 and an excellent agreement was found¹².

401 We next took sequential Raman spectra to characterize the full deliquescence process and the
402 transformation of different phases. The deliquescence was carried out in a similar mini
403 climatic chamber as used with the inverted Leica microscope, but with an aluminum slide
404 instead of a glass slide at the bottom. Furthermore, we also did single measurements on the
405 hydrated crystals that appeared during deliquescence.

406 **Statistical test to determine the best fitting model for the hole-levelling out**

407 In order to determine the driving mechanism behind the levelling out of the holes, we
408 compared the data of the measurements with four candidate models: (1) relaxation by a
409 viscous flow, (2) a diffusion process, (3) a bulk lattice flow, and (4) a rearrangement of
410 surface molecules by surface diffusion. The theoretical framework for these mechanisms is
411 given by Mullins.^{17,18} He showed that for a one-dimensional sinusoidal indentation
412 $u_0 \sin(kx)$, the profile evolves as $U(x, t) = u(t) \sin(kx)$, with x the direction perpendicular
413 to the hole, and t the time and $u(t)$ the amplitude which is described by:

$$414 \frac{\partial u}{\partial t} = -C_n(T)k^n u,$$

415
416 Where $C_n(T)$ is a temperature-dependent prefactor, k the wavenumber of the profile and n an
417 integer that represents the levelling physical mechanism:

- 420 • $n = 1$: a viscous-capillary flow, for an incompressible fluid and a low Reynolds
421 number.
- 422 • $n = 2$: a dissolution-precipitation process, driven by the Kelvin equation that states that
423 a curved surface has a higher saturation concentration and therefore can precipitate
424 ions dissolved in its vicinity.
- 425 • $n = 3$: a bulk diffusive lattice flow, can be caused by interstitial or substitutional
426 mechanisms in the bulk of the salt crystal.
- 427 • $n = 4$: surface diffusion, a rearrangement of surface molecules driven by the gradient
428 of the curvature of the profile.

429
430 Since Mullins also showed that the sum of two solutions is again a solution, we decomposed
431 the initial experimental profile at $t = 0$ as a Fourier sum and propagate each component
432 individually in time (where we used a cut-off after the 90 lowest-spatial-frequency
433 components). We found that the experimental data showed the best agreement with the first
434 order model describing the viscous flow process. This was obtained by a statistical test using
435 chi-square minimization between the u -values of the data and the four candidate mechanisms.
436 Note that in the simplification of the model the small slope approximation has been applied
437 ($du/dx \ll 1$) for both the projection of the molecular velocity and the curvature of the hole.

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443 **Methods references**

- 444 12. Steiger, M. & Asmussen, S. Crystallization of sodium sulfate phases in porous
445 materials: The phase diagram Na₂SO₄-H₂O and the generation of stress. *Geochimica*
446 *et Cosmochimica Acta* **72**, 4291–4306 (2008).
447 17. Mullins, W. W. Theory of Thermal Grooving. *Journal of Applied Physics* **28**, 333–339
448 (1957).
449 18. Mullins, W. W. Flattening of a nearly plane solid surface due to capillarity. *Journal of*
450 *Applied Physics* **30**, 77–83 (1959).

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457

458 **Author contributions**

459 Conceptualization: RW, NS

460 Methodology: RW, MD, DB, EJP, FA, NS

461 Investigation: RW, MD, DB, EJP, FA, NS

462 Visualization: RW, MD, EJP, NS

463 Project administration: NS

464 Supervision: NS, DB, FA

465 Writing – original draft: RW, MD, DB, NS

466 Writing – review & editing: RW, MD, DB, EJP, FA, NS

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468 **Competing interest declaration**

469 Authors declare that they have no competing interests

470

471 **Additional information (containing supplementary information line, corresponding**
472 **author line)**

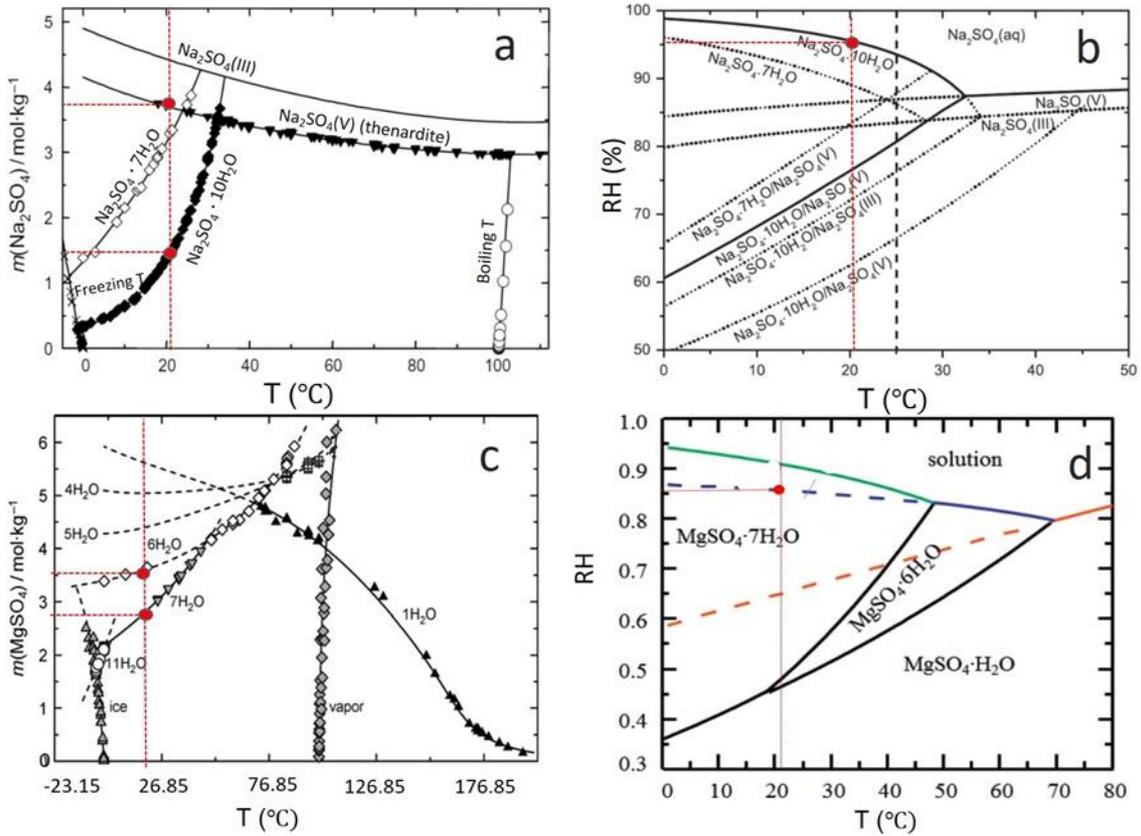
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474 **Supplementary Information** is available for this paper.

475 Correspondence and requests for materials should be addressed to Prof. dr. Noushine
476 Shahidzadeh.

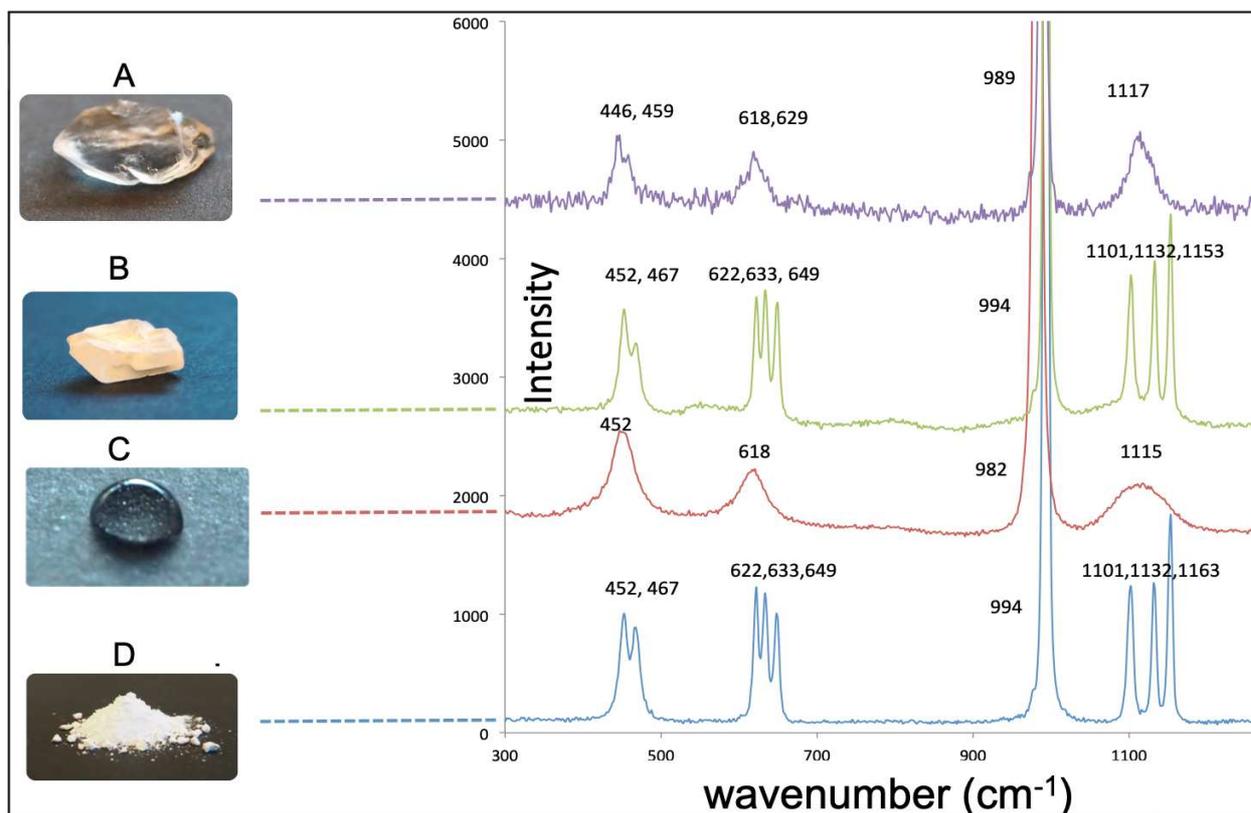
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Extended data figure/table legends



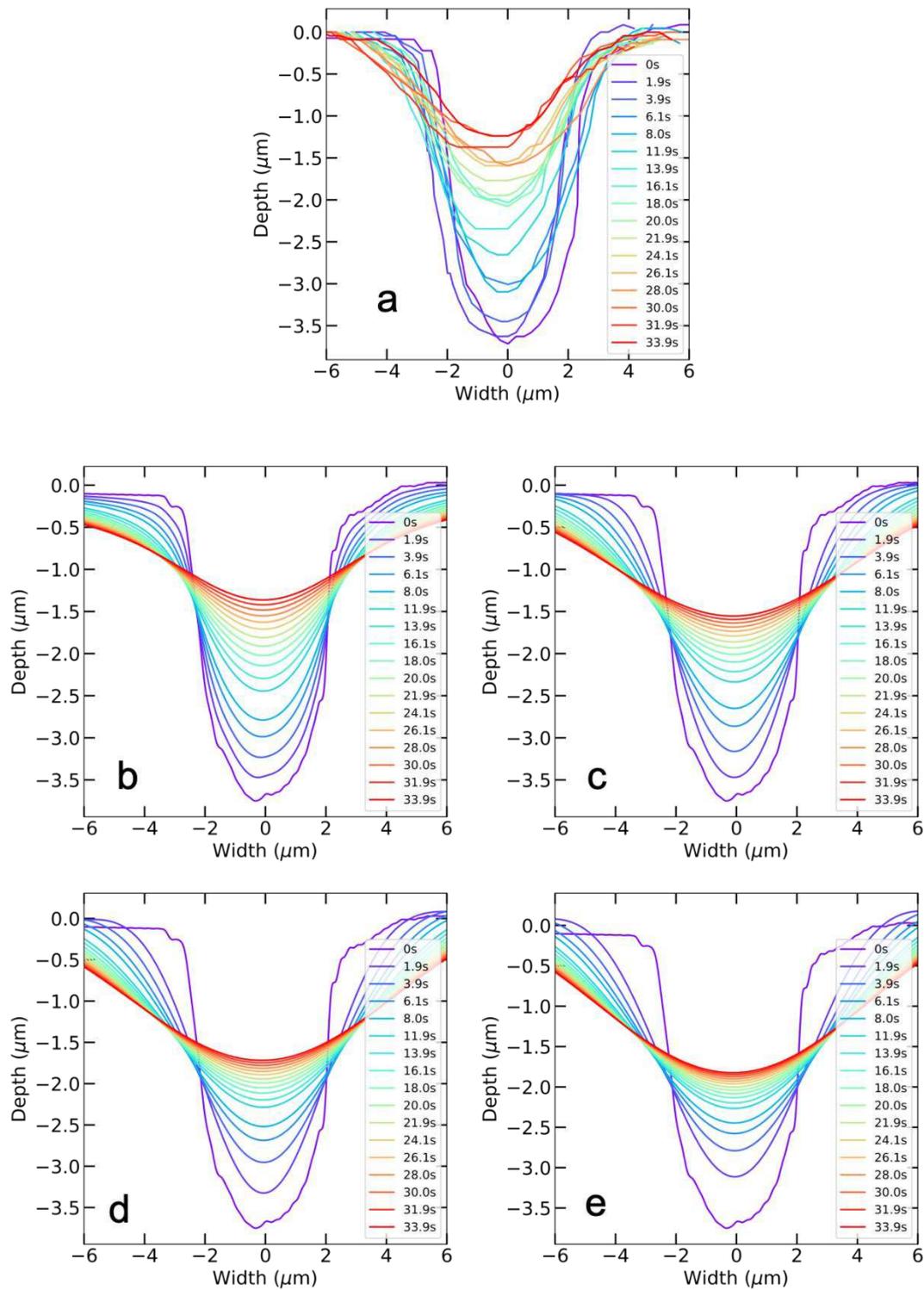
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Extended Data Figure 1 | Phase diagrams of sodium sulfate and magnesium sulfate from the literature. Solubility phase diagram (10) (a) and humidity phase diagram (13) (b) of sodium sulfate Mirabilite ($\text{Na}_2\text{SO}_4 \cdot 10\text{H}_2\text{O}$), hepta hydrate ($(\text{Na}_2\text{SO}_4 \cdot 7\text{H}_2\text{O})$, thenardite phase III and V (Na_2SO_4). Solubility phase diagram (14) (c) and humidity phase diagram (15) (d) of magnesium sulfate Epsomite ($\text{MgSO}_4 \cdot 7\text{H}_2\text{O}$), hexahydrate ($\text{MgSO}_4 \cdot 6\text{H}_2\text{O}$), kieserite ($\text{MgSO}_4 \cdot \text{H}_2\text{O}$). The red dotted lines represent our experimental conditions.



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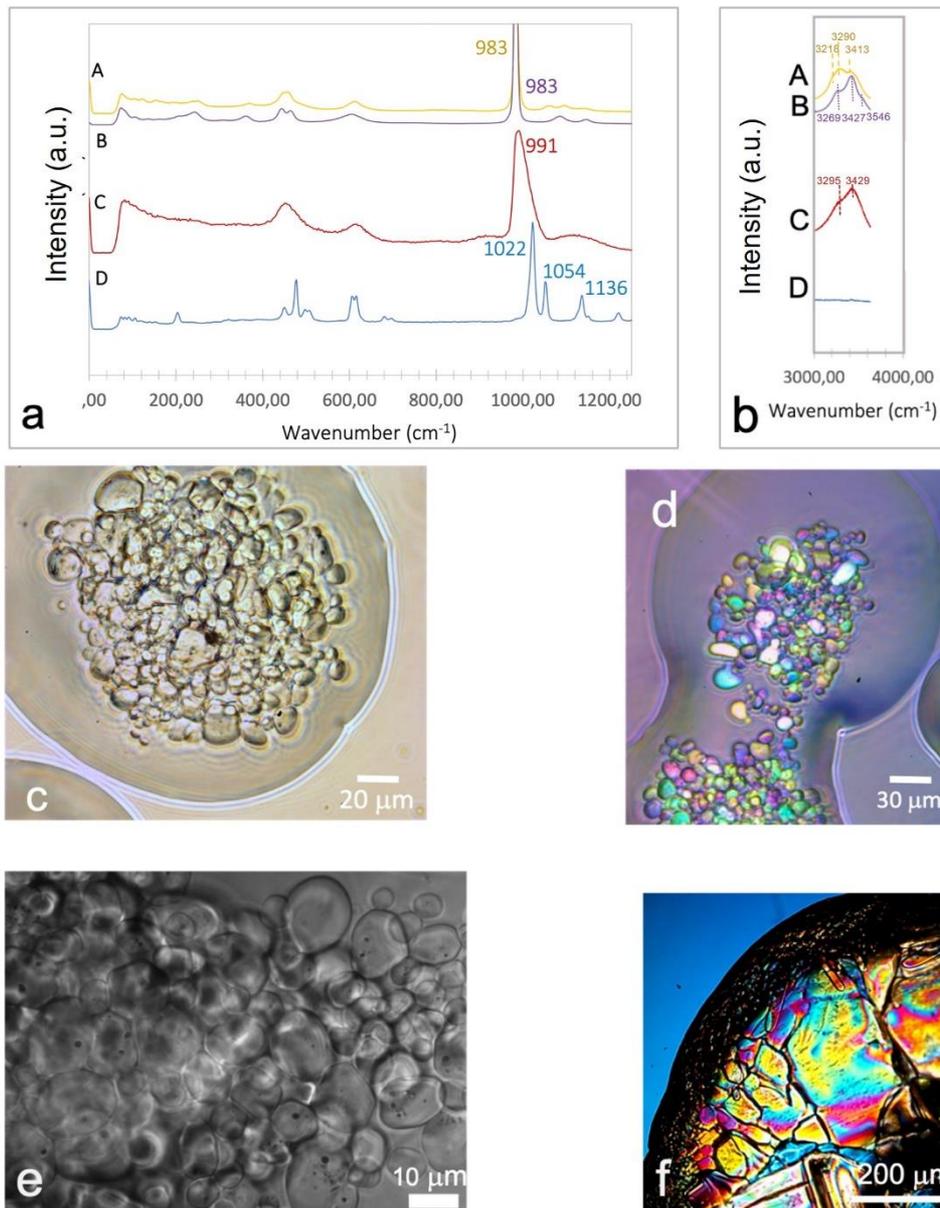
Extended Data Figure 2 | Raman spectra of different phases of sodium sulfate. (A) Mirabilite crystal (decahydrate of Na_2SO_4) (B) Dried Na_2SO_4 crystal, identified as thenardite (anhydrous Na_2SO_4). (C) Aqueous Na_2SO_4 solution. (D) Crushed dried Na_2SO_4 crystal, identified as thenardite (anhydrous Na_2SO_4). Spectra are vertically offset for clarity; intensities in arbitrary units.



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499

500 **Extended Data Figure 3 | Comparison of the four models to the mirabilite self-healing**
 501 **data.** (a) Raw data; (b) Effective flow model ($n = 1$); (c) Diffusion model ($n = 2$); (d) Bulk
 502 lattice flow model ($n = 3$); (e) Rearrangement of surface molecules by surface diffusion
 503 model ($n = 4$).

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508 **Extended Data Figure 4 | Raman and optical microscopy of magnesium sulfate. (a,b)**

509 Raman spectrum of the ‘floppy’ magnesium sulfate crystal (A) compared to that of the

510 hexahydrate crystal (B), a saturated solution of magnesium sulfate (C), and anhydrous

511 magnesium sulfate powder (D). (c-e) Floppy magnesium sulfate crystals under optical and

512 light polarizing microscope. (f) Hexahydrate of magnesium sulfate in a dried droplet under

513 the light polarizing microscope.

514

515 **Extended Data Table 1 | Levelling velocity of holes in mirabilite crystal walls for holes of**
516 **different depths**

Hole depth (μm)	Levelling velocity ($\mu\text{m/s}$)
4.2	0.205 ± 0.03
1.8	0.096 ± 0.03
4.2	0.535 ± 0.08
12.1	0.143 ± 0.01

517 This value is determined by the best fitting value $C_I(T)$ of the first order model.

518

519