

# Testing CPT symmetry in ortho-positronium decays with positronium annihilation tomography

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## Article

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# Testing CPT symmetry in ortho-positronium decays with positronium annihilation tomography

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## Abstract

Charged lepton system symmetry under combined charge, parity, and time-reversal transformation (CPT) remains scarcely tested. Despite stringent quantum-electrodynamic limits, discrepancies in predictions for the electron–positron bound state (positronium atom) motivate further investigation, including fundamental symmetry tests. While CPT noninvariance effects could be manifested in non-vanishing angular correlations between final-state photons and spin of annihilating positronium, measurements were previously limited by knowledge of the latter. Here, we demonstrate tomographic reconstruction techniques applied to three-photon annihilations of ortho-positronium atoms to estimate their spin polarisation without magnetic field or polarised positronium source. We use a plastic-scintillator-based positron-emission-tomography scanner as a high-acceptance photon detector to study CPT-prohibited angular correlation in ortho-positronium (o-Ps) annihilations. We record an unprecedented range of kinematical configurations

of o-Ps annihilations into three photons. Tomographic reconstruction of the annihilation points in a large medium allows single-event estimation of positronium spin orientation and, consequently, determination of the complete spectrum of angular correlation between the annihilation plane orientation and positronium spin, the non-vanishing expectation value of which would manifest CPT-violating effects. We find no violation at the precision level of  $10^{-4}$ , with a fourfold improvement on the previous measurement. This work enables application of positronium-imaging techniques to study discrete symmetries in positronium decays.

The symmetry under the combined operations of charge conjugation (C), spatial inversion (P), and reversal in time (T), referred to as CPT invariance, is the last of the fundamental discrete symmetries, a violation of which has never been observed. While numerous experimental tests have been carried out using baryonic matter systems (such as, neutral mesons that provide stringent limits [1–4]), and significant efforts have recently been made towards comparing the properties of hydrogen and anti-hydrogen atoms [5–7], certain CPT symmetry tests are also conceivable with purely leptonic systems.

Studies of neutrino oscillations in long-baseline experiments are prominent candidates in this sector. Discrepancies in mass and mixing parameters between neutrinos and antineutrinos would indicate CPT violation [8, 9], and while it has been argued that the sensitivity of such experiments could surpass even that of the neutral kaon system [8], recent results agree with CPT conservation [10].

Systems consisting of charged leptons, in contrast, are governed by quantum electrodynamics (QED), supported by experimental evidence of unparalleled precision. The most recent measurements of the electron anomalous magnetic moment [11], and its QED predictions based on the recent precise measurement of the fine structure constant [12] agree to 1 part in  $10^{12}$ , indicating that no new CPT-violating interactions coupling to the electron could be observed to this level of precision [13].

Nonetheless, the recently discovered deviations from QED predictions in the fine structure of the positronium atom, that is, the bound state of an electron and positron [14], suggest that diverse experimental studies of charged lepton systems should not be hastily abandoned. In fact, as the lightest matter–antimatter bound states, positronium atoms are considered in the context of CPT tests through searches for Lorentz noninvariance. Such tests, based on atomic spectroscopy of positronium have been recently proposed in the theoretical framework of Standard-Model Extension [15, 16].

A complementary, model-independent approach to searching for new CPT-violating interactions in the positronium system is constituted by searching for CPT-prohibited angular correlations in the annihilation of ortho-positronium atoms into three photons (o-Ps $\rightarrow$ 3 $\gamma$ ). This approach, first reported in 1988 [17], has seen relatively little activity in recent years, with the best result to date achieving a sensitivity of  $3 \times 10^{-3}$  [18].

In this work, we present a new experimental approach to search for a CPT-violating angular correlation in the three-photon annihilation of ortho-positronium atoms. We employ a positron-emission-tomography (PET) imaging device and demonstrate for the first time that three-photon annihilations of positronium can be reconstructed in a large volume, which allows for a new method of estimating its spin polarisation without the use of an external magnetic field. Using this technique, we measure the expectation value of

a CPT-prohibited angular correlation  $\hat{S} \cdot (\vec{k}_1 \times \vec{k}_2)$ , improving the sensitivity of the previous result [18] by a factor of four. Notably, the novel o-Ps $\rightarrow$ 3 $\gamma$  reconstruction technique presented here additionally opens the possibility of imaging of excited long-lived positronium states in studies of their gravitational fall as a matter–antimatter system [19].

#### CPT-PROHIBITED ANGULAR CORRELATIONS IN ORTHO-POSITRONIUM ANNIHILATIONS

The decaying ortho-positronium state is described by its unit spin along a given axis  $\hat{S}$ , whereas the final state of this annihilation is characterised by the momenta of the three photons  $\vec{k}_i$ ,  $i = 1, 2, 3$ , where we assume an ordering of  $|\vec{k}_1| > |\vec{k}_2| > |\vec{k}_3|$ . Certain angular correlation operators constructed using these observables are sensitive to the effects of violation of T, CP, or combined CPT symmetry if the corresponding operators are odd under the given symmetry transformation [20, 21]. Notably, these symmetry-violating correlations are not limited to positronium, and a similar approach has been employed to search for T noninvariance in the decays of  $Z^0$  bosons into hadronic jets [22].

Here, we focus on a study of the following operator:

$$\mathbf{O}_{\text{CPT}} = \hat{S} \cdot (\vec{k}_1 \times \vec{k}_2) / |\vec{k}_1 \times \vec{k}_2| = \cos\theta, \quad (1)$$

expressing the angular correlation  $\theta$  between the ortho-positronium spin and the orientation of the decay plane spanned by the three annihilation photons. This correlation operator constitutes a robust CPT-violation-sensitive observable defined through the fundamental geometry of an annihilation event, independent of particular kinematical configurations. The definition of the annihilation plane orientation through the two most energetic photons is merely a convenient convention and thus does not impose a bias on the space of kinematical configurations of three-photon annihilations, a problem affecting other conceivable angular correlations for o-Ps $\rightarrow$ 3 $\gamma$  events. Consequently, the choice of this observable for the CPT test minimises spurious asymmetries originating from the detector geometry that could mimic a violation [23]. Moreover, due to the neutrality of the final state of the o-Ps $\rightarrow$ 3 $\gamma$  process, radiative corrections that could cause fake CPT asymmetry, thus posing a natural limit on the sensitivity of such a test, are only expected at a precision level of  $10^{-9}$  [17, 20], thus leaving significant exploratory potential.

#### PREVIOUS MEASUREMENTS

Measurement of angular correlations in o-Ps $\rightarrow$ 3 $\gamma$  decays requires knowledge of the positronium spin as well as the recording of the momenta of the annihilation photons. Previous studies of CP- and CPT-violating operators achieved the former by either utilising a polarised positron beam [17] or polarising an o-Ps sample in an external magnetic field. The latter approach has been commonly used in studies of the CP-violation-sensitive correlation  $(\hat{S} \cdot \vec{k}_1)(\hat{S} \cdot \vec{k}_1 \times \vec{k}_2)$ , which has a sensitivity to symmetry violation that demands specific tensor polarisation of ortho-positronium [24, 25]. Notably, these experiments simultaneously recorded photons in only a single annihilation plane so that the magnetic-field-generation setup did not overlap with the detection devices. In turn, the measurements were limited to a scalar asymmetry

between two opposite configurations integrated over possible angles. Moreover, the need to reconfigure the detection setup to detect opposite decay geometries accounted for a significant source of systematic uncertainty in these measurements [24, 25].

The most recent search for CPT violation using the  $\cos\theta$  correlation, therefore, used the intrinsic polarisation of positrons from  $\beta^+$  decay combined with limiting the direction of their emission, which allowed for the use of a high-angular-acceptance photon detector capable of recording multiple scattering planes simultaneously. Thus, the first search for asymmetry in the distribution of the  $\cos\theta$  operator could be performed [18]. The statistical polarisation of o-Ps was obtained by allowing positrons from a  $^{22}\text{Ne}$  or  $^{68}\text{Ge}$  source to thermalise and form positronium only in a hemisphere of aerogel with the source at its centre, elucidating the positron emission direction with  $\Delta\vartheta = 90^\circ$ . The effective polarisation of the positrons was therefore given by  $\frac{v}{c}\frac{1}{2}(1 + \cos(\Delta\vartheta))$ , where  $\frac{v}{c}$  is the intrinsic polarisation of positrons from  $\beta^+$  decay along the direction of emission, and the remainder is a geometrical factor limiting the directional  $e^+$  polarisation, in this case to a cone with a  $2\Delta\vartheta$  opening angle [26]. Owing to the hemisphere used as the positronium production medium and the dimensions of the  $e^+$  trigger detector used, this geometrical factor amounted to 0.686. Moreover, despite the spherical symmetry of the detector, the chosen axis of the hemisphere required the source setup to be rotated in subsequent phases of the measurement to reduce systematic effects [18].

#### POSITRONIUM SPIN ESTIMATION AND RECONSTRUCTION OF O-PS $\rightarrow$ 3 $\gamma$ EVENTS IN J-PET

Our study was performed using the Jagiellonian Positron Emission Tomograph (J-PET) detector, which was conceived as the first PET imaging device based on plastic scintillators [27, 28]. In addition to exploring the path towards a cost-effective total-body PET scanner [28–30] and imaging with properties of  $e^+e^-$  bound states produced during scans of living organisms [31–33] as well as positron annihilation lifetime spectroscopy [34], J-PET constitutes a robust photon detector well suited to the study of angular correlations in o-Ps $\rightarrow$ 3 $\gamma$  annihilations [35, 36].

Figure 1 presents the methodology used in this work for studying o-Ps $\rightarrow$ 3 $\gamma$  events with positronium spin control without the use of an external magnetic field. Annihilation photons are detected through their Compton interaction in strips of plastic scintillators with an interaction time resolution at the level of 100 ps [27] and an angular resolution of approximately  $1^\circ$  (Fig. 1 (a)).

Data acquisition operates in a triggerless mode [37, 38] to minimise the initial sample bias for discrete symmetry tests and searches for rare positronium decays [39]. A detailed description of the J-PET detector can be found in the supplementary material A.

To estimate the ortho-positronium spin by utilising the imaging capabilities of the J-PET tomograph, we devised an extension of the spin control method used in Ref. [18] based on allowing positrons from a point-like  $\beta^+$  source to form o-Ps in a large volume following the symmetry of our photon detector. We used a  $^{22}\text{Na}$  source placed in the centre of a cylindrical vacuum chamber, the walls of which were coated with a layer of porous medium, enhancing positronium formation (Fig. 1(b)). This layer, consisting of

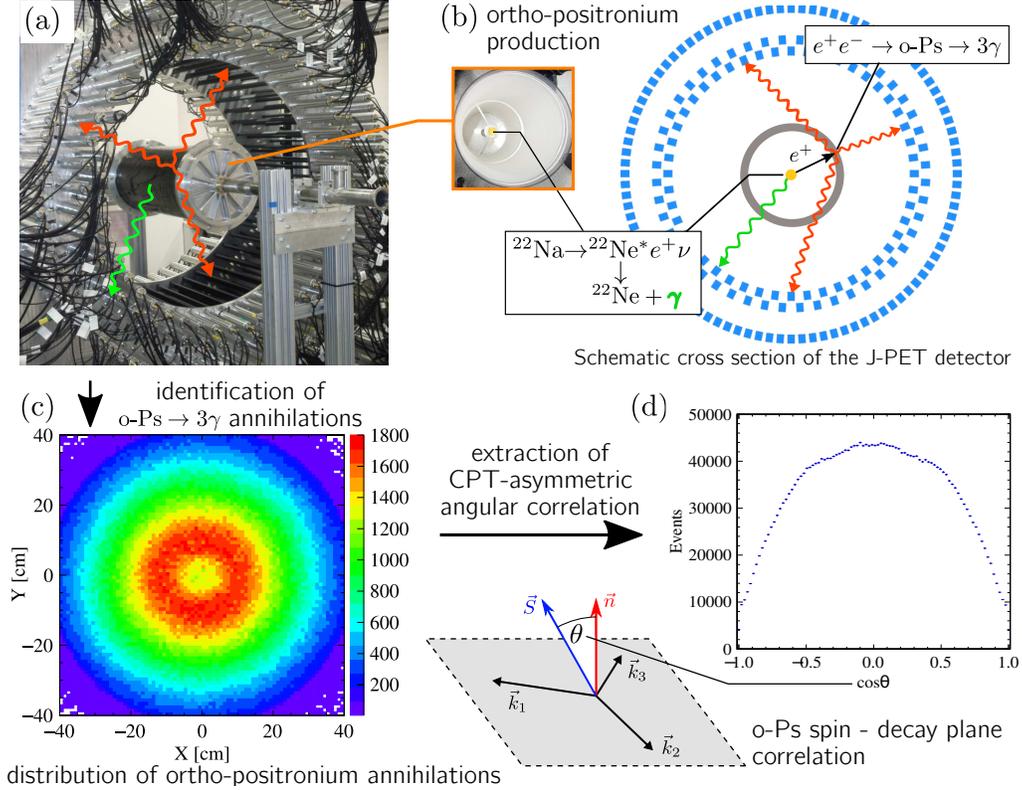


Fig. 1: Schematic of the research methodology leading to measurement of the CPT-violation-sensitive angular correlation between ortho-positronium spin and the plane of its three-photon annihilation. (a) Photograph of the J-PET tomograph with the cylindrical positronium production chamber. Plastic scintillators (black strips) record photons produced in  $o\text{-Ps} \rightarrow 3\gamma$  events (red arrows) and a prompt photon from  $\beta^+$  emitter de-excitation (green arrow). (b) Photograph of the interior of the positronium production chamber and a schematic view of the transverse cross section of the J-PET tomograph. Plastic scintillators (with rectangular cross sections (blue)) form three coaxial rings. Annihilating  $o\text{-Ps}$  atoms are produced by positrons from a  $^{22}\text{Na}$   $\beta^+$  source in the centre of a cylindrical vacuum chamber (grey), the walls of which are coated with porous silica, where positrons thermalise and form positronia. (c) Reconstructed spatial locations of the identified  $o\text{-Ps} \rightarrow 3\gamma$  events obtained for the first time using three-photon annihilations, allowing us to reproduce a tomographic image of the vacuum chamber. The maximum density distribution is in a ring with a radius of 12 cm, equal to the radius of the positronium production chamber. (d) For every selected event, the  $o\text{-Ps}$  spin axis is reconstructed, allowing calculation of the cosine of the angle between the spin ( $\vec{S}$ ) and annihilation plane orientation ( $\vec{n} = \vec{k}_1 \times \vec{k}_2$ ), the distribution of which is sensitive to CPT-violating asymmetries. The histogram presents the determined  $\cos\theta$  distribution for  $3.0 \times 10^6$  identified  $o\text{-Ps} \rightarrow 3\gamma$  events.

mesoporous silica and gypsum (similar to that used in thin-layer chromatography), allowed for preparation of a highly porous self-supporting structure of extensive size. The usage of gypsum enhances binding of

the silica gel and its expansion during setting allows for better packing of the material on a concave surface, such as the inner wall of the cylindrical chamber. The porous medium was prepared by multi-stage deposition through the evaporation of a silica–gypsum suspension in an alcohol solution distributed on the inner surface of the cylinder.

The thickness of the obtained silica coating was approximately 3 mm, well above the maximum implantation depth expected for unmoderated positrons from  $^{22}\text{Na}$  [40]. The vacuum chamber walls consisted of 4 mm thick polycarbonate, allowing transmission of photons from positronium annihilation at the level of 90 % [41]. A pressure of  $10^{-3}$  Pa was maintained inside the chamber to reduce the scattering of positrons travelling from the source to the positronium production medium on the chamber walls [42].

To maintain the symmetry of the setup, we did not limit the  $e^+$  emission direction to any solid angle besides the limits of the chamber. Instead, the positron flight direction was estimated separately for every recorded three-photon event using the position of the  $3\gamma$  annihilation vertex reconstructed with a novel trilateration-based technique [39, 43].

Initial identification of  $o\text{-Ps}\rightarrow 3\gamma$  event candidates in the J-PET data was based on photon energy deposition measured through time-over-thresholds (TOTs) of pairs of photomultiplier signals recorded at two ends of a scintillator strip for every registered photon interaction [44, 45]. As Compton interaction in the J-PET plastic scintillators yielded a continuous energy deposition spectrum, candidates for the identification of photons from an  $o\text{-Ps}\rightarrow 3\gamma$  annihilation were recognised using a TOT window located below the Compton edge of 511 keV photons, as indicated in Supplementary Fig. 2. The fast response of the plastic scintillators allowed a fine coincidence time window of 2.5 ns to be employed for the identification of  $3\gamma$  event candidates.

The high angular resolution of the detector at the level of  $1^\circ$ , coupled with exclusive registration of  $o\text{-Ps}\rightarrow 3\gamma$  decays, compensates for the lack of direct measurement of the photon energy. Their momenta are completely reconstructed using the recorded points of Compton interactions in the detector. Subsequently, the directions of  $\vec{k}_1$ ,  $\vec{k}_2$ , and  $\vec{k}_3$  span between the  $\gamma$  recording points and the annihilation point reconstructed using the latter combined with the corresponding interaction times. Consequently, the knowledge of the relative angles between the photon momenta can be used to unambiguously determine their energies [46].

Further selection of  $o\text{-Ps}\rightarrow 3\gamma$  events is based on the reconstructed photon energies and relative angles between their momentum vectors. Moreover, the background originating from secondary Compton-scattered photons recorded in the detector is discriminated by testing a hypothesis of a photon travelling directly between all pairs of detected  $\gamma$  interactions (for details, see Sec. C in the Supplementary Material).

The identified sample of three-photon annihilations of ortho-positronium atoms constitutes the first measurement of such events taking place in a medium of extensive size and featuring annihilation point reconstruction, which allows for determination of a three-photon tomographic image of the annihilation chamber using the spatial density of  $o\text{-Ps}\rightarrow 3\gamma$  events (Fig. 1(c)). In addition to the studies discussed in this work, this newly explored tomographic capability opens prospects for new medical imaging modalities with spatially resolved determination of positronium properties [29, 31–33].

SEARCH FOR CPT-VIOLATING ANGULAR CORRELATIONS IN  $o\text{-Ps} \rightarrow 3\gamma$

In August 2018, we collected a total of  $7.3 \times 10^6$  event candidates in a continuous 26-day measurement using the described setup with a 10 MBq  $^{22}\text{Na}$  positron source. Figure 1(c) presents the distribution of reconstructed  $o\text{-Ps} \rightarrow 3\gamma$  annihilation points in the transverse plane of the detector where the location of the porous silica allowing for positronium formation on the walls of the vacuum chamber is visible as the region with the highest density of reconstructed points. Notably, this distribution constitutes the first image of an extensive-size object obtained using  $o\text{-Ps} \rightarrow 3\gamma$  annihilations.

To ensure proper reconstruction of the ortho-positronium spin, the recorded photon interaction times were varied within their uncertainties so as to minimise the discrepancy between  $\rho = \sqrt{X^2 + Y^2}$  of the reconstructed annihilation point and the radius of the annihilation chamber ( $R = 12$  cm). Only events in which the total required variation, measured with a  $\chi^2$ -like variable (with two degrees of freedom), did not exceed  $\chi^2 = 0.5$  were retained for evaluation of the  $\mathbf{O}_{\text{CPT}}$  operator. This selection rejected poorly reconstructed events in which the geometrical uncertainty of the estimated spin axis could deteriorate the resulting statistical polarisation of the  $o\text{-Ps}$  sample, as well as spurious  $3\gamma$  annihilation points resulting from reconstruction applied to background events (a comprehensive discussion of background sources is enclosed in Sec. B of the Supplementary Material).

For approximately three million selected events, the estimated spin axis direction and photon momentum vectors were used to evaluate the  $\mathbf{O}_{\text{CPT}}$  correlation (equation 1), resulting in the distribution presented in Fig. 1(d). Extraction of the degree of potential symmetry violation from the angular correlation requires determination of the expectation value of the corresponding operator, corrected by the analysing power of the setup. The latter is dominated by the effective ortho-positronium polarisation resulting from the evaluation of the spin axis orientation for each event. The effective polarisation is determined by (i) average longitudinal polarisation of positrons from  $^{22}\text{Na}$  ( $P_{e^+} = \bar{\beta} \approx 0.67$ ), (ii) the amount of polarisation statistically transferred to the formed positronium ( $\frac{2}{3}$ ), (iii) polarisation loss during positron thermalisation amounting to approximately 8% [18, 35], and (iv) the geometrical uncertainty of the estimated direction of positron emission, which was reduced to approximately 2% in the presented setup [43]. The total effective polarisation was estimated to be  $P \approx 40.3\%$ .

It should be noted that while previous experiments [17, 18, 24, 25] concentrated on measuring the asymmetry of a given angular correlation, defined as  $A(|\mathbf{O}_{\text{CPT}}|) = (N(+|\mathbf{O}_{\text{CPT}}|) - N(-|\mathbf{O}_{\text{CPT}}|)) / (N(+|\mathbf{O}_{\text{CPT}}|) + N(-|\mathbf{O}_{\text{CPT}}|))$  (where  $N$  denotes the number of observed events for a given correlation value), it is only a special case of the expectation value of the corresponding operator in the case where a single value of  $\pm|\mathbf{O}_{\text{CPT}}|$  is measurable simultaneously.

Therefore, we use the expectation value  $\langle \mathbf{O}_{\text{CPT}} \rangle$  over the entire region of its definition as a measure of the observed asymmetry and obtain a value of

$$\langle \mathbf{O}_{\text{CPT}} \rangle = -0.00018 \pm 0.00031,$$

thus observing no significant asymmetry. The error is dominated by the statistical uncertainty of 0.00028.

For comparison with the measurements conducted to date, a parameter quantifying the level of observed CPT violation  $C_{\text{CPT}} \in [0, 1]$  can be used, as in Ref. [18], for which  $C_{\text{CPT}} = 1$  corresponds to a maximal asymmetry violating CPT. Correcting the above result for the analysing power ( $P$ ), we obtain

$$C_{\text{CPT}} = \langle \mathbf{O}_{\text{CPT}} \rangle / P = -0.00045 \pm 0.00077.$$

#### CONCLUSIONS AND PROSPECTS

We have demonstrated the application of tomographic reconstruction of three-photon positronium annihilations to determine the spin polarisation of ortho-positronium atoms, enabling precise studies of discrete symmetries in positronium decays without requiring a magnetic field. Through the exploratory measurement presented in this work, we have demonstrated that the J-PET detector based on plastic scintillator strips can exclusively register three-photon annihilations of ortho-positronium atoms occurring in an extensive-size medium with reconstruction of the spatial location of the annihilation and o-Ps spin-axis estimation on a single-event basis. With this 26-day measurement, the sensitivity to the CPT-violating angular correlation between ortho-positronium spin and orientation of the annihilation plane reached a statistical precision of  $10^{-4}$ , a factor of four better than the previous result [18].

Further sensitivity improvements to the presented CPT symmetry test are possible owing to an additional densely packed layer of plastic scintillators with a fully digital readout, which is currently being added to the J-PET detection setup. Providing up to quadruple enhancement in registration probability for a single annihilation photon, it can increase the overall detection efficiency of o-Ps  $\rightarrow 3\gamma$  events by a factor of 64. This will be further augmented by a new design of the annihilation chamber with spherical geometry, increasing the o-Ps formation probability by a factor of approximately 1.5 with a given  $\beta^+$  source activity and minimising potential spurious asymmetries from the geometry of the experimental setup [23]. These improvements will allow the J-PET setup to collect a dataset of exclusively recorded ortho-positronium three-photon annihilations of unprecedented size. Even with a moderate extension in measurement time compared to the experiment presented in this work, the statistics of the collected sample of o-Ps  $\rightarrow 3\gamma$  will allow for CPT symmetry testing at a precision level of  $10^{-5}$ . Moreover, the control over systematic effects in the presented measurement and coverage of a broad span of kinematic configurations of recorded events. This eliminates false asymmetries from the detector geometry, and opens the way to exploring angular correlation operators for which usage was not feasible in previous experiments, such as the  $\vec{S} \cdot \vec{k}_1$  correlation, which is odd under both CP and CPT transformations [23].

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#### AUTHOR CONTRIBUTIONS

The experiment was conducted using the J-PET apparatus developed by the J-PET collaboration. All authors participated in the construction and operation of the experimental setup or in the development of data reconstruction and analysis procedures. The J-PET detector and the techniques of the experiment were conceived by P.M. The data analysis was conducted by A.G. and M.M. The positronium production chamber was designed and constructed by M.G., J.G., B.J., and A.S. The manuscript was prepared by P.M and A.G and was then edited and approved by all authors.

#### COMPETING INTERESTS

The authors declare no competing interests.

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## APPENDIX

### A. J-PET detector

The (J-PET) detector comprises three sparse concentric layers of 50 cm long plastic scintillator strips arranged axially in a cylindrical geometry (see Fig. 1(a)). Each scintillator strip constitutes a detection module with two vacuum tube photomultipliers attached at its ends, which collect light produced as a result of Compton scattering of a  $\gamma$  quantum in the strip.

The longitudinal positions and time of  $\gamma$  interactions in the scintillator strips are reconstructed using times of light propagation from the point of  $\gamma$  Compton scattering to the ends of the strip, providing interaction time resolution at a level of approximately 100 ps owing to the fast signals of the plastic scintillators [27].

Synchronisation of time signals from both sides of each scintillator, as well as between detection modules and layers, is obtained from self-calibration of the recorded data using a statistical technique exploiting time differences between  $2\gamma$  annihilations and uncorrelated prompt photon emissions from the  $^{22}\text{Na}$  source [47].

### B. Identification of $o\text{-Ps} \rightarrow 3\gamma$ annihilation event candidates

For every recorded interaction of a  $\gamma$  quantum in the plastic scintillators, the two photomultiplier electric signals are sampled by the J-PET front-end electronics in the voltage domain at four predefined amplitude thresholds applied to their leading and trailing edges [48], as depicted in the left panel of Fig. 2. This technique provides a measure of the total size of the two signals through their TOTs, which is used to estimate the total energy deposited in the Compton scattering of the incident photon [45]. The right panel of Fig. 2 presents the spectrum of TOT values for all recorded photon interactions, in which the three visible structures correspond to an overlap of ortho-positronium annihilation and secondary Compton-scattered photons (hatched red), Compton spectrum edge for 511 keV photons from  $e^+e^-$  annihilations, and Compton edge for 1275 keV photons from de-excitation of the  $\beta^+$  source (dotted green), respectively, for increasing TOT values. Candidates for ortho-positronium annihilation photons are selected by TOT

values in a window of (15;55) ns to exclude noise and minimise the contribution of direct  $2\gamma$  annihilations that are abundant close to the middle Compton edge.

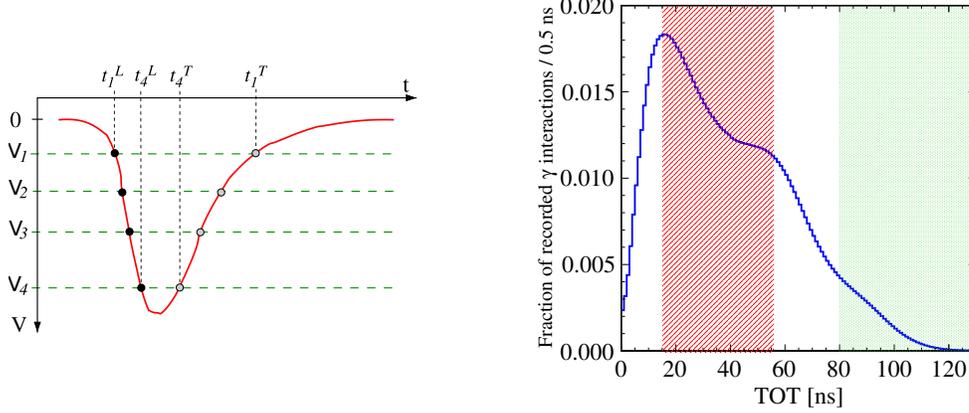


Fig. 2: Left: Photomultiplier electric signals in J-PET sampled at four predefined voltage thresholds, yielding four timestamps for the leading ( $t_i^L$ ) and trailing ( $t_i^T$ ) edges of the signal. The total TOT of a signal is calculated as  $\sum_{i=1}^4 (t_i^T - t_i^L)$ . Right: Distribution of total TOTs from both photomultiplier signals in J-PET detection modules, used as a measure of photon energy deposited in Compton scattering. The hatched red region is used to identify o-Ps annihilation photon candidates, whereas candidates for prompt photons from  $^{22}\text{Ne}^*$  de-excitation are found in the green dotted region.

Candidate events for o-Ps three-photon annihilations are identified as clusters of three-photon interactions recorded in a coincidence time window of 2.5 ns. At this stage, the background is mostly comprised of secondary Compton scatterings recorded in the detector, as well as  $2\gamma$  annihilations originating close to the  $\beta^+$  source recorded alongside a low-energy deposition of a de-excitation photon from the source itself. The event topologies of the signal and major background components are presented schematically in Fig. 3. The next section describes the treatment of particular background sources.

### C. Rejection of secondary Compton scatterings

The use of plastic scintillators in J-PET entails the possibility of recording secondary Compton scatterings mimicking the registration of primary photons (Figs. 3(b) and (c)). Back-to-back  $e^+e^- \rightarrow 2\gamma$  annihilations recorded with secondary scattering constitute a significant source of background for the identification of  $3\gamma$  events. However, the time resolution of plastic scintillators also allows such events to be distinguished by testing a hypothesis of a photon propagating between every two recorded interactions, quantified by

$$\delta_{ij} = |d_{ij} - c\Delta t_{ij}| = \|\vec{r}_i - \vec{r}_j\| - c|t_i - t_j|, \quad (2)$$

where  $d_{ij}$  is the distance between the  $i$ th and  $j$ th photon interactions recorded in an event candidate with recording times of  $t_i$  and  $t_j$ , respectively, and  $c$  denotes the velocity of light.

With such definition,  $\delta_{ij}$  assumes values close to zero for photon interaction pairs corresponding to a secondary Compton scattering. Consequently, its contamination among three-photon events is identified using  $\delta_{min} = \min\{\delta_{ij}\}$  for  $i \neq j=1,2,3$  calculated for every event. The contribution of events with secondary scattering, clearly visible in the right panel of Fig. 4, is rejected by a requirement of  $\delta_{ij} > 15$  cm.

The effect of the removal of secondary Compton scatterings is visible in the relative distribution of the sum and difference of the two smallest angles between the reconstructed momenta of the three photons, which is a convenient figure for evaluation of the o-Ps $\rightarrow$ 3 $\gamma$  sample purity [49]. As follows from a comparison of schemes (a) and (b-d) in Fig. 3, the genuine 3 $\gamma$  annihilations are characterised by large sums of the two smallest angles  $\theta_1 + \theta_2$  with a relatively small difference between them. Conversely, secondary scattering events are located closer to  $\theta_1 + \theta_2 \approx 180^\circ$  with a large  $\theta_2 - \theta_1$  value. Figure 5 presents the relative distribution of these figures before and after application of the criterion on  $\delta_{min}$ , demonstrating the purification of the event sample from secondary scattering events.

#### D. Rejection of two-photon annihilations coincident with a de-excitation photon

As indicated in Fig. 1 (b), the positron source used in this experiment emits prompt photons of 1275 keV owing to de-excitation of the  $\beta^+$  decay product. While such photons provide a useful start signal for evaluation of the positronium lifetime, their Compton interactions deposit energy from a continuous spectrum, which may mimic low-energy annihilation photons. As a result, the identified o-Ps $\rightarrow$ 3 $\gamma$  event candidates are contaminated by two-photon direct  $e^+e^-$  annihilations accompanied by a de-excitation  $\gamma$  quantum if it deposits low energy in a J-PET detection module (Fig. 3(d)). This background is recognised by the difference in azimuthal coordinates of the detection modules (labelled  $\theta^{XY}$ ) between two of the three interaction points being close to  $180^\circ$ . Moreover, if hypothetical lines of response are considered for every pair of recorded  $\gamma$  interactions and tomographic reconstruction of a 2 $\gamma$  annihilation point is performed, such events are characterised by one of the pairs yielding a hypothetical annihilation point at a small distance  $d_{LOR}$  from the  $\beta^+$  source.

Figure 6 presents the minimal values of  $d_{LOR}$  in each three-photon event as a function of the sum of the two smallest azimuthal locations of recorded  $\gamma$  interactions, where a structure corresponding to background events is removed by a cut marked by the black rectangle.

#### E. Evaluation of systematic uncertainties

False asymmetries can originate from the geometry of the positronium production setup or the detector, as well as from the contributions of background processes such as cosmic radiation recorded during measurement. Geometrical effects were evaluated by determining the following CPT-even and CP-odd angular correlation operator:

$$\mathbf{O}_{CP} = (\hat{S} \cdot \vec{k}_1)(\hat{S} \cdot \vec{k}_1 \times \vec{k}_2). \quad (3)$$

While it is potentially sensitive to CP-violation effects [20, 25], these could only be observed in the case of a tensor-polarised ortho-positronium sample available when an external magnetic field ( $\vec{B}$ ) was used.

As no  $\vec{B}$  fields were present in the J-PET measurements, all asymmetries manifested in the distribution of  $\mathcal{O}_{\text{CP}}$  must have originated from the detector geometry. The mean value of this operator was found to be  $(1.3 \pm 1.4_{\text{stat}}) \times 10^{-4}$ , showing no significant detector asymmetries at the level of  $10^{-4}$ .

The impact of cosmic rays recorded during the experiment was estimated in a dedicated measurement without the  $\beta^+$  source and positronium annihilation setup, where only cosmic radiation was recorded. It was estimated that cosmic ray events passing the event selection criteria may constitute at most  $3.9 \times 10^{-5}$  of the final event sample, which was considered as part of the systematic uncertainty.

The influences of subsequent event selection criteria applied in the analysis were tested by variation of each of the cut values within a few multiples of their corresponding experimental resolution. The only analysis cut found to have a statistically significant impact on the measured value of  $\langle \mathcal{O}_{\text{CPT}} \rangle$  was the upper bound on the TOT of candidates for annihilation photon interactions (see the right panel of Fig. 2). Consequently, an effect of  $1.26 \times 10^{-4}$  on  $\langle \mathcal{O}_{\text{CPT}} \rangle$ , corresponding to a variation of this bound by 1 ns, was included in the systematic uncertainty.

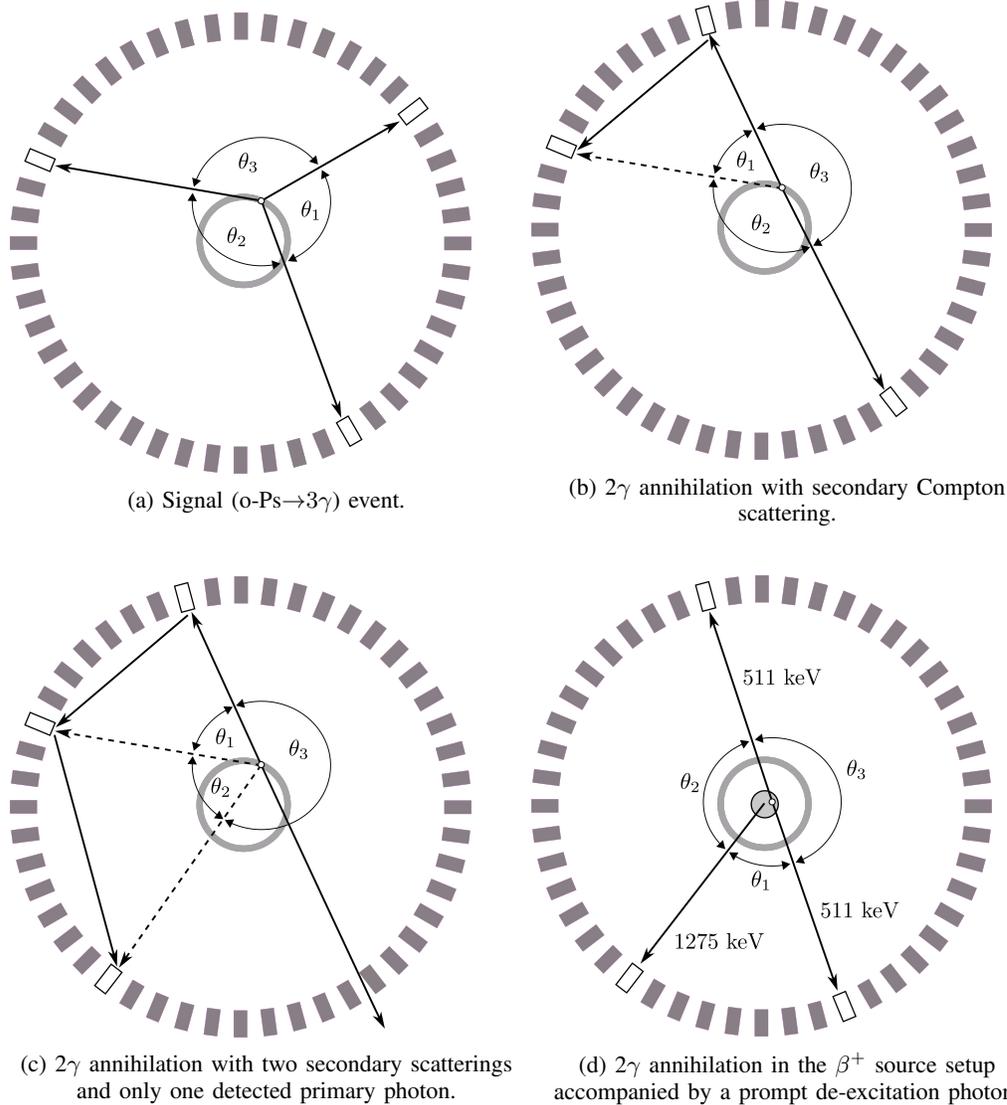


Fig. 3: Schematic presentation of signal (o-Ps $\rightarrow$  3 $\gamma$ ) and background events in the transverse view of the J-PET detector. The grey rectangles represent photon detection modules of plastic scintillators (only one layer is shown for readability). The modules in which photons were recorded are marked white. The annihilations take place in the wall of the vacuum chamber (grey band). In a signal event (a), all three photons from an ortho-positronium annihilation are recorded and the angles between their momenta, labelled according to ascending magnitude ( $\theta_1 < \theta_2 < \theta_3$ ), obey  $\theta_1 + \theta_2 > 180^\circ$ . The background is dominated by direct  $e^+e^- \rightarrow 2\gamma$  annihilations, where the third recorded photon comes from secondary Compton scattering in the detector (b), resulting in the identification of a spurious primary photon marked with a dashed arrow. Three-photon events may also be caused by two subsequent secondary scatterings in the case where one of the direct annihilation photons is not detected (c), and by two-photon annihilations originating in the  $\beta^+$  source setup (grey circle) accompanied by a prompt  $^{22}\text{Ne}^*$  de-excitation photon depositing low energy in Compton scattering (d).

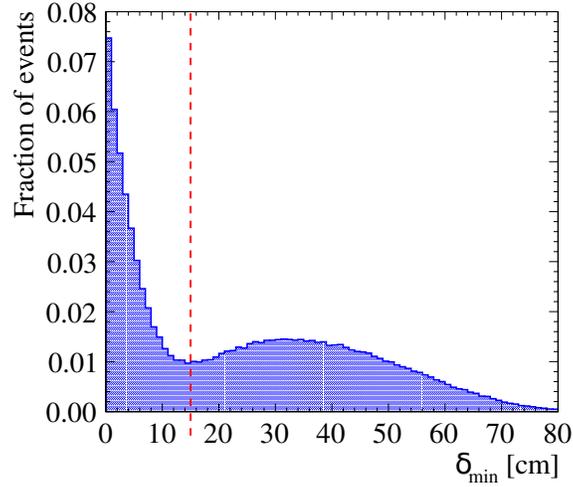


Fig. 4: Distribution of the minimal discrepancy between the travelled path and time of flight among all hypothetical secondary scattered photons considered in a single  $3\gamma$  event candidate. Events with  $\delta_{\min} < 15$  cm (marked with the red line) are discarded as containing secondary Compton scatterings in the detector modules.

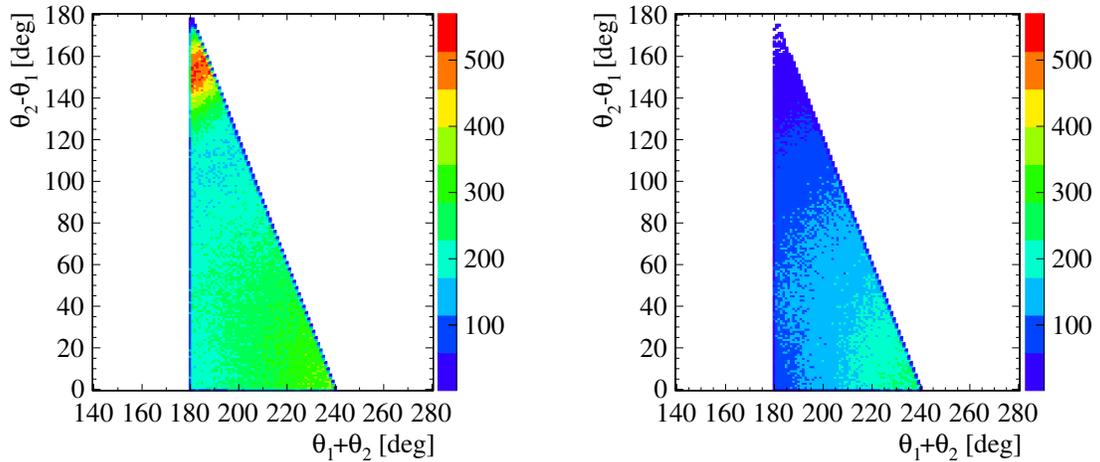


Fig. 5: Relative distribution of the sum and difference of the two smallest angles between photon momenta (see Fig. 3) in the identified  $3\gamma$  events before (left) and after (right) rejection of the secondary scattered photons. The  $\theta_1 + \theta_2 > 180^\circ$  limit results from momentum conservation in  $o\text{-Ps} \rightarrow 3\gamma$  annihilations.  $o\text{-Ps}$  annihilation events are expected in the right corner of the populated region [49], whereas the topmost region is specifically for secondary scattered photon events.

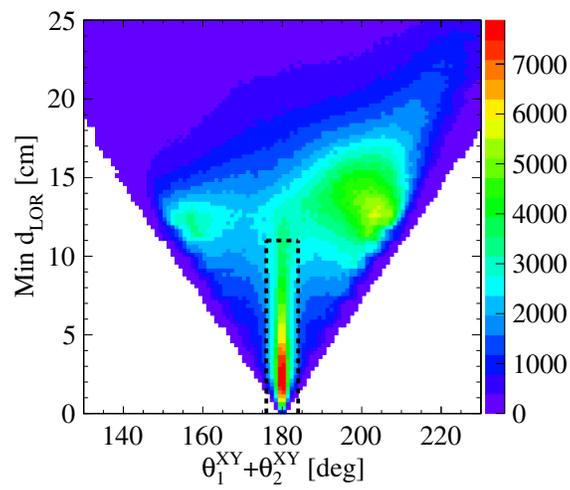


Fig. 6: Distribution of distance between the  $\beta^+$  source location and the closest hypothetical  $2\gamma$  annihilation point on a line of response between two recorded photon interactions vs. the sum of the two smallest angles between azimuthal coordinates of the recorded  $\gamma$  interaction points. Background events from  $2\gamma$  annihilations in the  $\beta^+$  source setup (see Fig. 3 (d)) are rejected by excluding the region marked with the black rectangle.

# Figures

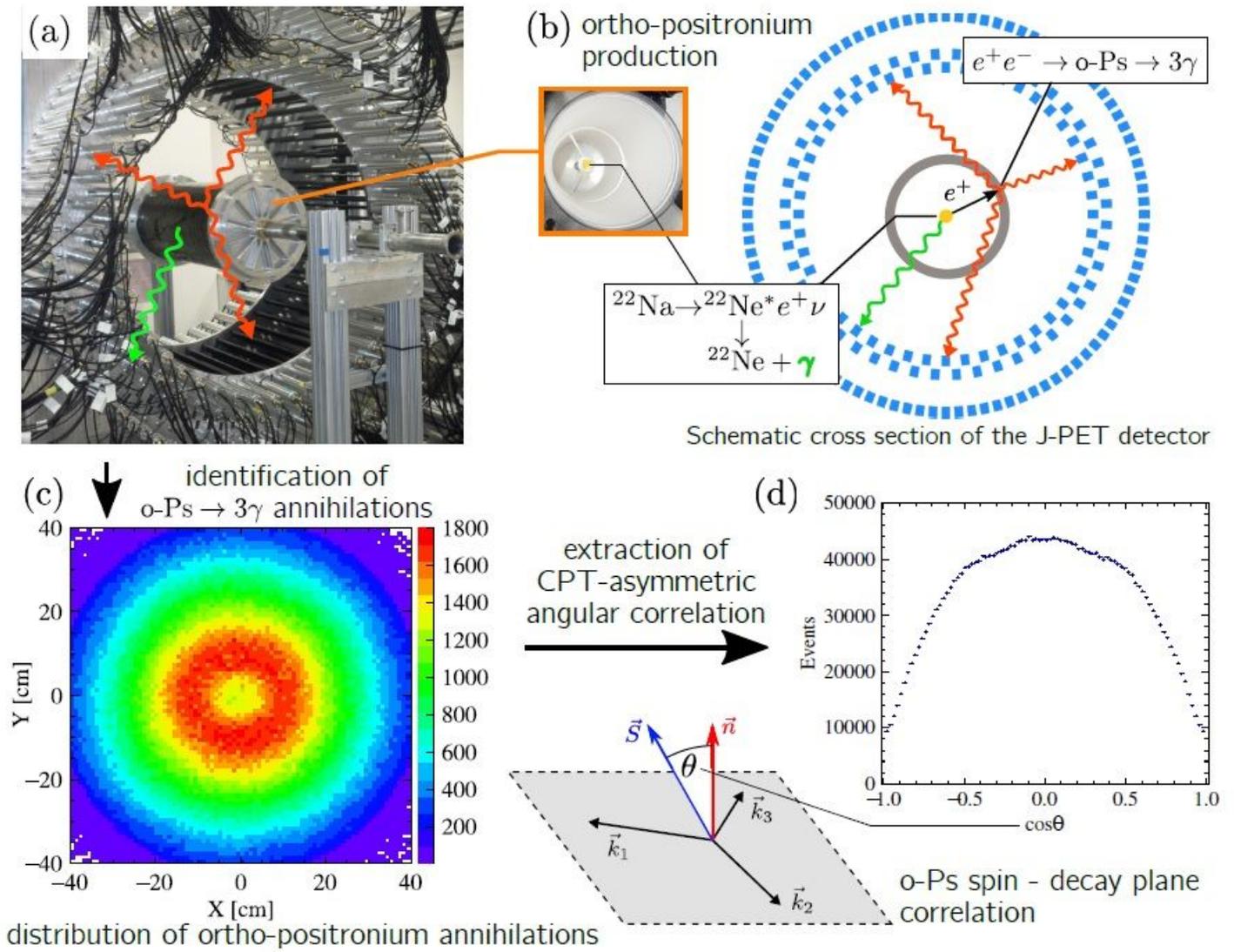
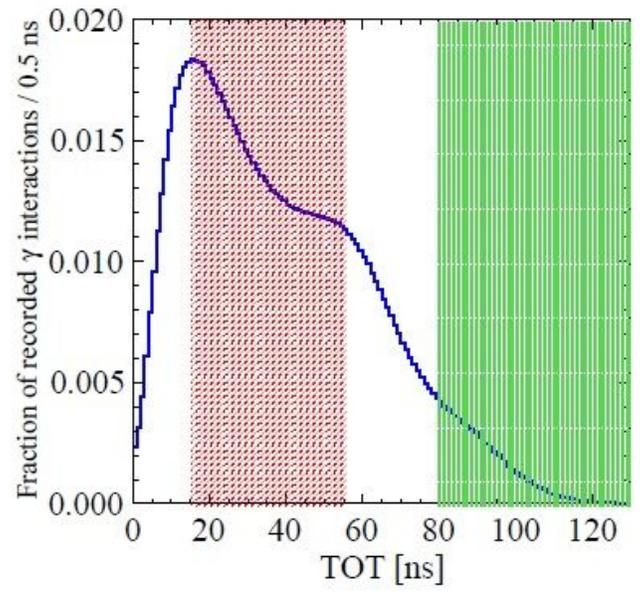
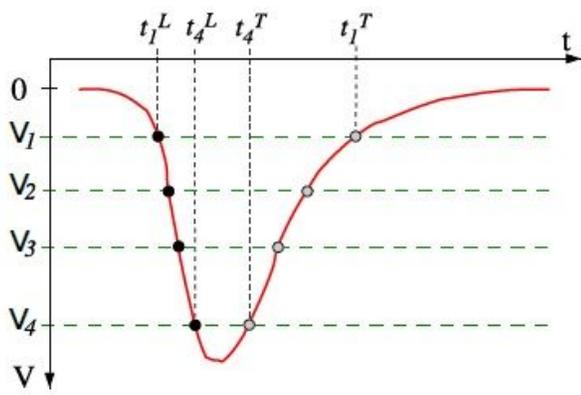


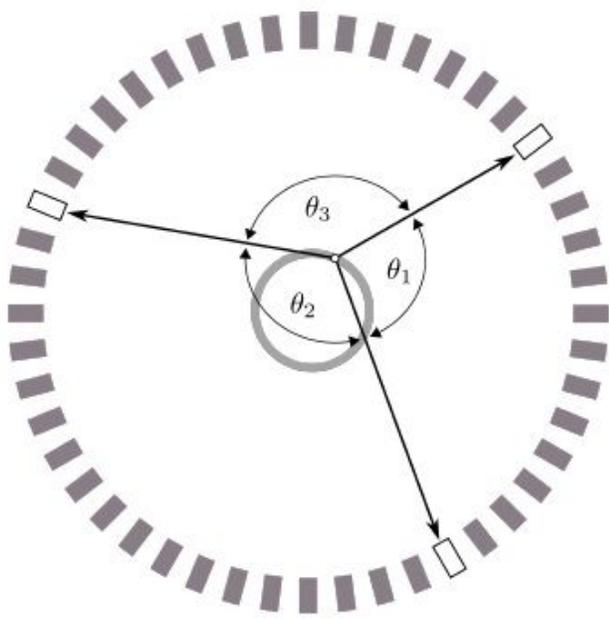
Figure 1

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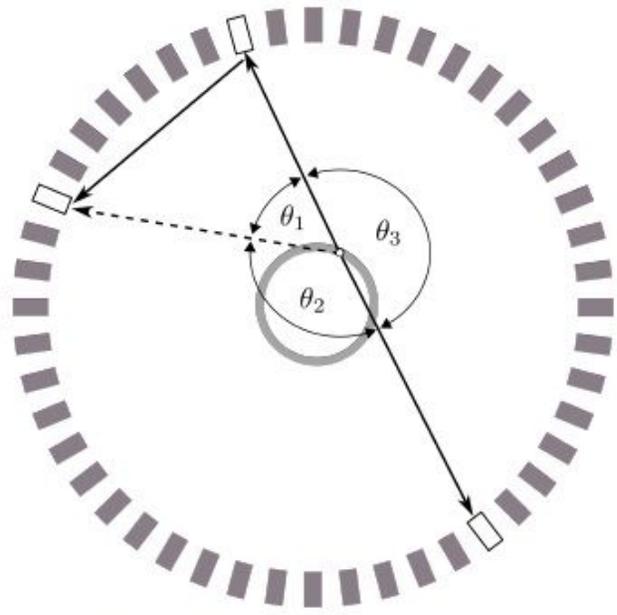


**Figure 2**

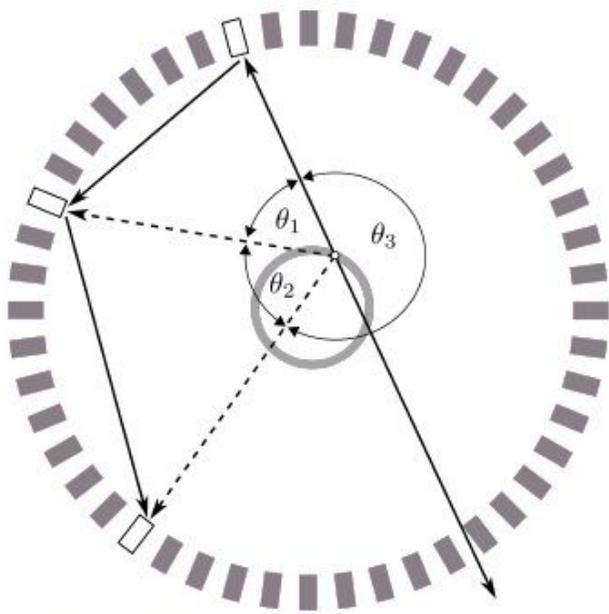
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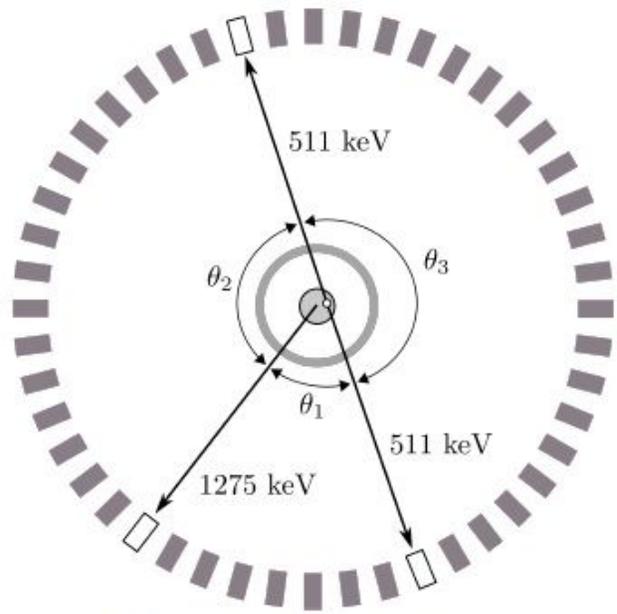
(a) Signal ( $o\text{-Ps} \rightarrow 3\gamma$ ) event.



(b)  $2\gamma$  annihilation with secondary Compton scattering.



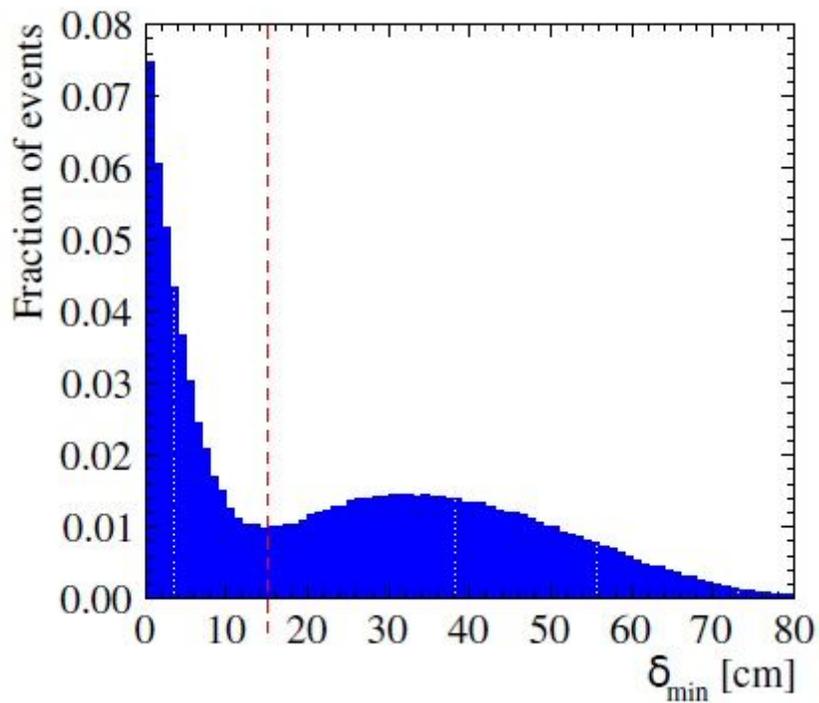
(c)  $2\gamma$  annihilation with two secondary scatterings and only one detected primary photon.



(d)  $2\gamma$  annihilation in the  $\beta^+$  source setup accompanied by a prompt de-excitation photon.

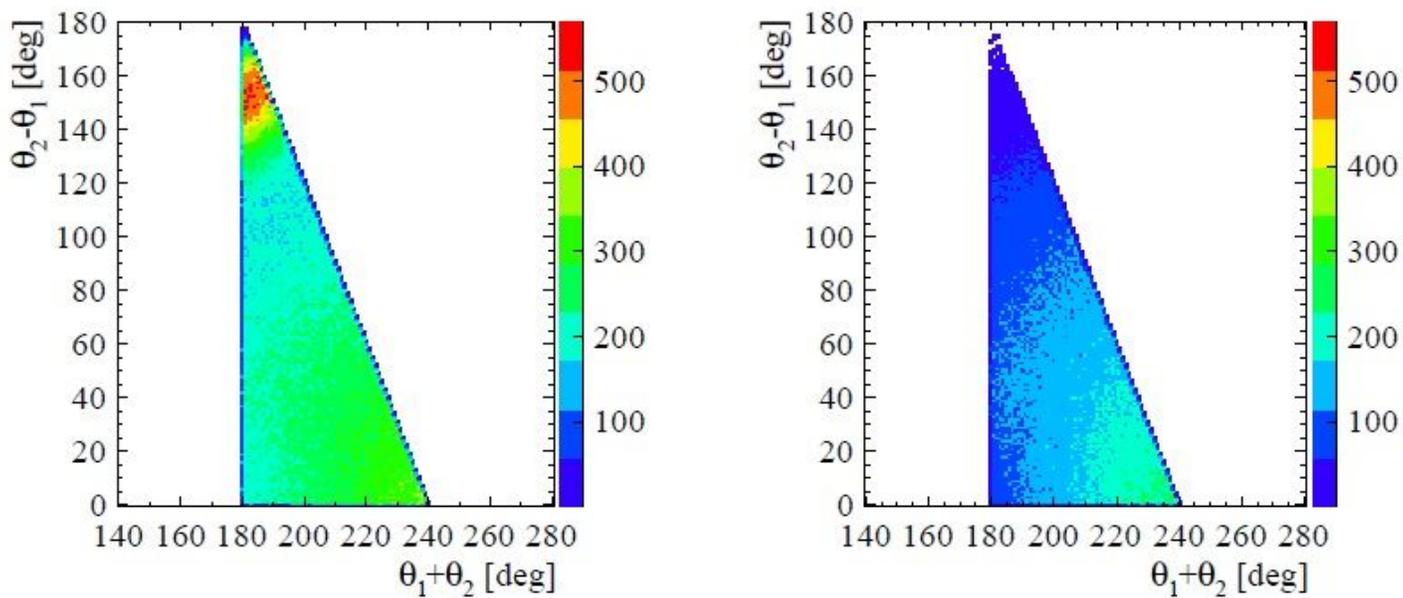
**Figure 3**

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**Figure 4**

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**Figure 5**

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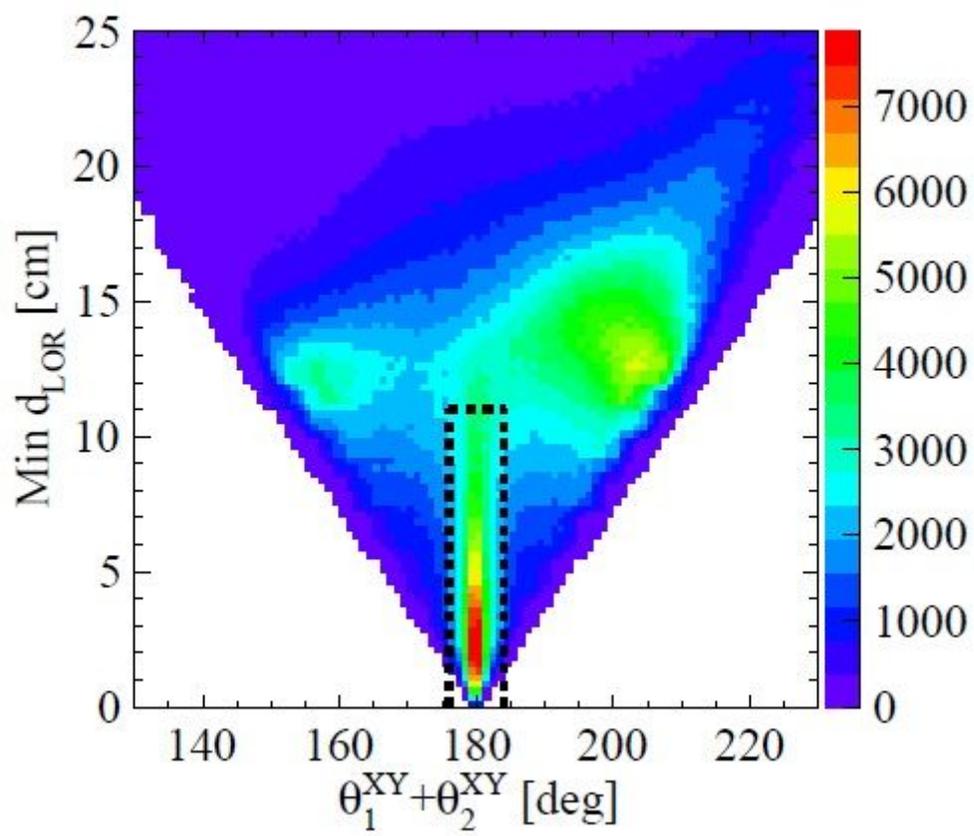


Figure 6

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