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Steps dominate gas evasion from a mountain headwater stream

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Abstract

Step-and-pools are dominant morphologic traits of high-energy streams, where climatically- and biogeochemically-relevant gases are processed, transported to downstream ecosystems or released into the atmosphere. Yet, capturing the signature of the small-scale morphological complexity of channel forms on large-scale river outgassing represents a fundamental unresolved challenge. Here, we combine theoretical and experimental approaches to assess the contribution of localized steps to the gas evasion from river networks. The framework was applied to a representative, 1 km-long mountain reach in Italy, where carbon dioxide concentration drops across several steps and a reference segment without steps were measured under different hydrologic conditions. Our results indicate that local steps lead the reach-scale outgassing, especially for high and low discharges. These findings suggest that step-pools are key missing components of existing scaling laws used for the assessment of gas fluxes across water-air interfaces. Therefore, global evasion from rivers may differ substantially from previously reported estimates.

Introduction

River networks transport, process and release a multitude of chemical substances, which are relevant to the biogeochemical functioning of stream ecosystems and eventually affect the fragile interconnections of the land-water-climate nexus¹⁻⁴. In particular, headwater streams are important greenhouse gas sources

23 to the atmosphere, owing to enhanced input of dissolved matter from the sur-
24 rounding landscape and high exchange rates across water-air interfaces. Conse-
25 quently, quantifying gas emissions from upland freshwater systems is regarded
26 as an important scientific challenge with multi-faced implications for a broad
27 range of disciplines including ecology, biology and climate sciences⁵⁻⁸.

28 Riffles represent a distinctive trait of most rivers worldwide⁹. In high-
29 gradient streams, where the granulometry is varied, riffles frequently give rise
30 to sequences of step and pools in which local hydraulic discontinuities of the
31 water flow are observed. Step-pool bedforms are thus widespread in different
32 regions of the globe, including humid areas, desert ephemeral systems, semiarid
33 environments and alpine settings^{10,11}. The important role of steps in regulating
34 fluvial sediment transport and channel morphology has been extensively stud-
35 ied¹²⁻¹⁹, but systematic knowledge about their contribution to gas exchange
36 between freshwater systems and the atmosphere remains elusive.

37 Several authors have argued that waterfalls, riffles, steps and cascades might
38 promote gas exchange with the atmosphere, owing to the enhanced turbulence
39 and air entrainment that are typically observed in correspondence of abrupt dis-
40 continuities of the flow field^{7,20-28}. However, available empirical data about the
41 outgassing produced by individual steps or cascades is relatively limited. Cirpka
42 et al.²⁰ and Natchimuthu et al.²⁹ have used tracer injections to demonstrate that
43 the presence of cascades and waterfalls significantly increases the reaeration co-
44 efficient in a set of tens-of-meters long river reaches located in Switzerland and
45 Sweden. More recently, Leibowitz et al.³⁰, Vautier et al.³¹ and Schneider et
46 al.³² have shown that mass evasion within channel stretches that contain wa-
47 terfalls is enhanced, causing a loss of carbon dioxide (CO_2) or injected tracers
48 within relatively short distances in the range of 15% to 50% of the initial mass.
49 However, spatial patterns of gas evasion were typically monitored at relatively
50 coarse spatial resolutions (i.e. some tens of meters), and empirical observations
51 at scales comparable to the step size are much rarer (see ref.³¹). High reso-
52 lution data, instead, represent a powerful means to reduce the uncertainty in
53 the characterization of small-scale patterns of stream outgassing, and enable
54 more robust assessments of gas evasion produced by local hydro-morphologic
55 heterogeneities of rivers³³.

56 Existing studies aimed at quantifying the relative contribution of cascades to
57 the total stream outgassing typically rely on spatial patterns of gas concentra-
58 tion within river reaches that contain steps or waterfalls³⁰⁻³². These patterns,
59 however, inherently mirror the unique morphologic characteristics of the case
60 studies selected for the analysis. Consequently, the existing estimates can not
61 be easily upscaled or extrapolated to different contexts. In other instances, gas
62 evasion produced by steps and cascades was investigated through the analysis
63 of the underlying mass transfer rate, k ^{27,29,31,34}. Though, this metric can be
64 strongly biased when the spatial resolution of the available data is insufficient³⁵.
65 In fact, the mean value of k in a given stream portion is not able to capture
66 internal peaks of gas transfer that might be observed when the size of the fluid
67 volume responsible for most of the evasion is smaller than the measurement
68 resolution - as in the presence of a falling jet (P.P., N.D. and G.B., in prep.).

69 Owing to these theoretical and practical limitations - in spite of the growing
70 awareness of the importance of local heterogeneity of the flow field in water-
71 air gas exchange - the relative contribution of steps to the total outgassing in
72 morphologically-complex reaches is not fully clear. This study aims at filling
73 this gap by developing a theoretical and experimental framework for the study
74 of gas emissions in heterogeneous streams. The approach was applied to a
75 high-gradient channel in the Italian Alps where water CO_2 concentrations were
76 measured across 19 natural and artificially-created steps, and along a reference
77 turbulent segment without steps. The main innovation of the experimental setup
78 is that we capitalize on direct CO_2 records gathered under different discharge
79 conditions, without relying on Schmidt-number scaling, which is problematic
80 in case of bubble-mediated transport of the type observed downstream of steps
81 and cascades^{25,34,36,37}. The major theoretical advance, instead, pertains the
82 development of a modular and scalable metric for disentangling the contribu-
83 tion of turbulent segments and local steps to the total stream outgassing. The
84 scalability of this metric was exploited to assess the role of local steps in the
85 reach-scale outgassing, taking into account both the temporal variations of the
86 discharge and the morphological characteristics of the river bed.

87 Concentration damping in streams

88 While we recognize that the morphological complexity of rivers might not fit
89 perceptual models of the type proposed here, we conceptualize high-gradient
90 stream networks as a heterogeneous sequence of two types of elements: steps
91 and segments (Figure 1b). Local steps are point-wise hydraulic discontinuities
92 of the flow, generated by a drop of the riverbed Δh higher than the typical
93 flow depth ($\Delta h > 10\text{ cm}$ in this case). In such circumstances, the presence of
94 a falling jet promotes air entrainment, bubbles and foaming which enhance gas
95 exchange with the atmosphere. Turbulent segments, instead, are continuous,
96 relatively regular river stretches located in between pairs of steps, in which the
97 flow is gradually varied. Therein, turbulence and gas-exchange are promoted by
98 heterogeneities of the velocity field, which are in turn produced by e.g. hurdles,
99 stones, bends and bed roughness.

100 A proper evaluation of the role of local steps in river gas evasion requires
101 the definition of a metric capable of objectively determining the separate con-
102 tribution to the outgassing of steps and turbulent segments. To this aim, we
103 utilize the concept of concentration damping, which is a dimensionless, scal-
104 able measure of gas evasion applicable to individual steps, single segments or
105 composite heterogeneous channels. Under steady state conditions, the $1D$ spa-
106 tial pattern of the concentration C of a dissolved gas advected downstream (x
107 direction) and evaded into the atmosphere is exponentially decreasing with a
108 spatially-heterogeneous decay rate (Methods):

$$C(x) = C_a + (C_0 - C_a) \exp \left[- \int_0^x \frac{K(x')}{u(x')} dx' \right] = C_a + (C_0 - C_a) \exp[-f(x)] \quad (1)$$

109 where $C(x)$ is the concentration in the position x along the streamline, C_0 the
 110 concentration in the upstream section ($x = 0$), $u(x)$ the water velocity, $K(x)$ the
 111 exchange rate (here defined as the gas transfer velocity k scaled to the water
 112 depth) and f is the exponential damping factor (hereafter damping factor).
 113 Starting from its definition as given in Equation (1), f can be expressed as the
 114 product between an effective exchange rate along the stretch $(0, x)$, $K_{eq}(x)$, and
 115 the corresponding water travel time $\tau(x)$, i.e. $f(x) = K_{eq}(x)\tau(x)$ (Methods).
 116 In particular, the damping factor of a continuous channel segment (say, c_i) of
 117 length ℓ_i can be written as:

$$f_{c_i}(\ell_i) = K_{c_i} \tau_{c_i}, \quad (2)$$

118 where $\tau_{c_i} = \int_0^{\ell_i} 1/u(x)dx$ is the travel time spent by water parcels in the segment
 119 (measurable via tracer experiments) and K_{c_i} the effective exchange rate therein.

120 Analogously, the damping factor of the step s_i , f_{s_i} , can be expressed as:

$$f_{s_i}(\Delta h_i) = K_{s_i} \tau_{s_i}, \quad (3)$$

121 where K_{s_i} is the effective exchange rate within the step i and τ_{s_i} the correspond-
 122 ing travel time. The notation emphasizes that f_{s_i} should depend on the step
 123 height Δh_i , which drives the amount of energy dissipated by the water flow and
 124 the ensuing outgassing process. In the case of steps and cascades, it is practi-
 125 cally unfeasible to separately measure τ_{s_i} and K_{s_i} . Nevertheless, step exchange
 126 rates (K_{s_i}) are expected to be very high owing to the local increase of energy
 127 dissipation and the enhanced air entrainment in correspondence of the falling
 128 jet^{20,38-40}. Consequently, f_{s_i} could be comparable to (or even larger than) f_{c_i}
 129 in many settings, in spite of the "local" nature of the outgassing process in
 130 correspondence of the steps (i.e., $\tau_{s_i} < \tau_{c_i}$).

131 The damping factor is additive and commutative (Methods), thereby im-
 132 plying that the value of f for a series of segments (steps) can be calculated as
 133 the sum of the damping factors associated to the individual segments (steps)
 134 involved, regardless of their specific order (SI). Consequently, we can evaluate
 135 the relative contribution to the total outgassing induced by all the steps em-
 136 bedded in a focus reach (or river network) using the dominance ratio r , which
 137 is defined as the ratio between the damping factor of the steps (f_s) and the
 138 damping factor of the segments (f_c) of that reach (see Figure 1b and SI):

$$r = \frac{\sum_i f_{s_i}}{\sum_i f_{c_i}} = \frac{f_s}{f_c}. \quad (4)$$

139 In Equation (4), f_c (f_s) is expressed as the sum of the damping factors of all
 140 the segments (steps) included in the reach (figure 1b and Methods). If the

141 dominance ratio is equal to unity, then the steps and the continuous segments
142 of the focus reach provide an equal contribution to the total river outgassing
143 (meaning that the gas concentration damping produced by a sequence of steps
144 would be the same as the concentration damping produced by the turbulent
145 continuous segments without the steps). Likewise, if $r > 1$ ($r < 1$) then the
146 steps proved a larger (smaller) contribution to the outgassing as compared to
147 the contribution generated by the continuous turbulent stretches. Crucially, r
148 is not affected by the specific order according to which segments and steps are
149 arranged, but only depends on key geometrical and hydraulic features of river
150 networks (e.g. step height/spacing, segment slope).

151 Steps dominate reach-scale gas evasion

152 The empirical estimates of the damping factors of the steps and the reference
153 segment belonging to the focus reach considered in this paper are summarized
154 in Fig. 2. In the reference segment, the damping factor f_{c_r} ranged from 0.1
155 (corresponding to approximately 10% of excess mass removed) to 0.32 (25%
156 of excess mass removed), depending on the underlying flowrate, Q . The non
157 monotonous dependence of f_{c_r} on Q was explained by the interplay between two
158 important drivers of gas transfer across water-air interfaces: i) the mean flow
159 velocity, which is positively correlated with the outgassing velocity and is an
160 increasing function of Q (Fig S4); and ii) the ratio between exchange area and
161 water volume, which is proportional to the exchange rate but decreased with Q
162 in the reference segment (Table S3). Consequently, f_{c_r} peaked for intermediate
163 discharges.

164 The estimated values of the damping factor in the natural and artificially-
165 generated steps (Methods and SI), instead, ranged from 0.03 to 0.27. In most
166 cases f_{s_i} had the same order of magnitude of f_{c_r} , demonstrating that a local
167 step may originate the same outgassing produced by turbulent river segments
168 with a length of tens of meters. Interestingly, f_{s_i} showed an almost linear
169 dependence on the step height. The increase of the damping factor with the
170 step height was explained by the fact that the turbulent energy dissipated by
171 the plunging jet - which drives the outgassing process - increases with Δh_i .
172 However, the physical reasons underpinning the linear dependence of f_{s_i} on
173 Δh_i are less trivial, and deserve more investigations. Instead, f_{s_i} was found
174 to be statistically independent on the discharge Q (SI). This is because the
175 rate of gas exchanged through the step is chiefly determined by the turbulent
176 kinetic energy for unit weight of fluid dissipated by the cascade, which is well
177 approximated by the potential energy loss across the drop (Δh_i) regardless of
178 the underlying discharge conditions.

179 The natural steps, where CO_2 production was impeded by manually scouring
180 and removing all the biofilm from the river bed and the downstream pool prior to
181 each measurement, showed a behaviour which was essentially undistinguishable
182 from that of the artificially created steps, suggesting that the procedure used
183 for simulating the steps with pipe diversions didn't introduce significant biases

184 in our analysis. The dependence of f_{s_i} on the step height (Δh_i) in natural
185 steps covered with the plastic film was in line with that observed for the other
186 steps, while the underlying absolute values were slightly lower - likely because
187 in this setting the falling jet adhered to the plastic film, thereby reducing the
188 reaeration rate and the gas evasion in the covered steps.

189 The ratios between f_s and f_c in the three reference reaches analyzed in this
190 study (the reaches *A* and *B*, and the virtual reach *A**, see Methods), were
191 found to be in the range [0.9 – 4.3], depending on the underlying discharge
192 level and the specific reach analyzed. These values were calculated taking into
193 account the number and heights of the steps within each reach, and using a
194 linear empirical function to link f_{s_i} to Δh_i as suggested by our experimental
195 data (dashed line in Figure 2a). Interestingly, most of the gas evasion induced
196 by the steps was associated to the smallest drop heights. In fact, about 60 %
197 of f_s was contributed by steps with heights smaller than 35 cm (SI). Provided
198 that f_s turned out to be independent on Q while f_c peaked for intermediate
199 discharge levels, r had a non monotonic dependence on Q , with higher values
200 observed for low and high streamflow conditions. Remarkably, in all the settings
201 analyzed r was systematically larger than (or close to) unity, thereby indicating
202 that the contribution to gas evasion provided by the steps in the Valfredda creek
203 was remarkable, particularly during low-flow and high-flow conditions.

204 The temporal changes of the dominance ratio - as induced by the hydrocli-
205 matic fluctuations experienced by the focus reach from July 2021 to November
206 2021 - are shown in Fig. 3b and 3c. The temporal pattern of r was reconstructed
207 from the underlying simulated discharge variations (Methods), using the em-
208 pirical relationship between r and Q shown in Figure 3a. Our results clearly
209 indicate that r remained consistently above unity during the whole monitoring
210 season for the three reference river reaches analyzed. The highest values of r
211 were observed during high flow conditions. The resulting average values of the
212 dominance ratio during the whole monitoring period were equal to 2.4 (for the
213 reach *A*), 2.5 (for the reach *B*) and 3.2 (for the reach *A**). We conclude that
214 the outgassing from the study reach of the Valfredda was largely dominated by
215 the local evasion induced by step and pools during the summer and early fall of
216 2021.

217 Implications for large-scale studies

218 On practical grounds, objectively decomposing morphologically-complex chan-
219 nels of the type investigated here into segments and steps may not be straightfor-
220 ward. Here, steps were identified in correspondence of sharp drops in the active
221 river bed, with heights that exceed 10 cm. In our experimental set-up, drops of
222 this type gave rise to aerated falling jets followed by bubbles and/or foaming in
223 the downstream pool, thereby enhancing gas exchange with the atmosphere. In
224 some cases, however, vertical drops could be very small and/or involve only a
225 portion of the active riverbed, creating spatially heterogeneous hydrodynamic
226 conditions which are quite difficult to describe. These hybrid elements were not

227 considered as actual steps in this paper - and their contribution to gas emis-
228 sions was neglected accordingly. Furthermore, owing to spatial heterogeneity of
229 key hydro-morphological characteristics such as slope, discharge and riverbed
230 composition/roughness, the continuous emissions in all the segments belonging
231 to the study reach might not be perfectly represented by the behaviour of the
232 reference segment as postulated by the upscaling procedure proposed in this pa-
233 per. Nevertheless, in spite of all these limitations and the inherent local nature
234 of our study, we believe that the results shown here can be of general validity
235 and properly describe the order of magnitude of the processes involved, as long
236 as the morphology of the mapped steps is in line with previously published data
237 gathered in different regions of the World⁴¹⁻⁴⁵. Therefore, the important contri-
238 bution of small-height steps to the total outgassing observed in the Valfredda is
239 expected to emerge in many high-slope settings, where step and pool bedforms
240 dominate the channel morphology.

241 Our study revealed that the footprint of gas emissions produced by local
242 steps do not vanish at the reach scale, owing to the pronounced concentration
243 damping in correspondence of each step and the high frequency of step-and-pools
244 typical of steep mountain rivers. This key finding prescinds from the precise
245 value taken by the dominance ratio, which might vary from site to site depending
246 on local morphologic and hydraulic features (e.g. channel slope, substrate, step
247 frequency). Besides, the weak dependence of r on the specific reach used for
248 the upscaling suggests that the dominance ratio could be a relatively smooth
249 function of space, as long as it integrates the effect of a sufficiently large number
250 of steps. This feature makes this dimensionless metric suitable to be scaled
251 across space or extrapolated along river networks, e.g. across different channels
252 with similar morphologic characteristics.

253 The damping factors and the dominance ratio quantify the potential out-
254 gassing of different stream elements, whereas the actual value of evaded mass
255 does depend on the spatial correlation between the sources of matter along the
256 stream and the spatial patterns of evasion, here encapsulated by the damping
257 factor²⁴. Given the local nature of gas emissions from step and pools, cas-
258 cades and waterfalls, these geomorphic elements could act as important emission
259 hotspots, where the excess mass transported downstream by the flow is quickly
260 released into the atmosphere (e.g. refs ^{22,26,30}). Therefore, there could be a
261 significant amount of outgassed mass in correspondence of steps, which could
262 be essentially undetectable because of the very short distances traveled from the
263 input (or production) site to the evasion point.

264 Step-and-pools represent crucial morphologic components of high-energy
265 streams, as they drive the transport of matter downstream, and regulate physical
266 or chemical exchanges at the interface with the landscape and the atmosphere.
267 Consequently, a proper characterization of the step frequency and geometry is
268 an important prerequisite for a robust assessment of gas evasion from channel
269 networks. We conclude that key geomorphological factors other than spatially-
270 averaged characteristics of a reach (e.g. mean slope, mean velocity) should be
271 properly incorporated in existing scaling laws currently in use for the prediction
272 of global-scale gas evasion from freshwater systems^{5,25,46-48}. In heterogeneous

273 high-energy streams, in fact, different stretches characterized by the same mean
 274 slope and velocity could lead to highly variable gas evasion rates depending on
 275 their internal configuration. For instance, the 13 *m* reference segment consid-
 276 ered in this study was able to evade - through a total elevation drop of 1.4 *m* -
 277 approximately 8 % to 27 % of the available excess mass, depending on the under-
 278 lying discharge rate. If the same mean slope and elevation drop were obtained
 279 by combining nearly horizontal segments with two steps (SI), the damping fac-
 280 tor would be much higher (from 30 % to 250 % higher than that observed in
 281 the reference segment), with a percentage of excess mass evaded close to 35 %.
 282 The example indicates that small-scale morphological traits of rivers, such as
 283 steps, riffles, cascades or waterfalls can significantly enhance gas exchange at
 284 the water-air interface and - therefore - need to be properly accounted for when
 285 observed patterns of gas evasion are extrapolated outside the specific context
 286 where these data were gathered³⁵. Better characterizing the local morphological
 287 traits of streams that regulate the mass exchanged through water-air interfaces
 288 at regional or global scales could help us to constrain the budget of focal chem-
 289 ical species (e.g. carbon, oxygen, nitrogen) relevant to the water-land-climate
 290 system.
 291

292 **Methods**

293 **Damping factor**

294 The equation governing the spatial patterns of gas concentration in a one-
 295 dimensional system with a curvilinear coordinate x aligned with the main flow
 296 direction, under the assumptions of stationarity (constant flow rate and time-
 297 invariant gas concentrations), no dispersion, no lateral input, absence of internal
 298 gas production reads:

$$u(x) \frac{dC(x)}{dx} + K(x)[C(x) - C_a] = 0, \quad (5)$$

299 where $u(x)$ is the local velocity in the streamline direction, C_a the atmospheric
 300 concentration, $C(x)$ the local water gas concentration, $K(x)$ the local, spatially
 301 variable exchange rate, which is equal to the mass transfer rate k scaled to the
 302 mean water depth. Crucially, the exchange coefficient K embeds the coupled
 303 effects of the mass transfer induced by the turbulence of the flow and that
 304 associated to gas transport mediated by bubbles and foams (if any, see ref⁴⁹).
 305 The solution of Equation (5) is given by Equation (1) in the main text. Therein,
 306 the exponential damping factor can be expressed as:

$$f(x) = \int_0^x \frac{K(x')}{u(x')} dx' = \int_0^{\tau(x)} K(t') dt' = K_{eq}(x) \tau(x), \quad (6)$$

307 where $K_{eq}(x)$ is a weighted spatial average of K in the stretch from 0 to x and
 308 $\tau(x) = \int_0^x 1/u(x) dx$ is the corresponding transit time - the time necessary to

309 travel from 0 to x (SI). Thanks to Equation (1), f_c - the damping factor of an
 310 ideal channel stretch of length L composed by all the continuous segments of
 311 the focus reach - could be expressed as:

$$f_c = \ln \left[\frac{C_0 - C_a}{C_L - C_a} \right], \quad (7)$$

312 where C_0 and C_L are the gas concentrations in the upstream ($x = 0$) and
 313 downstream ($x = L$) sections of the reach. Operationally, given the practical
 314 impossibility of measuring f_{c_i} within all the segments belonging to the focus
 315 reach, f_c was calculated based on the value of the damping factor of a represen-
 316 tative segment (see below) with length ℓ_r , f_{c_r} . The latter was estimated from
 317 Equation (7) through direct gas concentration measurements as:

$$f_{c_r} = \ln \left[\frac{C_0 - C_a}{C_{\ell_r} - C_a} \right], \quad (8)$$

318 where C_0 and C_{ℓ_r} are the concentrations in the upstream and downstream sec-
 319 tions of the reference segment. Then, f_c was calculated from f_{c_r} as:

$$f_c = f_{c_r} \frac{L}{\ell_r}, \quad (9)$$

320 exploiting the additivity of f_{c_i} across multiple segments and assuming that the
 321 exchange rate in the reference segment is equal to the average value of K across
 322 all the segments contained in the focus reach (SI).

323 Similarly, f_{s_i} was calculated from Equation (1) as:

$$f_{s_i} = \ln \left[\frac{C_{u_i} - C_a}{C_{d_i} - C_a} \right], \quad (10)$$

324 where C_{u_i} and C_{d_i} represent the water gas concentration upstream and down-
 325 stream of the step i . The damping factor of a sequence of N steps was then
 326 evaluated summing up the damping factors of all the individual steps f_{s_i} (i.e.
 327 $f_s = \sum_i f_{s_i}$, see SI), as in Equation (4) of the main text.

328 Study site and focus reach

329 The study site selected in this paper is a step-pool channel of the Rio Valfredda,
 330 a high-gradient headwater catchment of the Piave river basin, in the Italian
 331 Alps^{35,50}. The climate of the site is typically alpine: precipitation is relatively
 332 high throughout the year (annual rainfall > 1400 mm), with significant snowfall
 333 during winter and melting in spring⁵¹. The selected reach is 1.36 km long, and
 334 its elevation ranges from 1911 to 1720 m a.s.l., with a mean slope of 0.14 m/m
 335 (Fig. 1a). The reach was selected because of its accessibility and the significant
 336 CO_2 concentrations observed therein. The river bed is steeper in the upstream
 337 part, where it flows southwards. Then, the reach runs south-east across some
 338 pastures and a mixed larch-spur forest. The reach is fed at its source by a

339 groundwater spring. The discharge weakly increases downstream, owing to the
340 interplay between the losing bed and the hillslope lateral input. The stream bed
341 is silty and dominated by boulders, cobbles and wooden logs of different size
342 that originated several steps and pools. About 300 *m* upstream of its confluence
343 with the Valfredda, the channel was almost inaccessible due to the presence of
344 a landslide and several fallen trunks. Therefore, the analysis was concentrated
345 in the upper portion of the channel (reach *A*, with a total length $L_A = 1.06$ *km*
346 (blue line in Fig. 1a)). To capture the footprint of the bed geomorphology
347 on gas evasion, the reach *A* was decomposed into 270 segments and 271 steps
348 (Fig. 1b). A representative segment with a slope (discharge) similar to the
349 mean slope (mean discharge) of the continuous segments within the study reach
350 was identified in the middle part of the reach (46°22'50"N, 11°49'39"E, see Fig.
351 1a). This representative segment has an average slope $i_{c,r}$ of 0.108 m/m and a
352 length ℓ_r of 13 *m* (inset of Fig. 1a and 1c).

353 Discharge measurements

354 We performed several volumetric measurements of the discharge rate, Q , at the
355 two end points of the representative segment. This was done recording the filling
356 time of a graduated tank in correspondence of the upstream and downstream
357 section of the representative segment. We performed 10 measurements between
358 July and October 2021, with observed discharge values between 0.2 and 3.2 *l/s*
359 (see Table S1). Each measurement was the average among at least 5 different
360 replicas performed in both the locations within one hour.

361 Travel time measurements and estimation of the relevant hydraulic 362 properties

363 The water travel time along the representative segment, τ_{ℓ_r} , was measured
364 through instantaneous injections of a diluted sodium chloride (NaCl) solution in
365 the upstream cross section of the segment. The temporal variations of specific
366 conductivity at the outlet of the segment were then measured using a multi-
367 parameter sonde (YSI EXO2). The travel time was recorded both under natural
368 conditions and after having covered the stream bed with the plastic film (Fig.
369 S2 and S3). The longitudinal mean velocity, u , was estimated as the segment
370 length, ℓ_r , divided by the observed travel time. The procedure allowed us to
371 verify that the mean velocity along the reference segment was not significantly
372 impacted by the presence of the plastic film. Moreover, after having covered
373 the river bed with the plastic film, we also measured the mean width, W , and
374 water depth, H of the flow. This was done by taking spatial averages of the
375 local values observed in different cross sections along the segment. The obtained
376 hydraulic geometry scaling relationships (Fig. S4) were found to be in line with
377 the relationships proposed in the literature for mountain streams⁸, thereby sug-
378 gesting that the flow conditions in place during the CO_2 measurements were
379 nearly-natural.

380 **CO_2 concentration measurements**

381 Paired upstream and downstream CO_2 concentration measurements were taken
382 in the representative segment and in 19 steps (see SI), using a membrane-based
383 NDIR sensor, the Mini CO_2 TM designed by ProOceanus Systems Inc., Bridgewater,
384 Canada. The instrument has a tubular shape, 370 mm long and 53.4 mm
385 in diameter, and uses infrared detection to measure the partial pressure of dis-
386 solved CO_2 . Once the internal gas is fully equilibrated with the surrounding
387 water (typically 10-15 min after the deployment), NDIR measurement on the
388 equilibrated internal gas is taken at a wavelength of 4.26 μm close to the ab-
389 sorption band of CO_2 at a controlled optical cell temperature. The time of
390 deployment in a given position ranged from 30 to 90 minutes. To eliminate the
391 confounding effect of high-frequency fluctuations in the recorded CO_2 signal of
392 the Mini CO_2 TM sensor (see Figure 1c-1e for an example), at least 20 minutes
393 of continuous measurements at steady-state were gathered, from which we esti-
394 mated the probability density function of CO_2 concentrations and the related
395 mean. Steady-state conditions were pre-identified based on the temporal pat-
396 terns of the long-term average of the signal. The steady state mean was then
397 taken as representative of the equilibrium carbon dioxide concentration in water
398 (Fig. S8 and S9). Paired upstream-downstream concentrations were gathered
399 within 120 mins from each other, so as to reduce as much as possible spurious
400 effects induced by diel variations of water CO_2 concentration.

401 **Estimating f_{s_i} AND f_{c_r}**

402 In this study the damping factors were evaluated on a purely experimental
403 basis. The damping factor of individual steps f_{s_i} was estimated from up-
404 stream/downstream CO_2 measurements via Equation (10), considering three
405 different step types: i) 11 simulated steps, which were created forcing the water
406 into pipes and then letting the water flow hit a covered portion of channel bed
407 from a given height, so as to reproduce the behaviour of a falling jet of a nat-
408 ural drop with the desired Δh (Fig. 1d and S5); ii) 4 covered steps, obtained
409 by folding natural steps with a thin plastic film, which was carefully shaped
410 around the actual channel bed in the ramp, the step and the downstream pool
411 (Fig. S6e); iii) 4 natural steps belonging to the focus reach (Fig. 1e and S6f),
412 which were scoured to remove the existing biofilm prior to each measurement.
413 These precautions allowed us to eliminate the effect of CO_2 production in all the
414 analyzed steps, while assuring natural hydraulic conditions in all the measured
415 steps. Further details on the experimental setup are available in SI.

416 As per the estimate of the damping factor in the reach segments, owing to
417 the practical impossibility to quantify the outgassing within all the segments
418 contained in the study reach, f_c was estimated from the observed concentration
419 drop in the representative segment by means of an upscaling procedure (Equa-
420 tion (9)). To measure f_{c_r} , prior to each field measurement the stream bed was
421 covered by a plastic film, to avoid lateral input of water and CO_2 and internal
422 production induced by the ecosystem metabolism (Fig. S7). Then upstream vs

423 downstream CO_2 concentrations were measured under different hydrologic con-
 424 ditions (discharge range: from 0.19 to 2.11 l/s), and Equation (8) was applied.
 425 In the light of the constraint placed by the specific slope of the representative
 426 segment and the mean slope of the segments contained in the reach A , we an-
 427 alyzed three different scenarios, in which the damping factor of the reference
 428 segment was upscaled via Equation (9) referred to three different reaches: i) the
 429 whole reach A , which has a length L_A of 1060 m and a mean slope of its contin-
 430 uous segments $i_{c,A}$ of 0.081 m/m ; because $i_{c,A}$ is significantly smaller than the
 431 slope of ℓ_r , in this case f_c should be overestimated; ii) the reach B , including 130
 432 steps, which has a length L_B equal to 543 m and is characterized by an average
 433 slope of its continuous segments $i_{c,B} = 0.101 m/m$ - quite close to the slope of
 434 the representative segment; iii) the reach A^* , an idealized reach characterized
 435 by: i) the same elevation drop as that observed in between the two endpoints
 436 of reach A ; and ii) segments that have the same slope of the reference segment.
 437 A^* is 769 m long, has a slope of its segments equal to 0.108 m/m (by definition
 438 equal to $i_{c,r}$) and contains all the 271 steps of reach A (further details in SI).
 439 Since the mean slope of the segments included in B and A^* is closer to the
 440 actual slope of the reference segment, the corresponding estimates of f_c and r
 441 should be more reliable in this case. Note that the mean slope of the segments
 442 included in a reach was calculated starting from the mean slope of the overall
 443 reach, taking into account the heights of all the steps included in that reach and
 444 assuming a longitudinal size of 10 cm for each step (SI). This implies that the
 445 length of the segments included in a reach is slightly smaller than the length of
 446 the whole reach (SI).

447 **Morphological survey**

448 The dominance ratio depends on the spatial frequency and the height distribu-
 449 tion of the steps in the focus reach. We collected geometrical data about the step
 450 geometry during field surveys performed under very dry conditions. For each
 451 step we measured the step drop height, Δh , corresponding to the gap of water
 452 surface elevation in each nearly vertical fall with a drop higher than 10 cm . We
 453 mapped 271 steps in 1060 m with an average Δh equal 23.7 cm (Fig. S10). The
 454 frequency distribution of the step height was monotonically decreasing, with
 455 47.6 % of steps in range 0 – 15 cm , and 70.5 % of steps in range 0 – 25 cm .
 456 Elevations, lengths and slopes of the relevant reaches were estimated through a
 457 high resolution (1 m) *DTM*.

458 **Streamflow regime**

459 The streamflow regime was estimated from rainfall data using a simple rainfall-
 460 runoff model. In particular, we used daily precipitation heights P [mm] recorded
 461 during the summer and fall seasons (Jun to October 2021) collected by a weather
 462 station of the Veneto Region Environmental Protection Agency (ARPAV) lo-
 463 cated in Falcade, 4.5 km far away from the catchment centroid. Discharge
 464 timeseries in the focus reach were then simulated using an exponential IUH

465 applied to the censored precipitation time series as:

$$Q(t) = A \sum_i j_i k e^{-k(t-t_i)}, \quad (11)$$

466 where A is the catchment area, k is the recession rate, t_i is the occurrence
467 time of effective rain events, j_i is the effective (i.e. censored) rain depth ($j_i =$
468 $\max(P_i - \phi, 0)$), with P_i representing the total rain depth and ϕ the censoring
469 threshold embedding soil moisture dynamics⁵². ϕ and k were calibrated against
470 the discharge observations available (Figure 3c and SI).

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475 at <http://researchdata.cab.unipd.it/id/eprint/619>, reference number 619.

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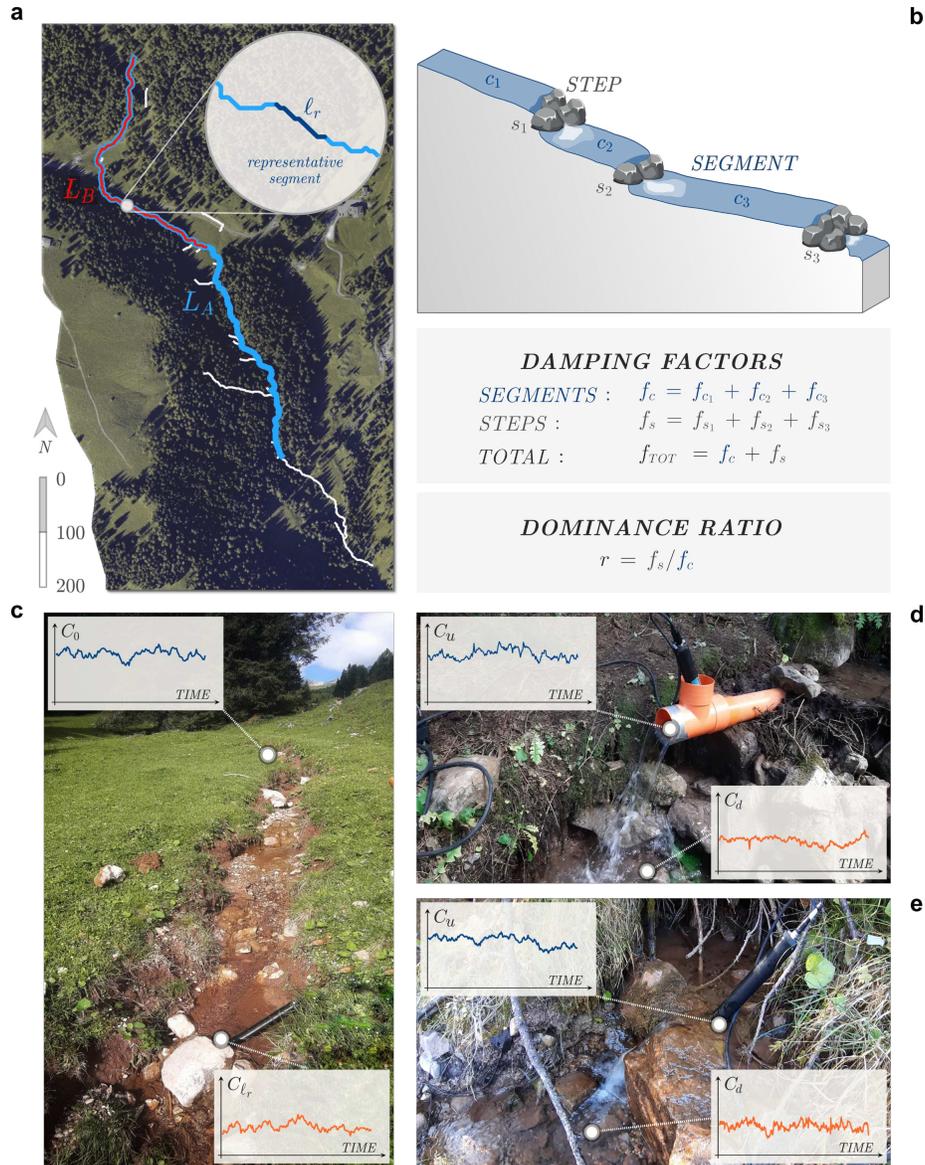


Figure 1: **Focus reach of the Valfredda creek and experimental set up.**

a) planar view of the reach selected for this study, shown with an orthophoto of the eastern part of the Valfredda catchment in background. In particular, the reach A, with length $L_A = 1060\text{ m}$, is indicated in light blue, while the reach B, with length $L_B = 543\text{ m}$ is shown in red. The inset instead shows a planar view of the representative segment without steps, which has a length $\ell_r = 13\text{ m}$; b) schematic of the decomposition of a reach into segments and steps, and aggregation of the damping factors f within a complex system, which leads to the definition of the dominance ratio r ; c) overview of the reference segment, with the insets showing temporal dynamics of the water CO_2 concentration upstream (C_0) and downstream (C_{ℓ_r}) of the segment; d) artificial step, created by forcing the stream into a pipe, and then covering the downstream river bed with a plastic film; the insets show representative temporal dynamics of the water CO_2 concentration upstream (C_u) and downstream (C_d) of the step. e) same as d) but for a scoured natural step.

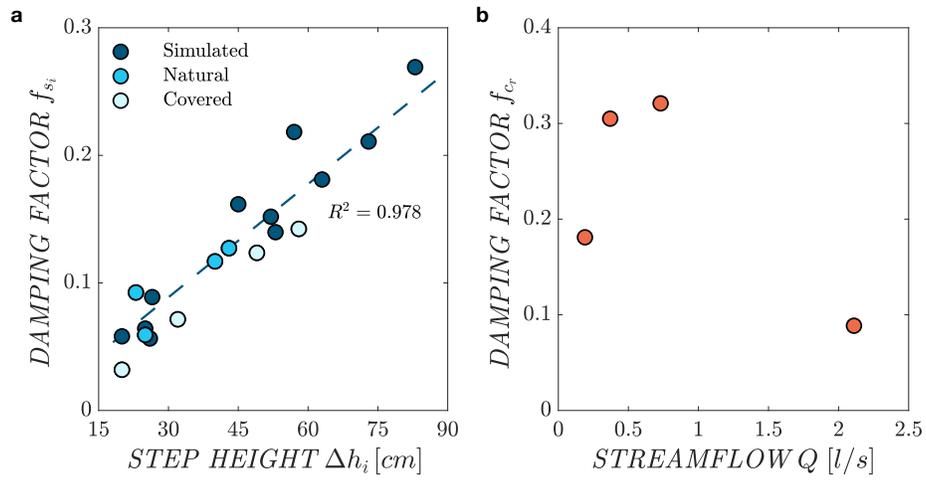


Figure 2: **Damping factors for the steps and the reference segment.** a) damping factor of the steps f_{s_i} as a function of the step height Δh_i for scoured natural steps (pale blue circles), natural covered steps (white circles) and artificially simulated steps (navy blue circles). Also shown is the linear regression $f_{s_i} = 0.3 \Delta h_i$ (with Δh_i in m), and the regression coefficient R^2 ; b) damping factor in the reference segment, f_{c_r} , for different discharge conditions in the range from 0.19 to 2.11 l/s .

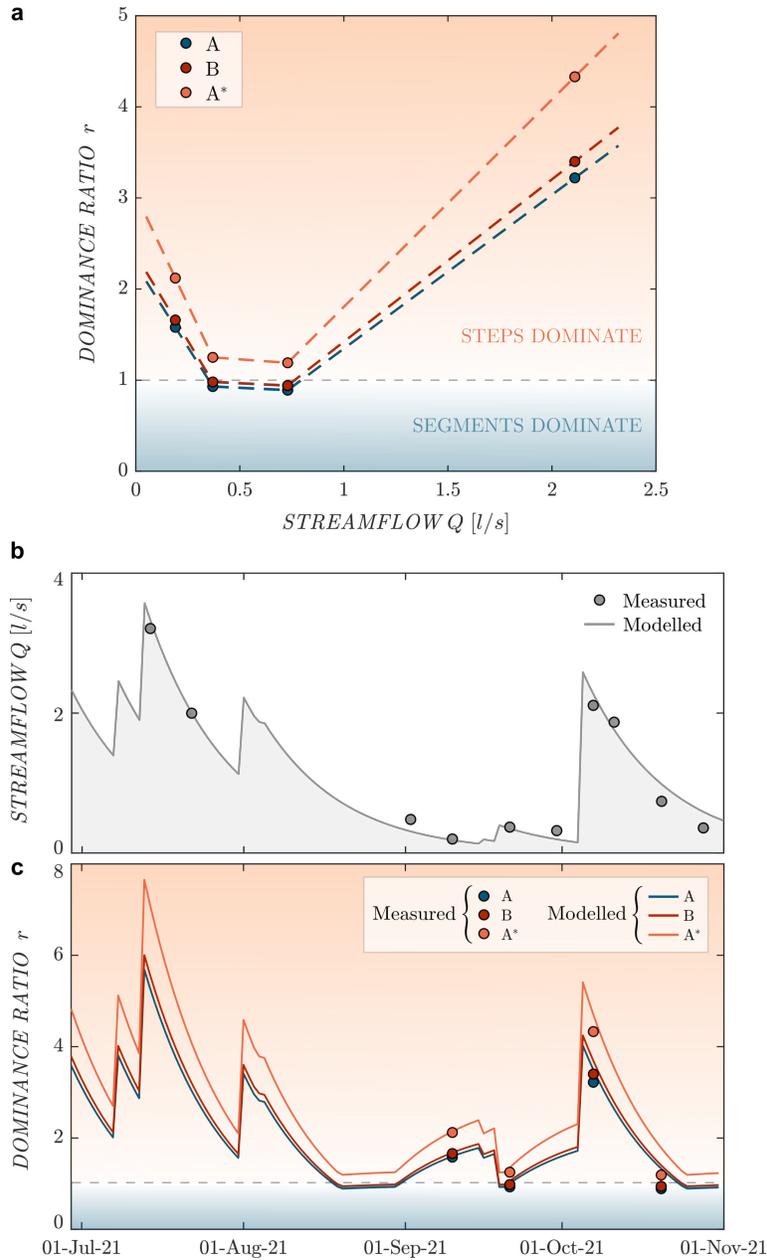


Figure 3: **Dominance ratio variations induced by changes in the underlying hydrologic conditions.** a) dominance ratio as a function of the stream discharge Q in the reaches A (blue petroleum), B (red) and A^* (orange). Circles represent the observed values of r , the dashed lines show a linear piece-wise interpolation. When $r > 1$ the steps dominate the outgassing (upper ibis-shaped region of the plot); b) temporal dynamics of the discharge in the reference segment of the focus reach from July 1 to Nov 1, 2021. Observed streamflows are indicated as grey dots, while the simulated discharges (see Methods) are shown by the shaded-grey region; c) temporal dynamics of the dominance ratio during the same time window shown in b) for the reaches A (blue petroleum), B (red) and A^* (orange). Circles refer to the observed values, while the solid lines refer to simulated values estimated based on the simulated discharges (shown in b) using the piece-wise linear interpolation between r and Q (shown in a). When $r > 1$ the steps dominate the outgassing (upper ibis-shaped region of the plot)

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