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Denise Degen (✉ denise.degen@cgre.rwth-aachen.de)

RWTH Aachen University <https://orcid.org/0000-0002-7932-6251>

Mauro Cacace

Helmholtz Centre Potsdam (GFZ) - German Research Centre for Geosciences <https://orcid.org/0000-0001-6101-9918>

Florian Wellmann

RWTH Aachen University

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3D Multi-Physics Uncertainty Quantification using Physics-Based Machine Learning

Denise Degen^{1,*}, Mauro Cacace², and Florian Wellmann^{1,3}

¹RWTH Aachen University, Computational Geoscience, Geothermics and Reservoir Geophysics (CGGR), Mathieustraße 30, 52074 Aachen, Germany

²GFZ German Research Centre for Geosciences, Telegrafenberg, 14473 Potsdam, Germany

³Fraunhofer Research Institution for Energy Infrastructures and Geothermal Systems (IEG), Am Hochschulcampus 1, 44801 Bochum, Germany

*denise.degen@cgre.rwth-aachen.de

ABSTRACT

Quantitative predictions of the physical state of the Earth's subsurface are routinely based on numerical solutions of complex coupled partial differential equations together with estimates of the uncertainties in the material parameters. The resulting high-dimensional problems are computationally prohibitive even for state-of-the-art solver solutions. In this study, we introduce a hybrid physics-based machine learning technique, the non-intrusive reduced basis method, to construct reliable, scalable, and interpretable surrogate models. Our approach, to combine physical process models with data-driven machine learning techniques, allows us to overcome limitations specific to each individual component, and it enables us to carry out probabilistic analyses, such as global sensitivity studies and uncertainty quantification for real-case non-linearly coupled physical problems. It additionally provides orders of magnitude computational gain, while maintaining an accuracy higher than measurement errors. Although in this study we use a thermo-hydro-mechanical reservoir application to illustrate these features, all the theory described is equally valid and applicable to a wider range of geosciences applications.

1 Physics-based models, meaning models that are governed by partial differential equations, are extensively used in geosciences.
2 These models are all based on well-established physical principles, which provide a high-fidelity approximation of the Earth's
3 system dynamics. Typical problems in geosciences rely on a large number of parameters in order to describe the multidimen-
4 sional character of the underlying processes in both space and time¹⁻⁶. Parameter estimation is often based on limited data,
5 which, together with our incomplete knowledge of the heterogeneous physics at play, degrades the performance of currently
6 adopted solutions⁷. The challenge here is to understand both the structure of the prediction system and to quantify its uncertainty.
7 Uncertainty quantification (UQ) requires to develop a confidence metric to measure predictions, their validation against available
8 data, and their sensitivity upon variations in the parameter space. This prerequisite often results in high-dimensional problems,
9 which become computationally intractable if based on traditional probabilistic methods such as Markov chain Monte Carlo⁸.
10 Any advancement in UQ analyses requires to improve upon traditional statistical error analysis. A common strategy is to rely
11 on surrogate models, constructed either for the entire state^{9,10} or limited to the observation space alone¹¹⁻¹³. Surrogate models
12 for the entire state are based on the physical model, and share the advantage of preserving the underlying physics. However,
13 these models are limited in their range of applications, e.g., they can not generally be applied to hyperbolic and nonlinear
14 problems as typically encountered in geoscientific applications^{10,14}. Surrogate models for the observation space alone are not
15 applicable for nonlinear problems either. In addition, these models do not preserve the physics.
16 Recent advances in statistical and machine learning (ML) methods, together with increased computational resources offer a
17 new opportunity to expand our knowledge about the system's Earth from available data¹⁵⁻¹⁷. Several free computational ML
18 libraries (e.g. PyTorch¹⁸, TensorFlow¹⁹) are nowadays available to allow automatic management and integration of the steadily
19 increasing stream of geospatial data. However, the requirement of a reliable and robust surrogate model has so far limited a
20 straightforward application of ML techniques for UQ analysis in geosciences.
21 In this study, we propose the non-intrusive reduced basis (NI-RB) method, a hybrid physics-based machine learning tech-
22 nique^{14,17,20}, for the construction of reliable and robust surrogate models. By combining physics-based modeling with
23 data-driven concepts it extends the range of applicability of each individual method, provides substantial computational gain
24 with negligible degradation in the model accuracy, and enables the application of UQ for real case 3D multi-physics simulations.
25 In addition, its theoretical formulation is independent of the specific physics at play, and therefore the NI-RB is applicable to a
26 wide range of geoscientific applications.

27 Construction and Evaluation of the Surrogate Model

28 To illustrate the applicability of the NI-RB method, we selected a real-case reservoir problem: a hydrofracture treatment
 29 performed at the geothermal research facility of Groß Schönebeck in northern Germany²¹, Fig. 1. The targeted problem
 30 exhibits all features of interest, that is, high dimensionality (heterogeneous distribution of material properties) and tightly
 31 coupled multi-physics (fluid dynamics coupled to non-linear solid thermo-mechanics). For the model validation, we rely on
 32 matching the pore pressure response at the monitoring well E GrSk 3_90 (marked in blue in Fig. 1b), which we consider
 33 as representative of the far-field response over the entire reservoir to the hydraulic stimulation at the opposite well (500 m
 34 distance)²². The evolution in time of the monitored over-pressure (Fig. 1a) displays a rather complex pattern, with local rapid
 35 changes superimposed on a monotonic long-term diffusive trend. This is an important aspect since deep learning algorithms
 36 would perform poorly in matching such a complex time series (with the amount of data at hand) as they would for example fail
 37 to predict the timing and wavelength of the rapid bursts in the system response²³. Instead, we are able to replicate the pressure
 38 evolution curve both with respect to the long-term and short-term characteristics. The constructed surrogate model provides
 39 an accuracy of $2.29 \cdot 10^{-7}$ for the training set, and $2.48 \cdot 10^{-4}$ for the validation set, and it requires five basis functions for its
 40 construction. Deviations between measurements (solid black curve in Fig. 1a) and the FE solutions (colored solid curves in Fig.
 41 1a) are instead in the order of 10^{-1} . Our surrogate model, therefore, provides a higher accuracy than the accuracy of the entire
 42 model, and we can consider the approximation errors from the surrogate model as negligible.
 43 So far, we have limited our discussion to global values, that is we have been interested in the accuracy of the RB model for
 44 the entire training and validation data set. In an attempt to understand the time evolution in the error datasets, we investigate
 45 five simulations that we chose randomly from the validation dataset. For all five realizations, we obtain maxima in the errors
 46 within the time period between 2.1 and 2.8 days, which is exactly during the time window that features the highest variability in
 47 the model response. However, also for this time period, the introduced errors are smaller than the deviations between the FE
 48 solutions and the observation data. We consider this as proof that the NI-RB method is suited for the construction of reliable
 49 surrogate models for tightly coupled multi-physics problems even for systems that showcase a rather complex and variable
 response in time and space.

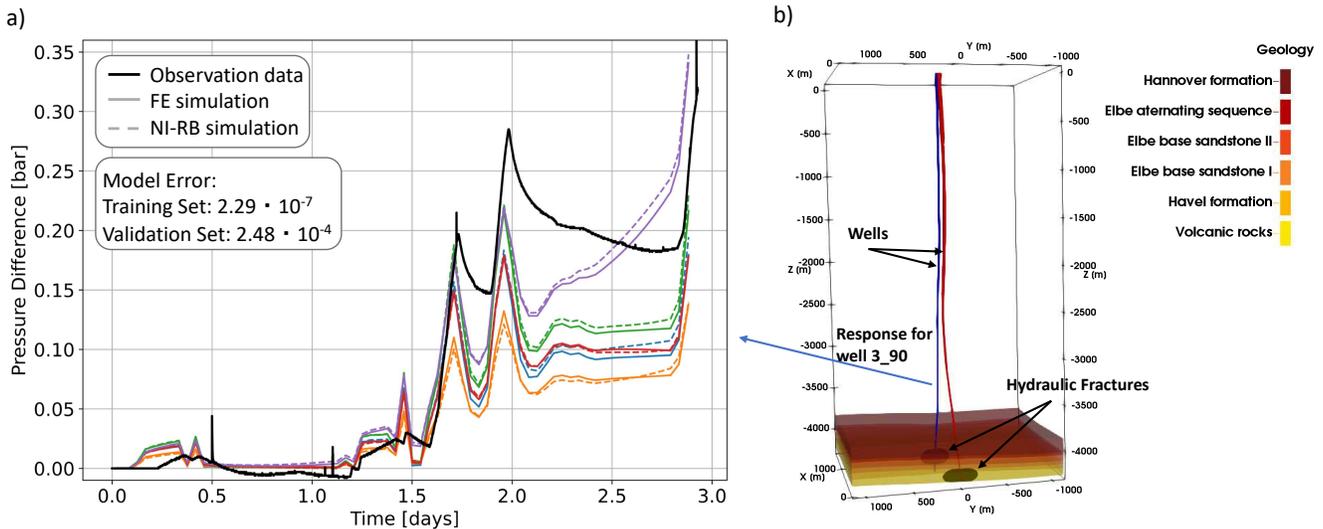


Figure 1. a) Representation of the accuracy of the surrogate model of Groß Schönebeck for five different parameter realizations of the validation data set. The solid black line denotes the observation data, the colored solid lines the FE solutions and the dashed lines the RB solutions. b) Overview of the geological reservoir model of Groß Schönebeck. Next to the six geological units also the two wells and the hydraulic fractures are indicated.

50

51 Global Sensitivity Analysis

52 In order to identify the model parameters having the highest influence on the resulting pressure evolution, we carry out a global
 53 sensitivity analysis. For our problem we allow six parameters to vary, that is, the solid bulk modulus, the thermal expansion
 54 coefficient, the permeability and the porosity of the Elbe base sandstone I and the Volcanic rocks, the latter being the reservoir
 55 layers targeted during the hydraulic stimulation. In Fig. 2 we list all parameters with their range of variations used in the
 56 sensitivity analysis. The goal of the sensitivity analysis is to identify the parameters that minimize the misfit between the

57 observation data and the modeling results the most. To this end, we carry out our sensitivity analysis with respect to the mean
 58 squared difference between simulated and measured data. Given the error of the surrogate models, we consider a threshold
 59 value of $5 \cdot 10^{-2}$ for the tolerance in our analysis.

60 We found that the parameters that exert the highest influence on the system response are the permeability and the porosity of
 61 the Elbe base sandstone I and the permeability of the underlying volcanic rocks, with the sandstone permeability exerting the
 62 highest influence and that of the Volcanics the lowest influence. We also note that the obtained difference between the first-
 63 and total-order contributions are negligibly small for all these three parameters and that consequently, we have only a minor
 64 parameter correlation.

65 We can further observe that thermal effects (quantified by the impact of assumed variations in the thermal expansion coefficient)
 66 do not affect the pressure distribution significantly and that also the mechanical parameter considered (solid bulk modulus) only
 67 has a minor influence. Hence, the thermo-elastic stress transfer is not impacting the pressure response, which is also confirmed
 by a study matching the microseismicity following the treatment²⁴.

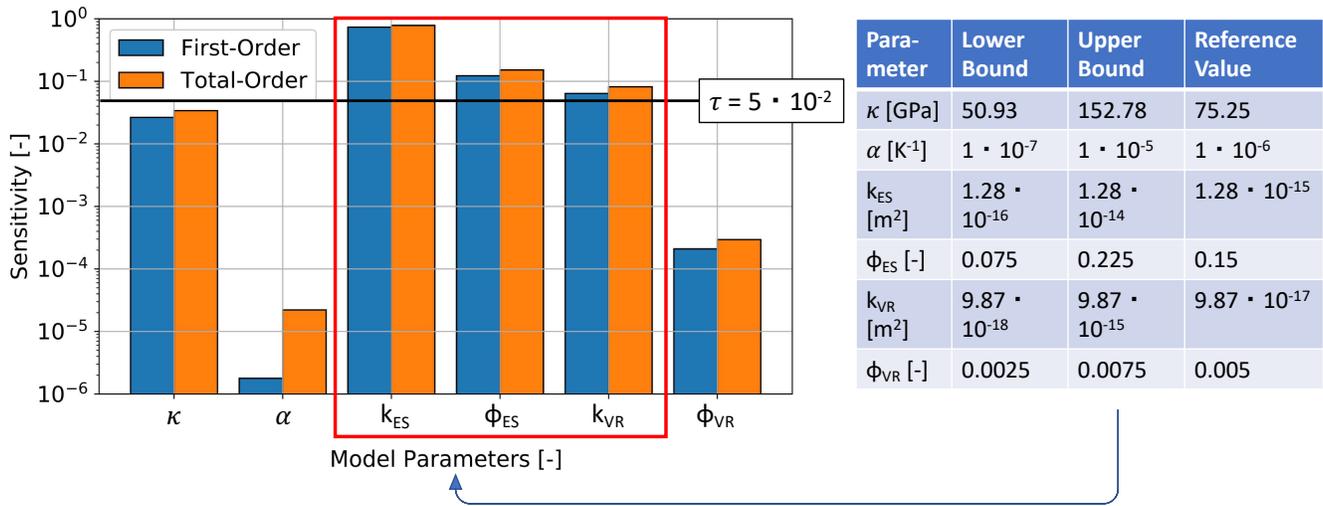


Figure 2. Global sensitivity analysis for the model of Groß Schönebeck with a threshold value of $5 \cdot 10^{-2}$, as well as the lower and upper bounds presented on the right side. Note that κ denotes the solid bulk modulus, α the thermal expansion coefficient, k the permeability, ϕ the porosity, and the subscripts ES and VR the Elbe sandstone I layer, and the Volcanic rocks, respectively.

68

69 Uncertainty Quantification

70 The global sensitivity analysis has led to a reduction of the dimension of the parameter space, and accordingly of the sampling
 71 ensemble required for UQ. Therefore, we focus our UQ on the resulting three most influencing parameters; the permeability
 72 and porosity of the Elbe base sandstone I layer and the permeability of the Volcanics. All other parameters are kept constant,
 73 according to their reference values (Fig. 2). We start our UQ based on the pressure response to the stimulation at the well
 74 3_90 as computed before and after the stochastic model calibration. The results are summarized in Fig. 3a. The green curve
 75 represents the “trial-and-error” calibration obtained by a previous work²², which we take in this study as the mean of the prior.
 76 The solid orange curve is our posterior mean, while the dashed orange curve is the 95 % quantile of the pressure. Comparing
 77 the simulations to the measurements, we note the presence of a source of epistemic uncertainties, which is resulting from the
 78 input fracture model used for the modeling. This overestimates the fracture closure after each stimulation stage, visible as the
 79 systematic overshooting of the modeled pressure relaxation after each pressure peak during each cycle. However, this source of
 80 uncertainty does not impact any results presented in this manuscript.

81 In Fig. 3b-e, we show exemplarily the posterior analysis for the porosity of the Elbe base sandstone I. The complete
 82 information for all remaining parameters is provided in the Supplementary Material. We observe a shift in the posterior mean
 83 of the permeability and a significant reduction in the uncertainty. This indicates that incorporating observation data yields a
 84 reduction in uncertainties. The z-score values do not exceed a range between -2 to +2. The low autocorrelation values, as well
 85 as the trace, demonstrate a suitable mixing and that the performed uncertainty quantification is robust.

86 With this approach, we can demonstrate an efficient and robust UQ workflow, even for this complex non-linear physical
 87 simulation. Note that the execution of the UQ requires about 20 min for 300,000 iterations. As a comparison the same analysis
 88 using the full model would have taken more than 50 compute-years, already accounting for a possible parallelization on

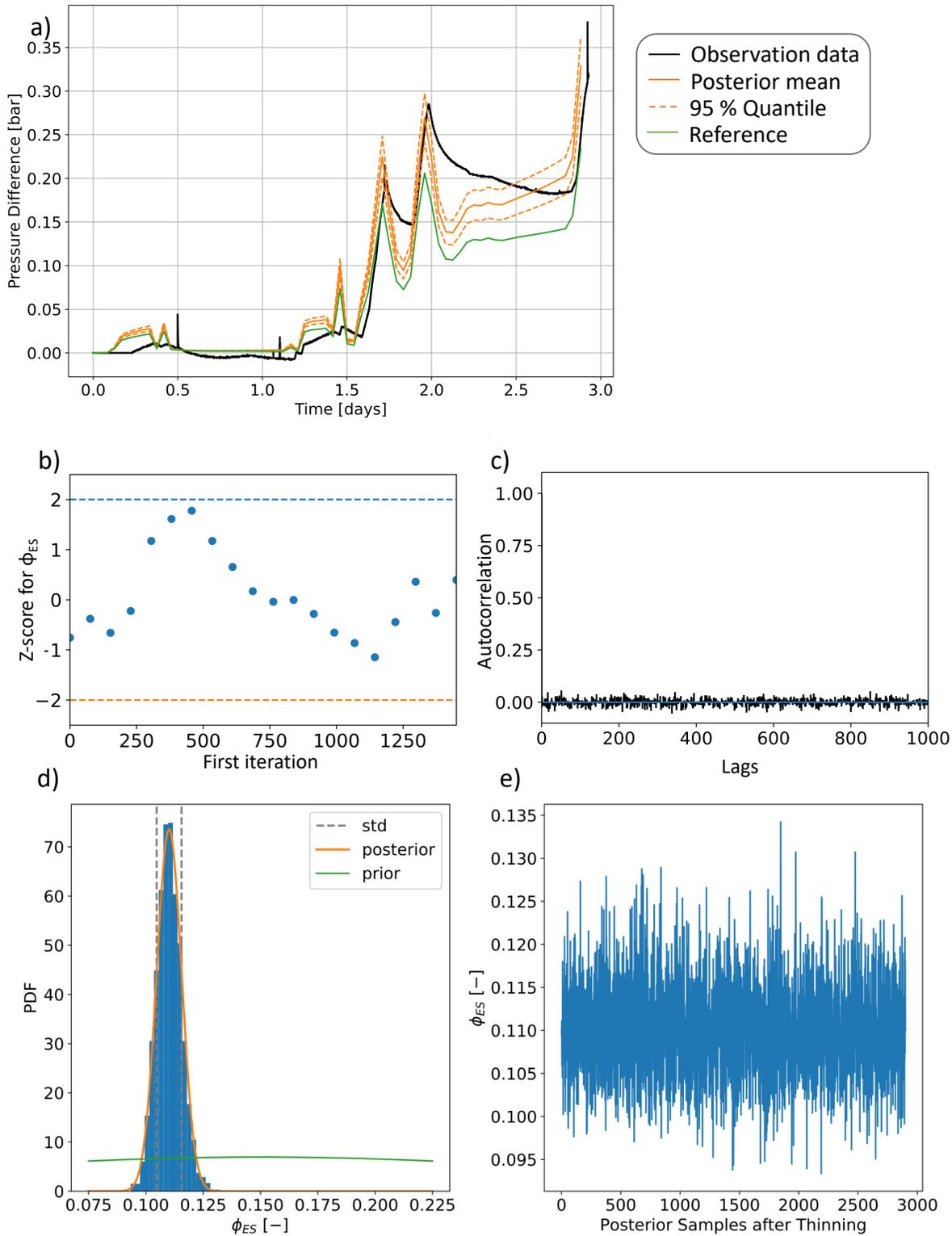


Figure 3. Posterior analysis of pressure response and exemplars of the porosity of the Elbe base sandstone I (ϕ_{ES}). Shown are the a) pressure distribution for the well 3_90 (both posterior mean and 95 % quantile), b) Geweke Plot, c) autocorrelation, d) posterior parameter distributions, and e) the trace.

high-performance infrastructures. Except for the large reduction in computation time, we observe that the UQ greatly reduces the uncertainties of the input parameters, as seen in the posterior distribution of the parameters and the small variation ranges in the distribution of the pressure response. For the UQ, the usage of a physics-based ML approach has clear advantages since we preserve the main physical characteristics and use the ML technique for the projection step only. Using either purely data-driven approaches or other physics-based methods, such as physics-informed neural networks (PINNs), yields degradations of the UQ analysis, especially when considering noisy data²⁵. This is why approaches such as BPINNs have been developed²⁵, which overcome this issue but have the disadvantage of including the hyperparameters in the UQ analysis, which results in a loss of the physical meaning of the system. In contrast, the non-intrusive RB method results in surrogate models that can be used in a flexible way for common parameter estimation methods, SAs, and UQ.

Discussion

In this paper, we presented the non-intrusive RB method, a hybrid physics-based machine learning method that enables the construction of robust, scalable, and interpretable surrogate models for a wide range of geoscientific applications where multiple simulations are required. To showcase the potential of the method, we use the multi-physics study of Groß Schönebeck because it compromises the typical challenges of geoscientific applications: being high-dimensional and nonlinear, and considers a tightly coupled multi-physics setting. For this model, we reduce the computational time from about 1.5 h (already using parallelization on High-Performance infrastructures) to about 3 ms (desktop computer). This reduction in computation time is a result of an enormous reduction in the dimensionality from 278,095 degrees of freedom to only five, yielding a speed-up of six orders of magnitude.

In contrast to classical machine learning or PINNs, the non-intrusive RB method does not treat physics simply as data or as a constrain. The POD step performed as a first step extracts the characteristic physical behavior and the machine learning technique is used to determine the weight of these characteristics only. This has the advantage that i) the amount of data required for the methodology is greatly reduced, and ii) the approximation errors are mainly related to the weighting and therefore the main characteristics of the response are preserved. The aspect of data is especially important since this application shows that the data generation is costly, which holds for geoscientific applications in general. In this work, 150 samples are sufficient for the generation of the surrogate model. For data-driven approaches, we would expect the amount of data to be magnitudes higher. Generally, machine learning has the disadvantage of not being rigorously proven, which is problematic for predictions and their reliability. This problem is for the non-intrusive RB method also the case, although it does affect the weighting step only. But the non-intrusive RB method is based on a rigorously proven method^{9,10,26}. So, by simplifying the problem setting to an elliptic or parabolic PDE, we can always fall back on the proven methodology. Finally, the presented non-intrusive RB method constructs a map from the model parameters (e.g., rock properties) to the model response making it ideally suited for parameter estimation studies. Other physics-based machine learning methods such as PINNs are meant for state estimation problems²⁷. Hence, modifications are required before being applied to parameter estimation studies. To sum up, the non-intrusive RB method is well suited for the construction of surrogate models for a wide range of geoscientific applications, enabling the possibility to gain new insights into this field.

Methods

In the following, we introduce the governing equations, and the physics-based machine learning method, namely the non-intrusive reduced basis method. We use the method to construct reliable and rigorous surrogate models, which is of utmost importance for scientific machine learning¹⁷. Furthermore, we briefly explain the concepts of uncertainty quantification and global sensitivity studies.

Governing Equations

For the case study of Groß Schönebeck, we consider a coupled thermo-hydro-mechanical forward problem as implemented in the software GOLEM¹, which is based on the finite element solver MOOSE²⁸. The formulation of the fluid pressure p_f is derived from the fluid mass balance and can be expressed as¹:

$$\frac{1}{M_b} \frac{\partial p_f}{\partial t} + \nabla \cdot q_D = 0, \quad (1)$$

$$\text{with } q_D = -\frac{k}{\mu_f} (\nabla p_f - \rho_f g). \quad (2)$$

Here, M_b denotes the Biot modulus, t the time, k the permeability, μ the dynamic viscosity, ρ the density, g the gravitational acceleration, and the subscript f the fluid component.

The equation for the temperature is an expression of the energy balance, and is given by¹:

$$(\rho c)_b \frac{\partial T}{\partial t} + \nabla \cdot \left((\rho c)_f q_D T - \lambda_b \nabla T \right) = 0, \quad (3)$$

131 where c is the specific heat capacity, T the temperature, λ the thermal conductivity, and the subscript b denotes the bulk
132 component.

We take the momentum balance to derive the equation for the effective stress σ' , yielding¹:

$$\nabla \cdot (\sigma' - \beta p_f \mathbb{1}) + \rho_b g = 0. \quad (4)$$

133 The Biot coefficient is denoted by β , and $\mathbb{1}$ is the rank-two identity matrix. For further information regarding the coupling
134 scheme, we refer to Cacace and Jacques¹.

135 Non-Intrusive Reduced Basis Method

136 The field of machine learning is rapidly growing and more and more incorporated into scientific applications. This growth is
137 also tightly linked to large open-source Python libraries such as TensorFlow and PyTorch, making the methodologies accessible
138 in a user-friendly black-box approach¹⁶. Nonetheless, machine learning for scientific purposes faces major challenges, when
139 applied in a black-box model^{17,20,29}. To ensure the scientific merit, models resulting from the applied methodologies must be
140 rigorous, reliable, scalable, generalizable, and interpretable^{17,20,29}.

141 Machine learning methods are commonly entirely data-driven methods. This poses a major challenge for many geophysical
142 applications that investigate subsurface processes since they face the problem of data sparsity, as common in many other
143 physical applications. Instead of data, these applications usually have a good understanding of the governing physical processes
144 (although subjected to uncertainties). This is the reason they classically focus on physics-based approaches¹⁷. However,
145 computing these models is computationally very demanding, making intensive parameter estimation problems prohibitive^{30,31}.
146 Therefore, a common procedure is to construct surrogate models. These surrogate models can be physics-based^{9,10,26} or data-
147 driven^{11,12,16,32}. We already addressed the shortcoming of data-driven approaches in our field. However, also physics-based
148 surrogate models are problematic since they often have only limited applicability to nonlinear hyperbolic partial differential
149 equations (PDEs). Hence, neither of the methods is suitable for the application presented here. Therefore, we propose the usage
150 of a physics-based machine learning method, namely the non-intrusive reduced basis method^{14,33}.

151 The non-intrusive reduced basis method (NI-RB) can be seen as a hybrid approach combining physics-based and data-
152 driven techniques to overcome the limitations of both approaches. Another commonly used physics-based machine learning
153 method is the physical informed neural network (PINN)²⁷. Although, the principle idea behind both methods is similar the
154 implementations differ.

155 The NI-RB method is a modification of the projection-based model order technique, from now on referred to as the intrusive
156 reduced basis (RB) method. Both methods aim to significantly reduce the degrees of freedom while maintaining the input-output
157 relationship. The intrusive version has been extensively studied in the field of mathematics^{9,10,26} and also for synthetic and
158 real-case geophysical applications^{30,31}. Furthermore, an goal-orientated error estimator for THM simulations is available³⁴.
159 However, this estimator is derived considering the conductive part of the heat transport and not in addition the advective part, as
160 required here. Furthermore, having an estimator that focuses on a quantity of interest alone is not always desirable. Especially,
161 when the prime interest of the study is to investigate the driving forces for the entirety of the model. Therefore, we use here an
162 approach that is generally applicable. The following presentation focuses on the non-intrusive RB method and the differences
163 to the intrusive method. We will not discuss the intrusive RB method in detail. For details regarding this method we refer to
164 Hesthaven et al.⁹, Benner et al.¹⁰, and Quarteroni and Rozza²⁶.

For the non-intrusive RB method, we take advantage of the circumstance that we can express the reduced solution as¹⁴:

$$\mathbf{u}_{rb}(\mu) = \sum_{i=1}^r \theta_{rb}^{(i)}(\mu) \psi_i \quad \in V_{rb}, \quad (5)$$

165 where, \mathbf{u}_{rb} is the reduced solution, r the size of the reduced basis, θ_{rb} the reduced coefficients (later on referred to as
166 weights), and ψ_i are the basis functions that span the reduced space V_{rb} .

The non-intrusive RB method consists of two stages: the offline and online stage. During the offline stage, the surrogate
model is constructed, involving all computational expensive steps. Note that this stage needs to be performed only once. In
contrast, during the online stage, the surrogate model is used, enabling a fast and efficient computation. The offline stage
consists of two steps. In the first step, the basis functions for the surrogate model are selected. This is commonly done
via the proper orthogonal decomposition¹⁴. Here, we perform a singular value decomposition and truncate the eigenvectors
after reaching a user-defined tolerance ε . This tolerance is set to $1 \cdot 10^{-4}$ in our case to ensure that the errors induced by the

Table 1. Hyperparameters of the neural network for the case study of Groß Schönebeck. Note that hl denotes hidden layers

Hyperparameter	Value
Number of hidden layers	5
Number of neurons per hidden layer	15 (hl 1), 25 (hl 2), 30 (hl 3), 35 (hl 4), 15 (hl 5)
Number of epochs	40000
Learning rate	1e-3
Batch size	25
Loss Function	sigmoid
Optimizer	Adam

approximation are negligible. The method aims to retrieve an optimal low-rank approximation. For the evaluation of the error, we are using an “energy” term according to the following definition^{10,20}:

$$\frac{\sum_{i=1}^r \sigma_i^2}{\sum_{i=1}^N \sigma_i^2} \leq \varepsilon. \quad (6)$$

167 Here, σ is the eigenvalue, and N is the total number of training samples.

168 The second step is the projection step. For the conventional intrusive RB method, we would employ a Galerkin projection
 169 (analogously to the finite element problem)⁹. In the case of the non-intrusive RB method, used here, the Galerkin projection is
 170 replaced by a machine learning method, such as Gaussian Process Regression or Neural Networks¹⁴. In this paper, we use a
 171 Neural Network since we solve for a nonlinear PDE, and Neural Networks are known for performing superior to Gaussian
 172 Process Regression in the case of nonlinearities²⁷. In this projection step, we determine the weight of each basis function.

173 In contrast to PINNs, the non-intrusive RB method first limits the physical plausible range and then uses the neural network
 174 to determine the weight of the basis functions, whereas PINNs use the physics as a constraint inside the loss function of the
 175 neural network²⁷. PINNs are classically employed for state estimation problems, whereas we need to construct a surrogate
 176 model for a parameter estimation problem. Hence, we need to define the mapping from the input parameter to the output
 177 solution, making the non-intrusive RB method advantageous.

178 An important aspect of the reliable construction of the surrogate model is the choice of the training set. This set should
 179 be representative of the entire parameter range but as small as possible since each sample in the training set corresponds to a
 180 costly finite element simulation. In this paper, we use the Latin hypercube sampling strategy³⁵ to efficiently sample the given
 181 parameter space and produce a training set of 150 samples. The validation data set is computed separately and not extracted
 182 from the training set. It consists of 20 samples, generated with a random sampling strategy, where we ensure that the validation
 183 samples are not identical to any training samples.

184 Another important point is the preprocessing of the data. Note that for the pressure output we focus on the pressure
 185 differences, thus no additional scaling is required. The input parameters (i.e., the rock properties) are scaled with a normal
 186 score transformation.

187 The offline stage of the non-intrusive RB method is computationally expensive for two reasons. The first reason is the
 188 construction of the training set, which involves solving numerous expensive finite element simulations. We already explained
 189 above how intelligent sampling strategies can reduce the cost of this stage. Another reason why the offline stage is expensive
 190 is the determination of the hyperparameters of the neural network prediction step (e.g., number of layers, neurons, epochs).
 191 Especially for complex studies, this can be a time-consuming process. To reduce the cost, a preprocessing of the input (see
 192 explanation above) is essential. Additionally, the cost can be further reduced by using a Bayes optimization scheme³⁶. For
 193 the given example the preprocessing and the Bayes optimization yield a satisfactory surrogate model. Note that for studies
 194 with a higher complexity, a Bayes optimization alone might not suffice to avoid unreasonable high offline times. In these
 195 cases, we advise using Bayes optimization with hyperband (BOHB) as the optimization method for the determination of the
 196 hyperparameters. We will not discuss the method in detail here. The general idea is to accelerate the convergence for the tuning
 197 of the hyperparameters and to enable parallel computations yielding a scalable approach. For further details, we refer to Falkner
 198 et al.³⁷. A list of the hyperparameters used in this study is presented in Table 1.

199 Global Sensitivity Analysis

200 In this paper, we perform a sensitivity analysis (SA) prior to the uncertainty quantification (UQ). The sensitivity analysis aims to
 201 determine which model parameters (e.g., rock properties such as permeability and porosity) influence the model response (e.g.,
 202 pressure) the most. This step is crucial to eliminate non-influential model parameters prior to the uncertainty quantification to
 203 avoid stability issues of the uncertainty quantification and to reduce the cost of the UQ by reducing the number of required
 204 samples³¹.

205 Sensitivity analyses can be distinguished into two categories: local and global sensitivity analyses. Here, we employ a
206 global SA. A detailed comparison between local and global SAs is found in Wainwright et al.³⁸ for hydrological models and in
207 Degen et al.³¹ for geothermal applications.

208 We use a global sensitivity analysis to investigate the entirety of the parameter space. Another reason for using the global
209 SA is that we want to consider not only the influence of the parameters themselves but also their correlations. As a global SA
210 method, we use the variance-based Sobol sensitivity analysis with a Satelli sampler. The sensitivities are expressed as the
211 ratio between the partial and total variance. To give an example, the first-order index describes the influence of the parameters
212 themselves and is mathematically the ratio of the variance of the parameter and the total variance. The total-order index
213 captures all parameter correlations in addition. Second-order indices define the correlations between two parameters, whereas
214 higher-order indices express the correlation between multiple parameters. For detailed information regarding the Sobol method,
215 we refer to Sobol³⁹ and for detailed information regarding the sampling method, we refer to Saltelli⁴⁰, and Saltelli et al.⁴¹.

216 In this work, we use the Python library SALib⁴² to perform the global sensitivity analysis. To reduce statistical errors we
217 use 100,000 realizations per parameter resulting in a total number of required forward solves of 140,000. The entire execution
218 of the global SA requires about 11 s for the non-intrusive RB approach. Note that we vary in total six parameters, including the
219 permeabilities of the Elbe base sandstone I, and the Volcanics layer. The study uses an anisotropic permeability with different
220 values of the permeability for the three main directions. We vary the dominant permeability of those two layers and keep the
221 ratio to the other directions constant. So, the two other permeability values per layer vary dependent on the main value. For the
222 permeabilities of the remaining layers, we keep the ratio to the permeability of the Elbe base sandstone layer fixed.

223 Uncertainty Quantification

224 In this work, we perform an uncertainty quantification for the model parameters influencing the pressure response in the
225 monitoring well E GrSk 3_90 (for the reservoir model of Groß Schönebeck). The uncertainty quantification is both important
226 to better determine the mean of the hydraulic parameters and their associated uncertainties. For the forward model, we consider
227 thermal, hydraulic, and mechanical parameters. However, the prior global sensitivity analysis showed that the pressure response
228 in the monitoring well is only sensitive to the hydraulic parameters.

229 As measurement data, we use the pressure data from the monitoring well E GrSk 3_90. This data has been conducted
230 during a cyclic hydraulic treatment in August 2007²¹, and we use the data as presented in Jacquey et al.²². During the treatment,
231 the pressure has been measured for about 4.5 days with a constant time step of 10 s, yielding 39,129 data points.

The UQ method is based on Bayes Theorem⁴³:

$$P(u|y) \propto P(y|u) P(u). \quad (7)$$

232 Here, $P(u|y)$ is the posterior, defining the knowledge about the unknown variable u given the data y , $P(y|u)$ is the likelihood,
233 and $P(u)$ is the prior (i.e., the knowledge of the variable u without any information about the data y)⁴³. For the sampling from
234 the posterior, we use the Markov chain Monte Carlo (MCMC) method as implemented in the Python library pyMC⁴⁴. We use
235 the MCMC method since it is a standard tool of comparison. The number of required samples could be reduced by using more
236 efficient sampling methods such as Hamiltonian Monte Carlo⁴⁵ but this does not change the general problem that this study is
237 computationally demanding.

238 Analogously, to the non-intrusive RB method, we have to set a couple of hyperparameters. We set the number of total
239 forward evaluations for the three influencing model parameters to 300,000. In addition, we use a thinning of 100 and 10,000
240 burn-in simulations. The parameters for the hydraulic properties during the MCMC run are listed in Table S1.

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329 **Author contributions statement**

330 Denise Degen has been responsible for the conceptualization, the methodology, the software, and for writing the original draft.
331 Mauro Cacace and Florian Wellmann were responsible for the reviewing of the concepts, and the reviewing and editing of the
332 manuscript.

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