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Abstract

Hydrogen-bonded organic frameworks (HOFs) are versatile materials with potential applications in proton conduction. Traditional approaches involve incorporating humidity control to address grain boundary challenges for proton conduction. This study finds vitrification as an alternative strategy to eliminate grain boundary effect in HOFs by rapidly melt quenching the kinetically stable **HOF-SXU-8** to glassy state **HOF-SXU-8-g**. Notably, a remarkable enhancement in proton conductivity without humidity was achieved after vitrification, from $1.31 \times 10^{-7} \text{ S cm}^{-1}$ to $5.62 \times 10^{-2} \text{ S cm}^{-1}$ at 100°C . Long term stability test showed negligible performance degradation, and even at 30°C , the proton conductivity remained at high level of $1.2 \times 10^{-2} \text{ S cm}^{-1}$. Molecule dynamics (MD) simulations and X-ray total scattering experiments reveal the **HOF-SXU-8-g** system is consisted of three kinds of clusters, i.e., 1,5-Naphthalenedisulfonic acid (1,5-NSA) anion clusters, N,N-dimethylformamide molecule clusters, and H^+ - H_2O clusters. In which, the H^+ plays an important role to bridge these clusters and the high conductivity is mainly related to the H^+ on H_3O^+ . ^1H magic angle spinning (MAS) solid-state nuclear magnetic resonance (NMR) revealed that the H^+ mobility was increased significantly after the transformation. These findings provide valuable insights for optimizing HOFs, enabling efficient proton conduction, and advancing energy conversion and storage devices.

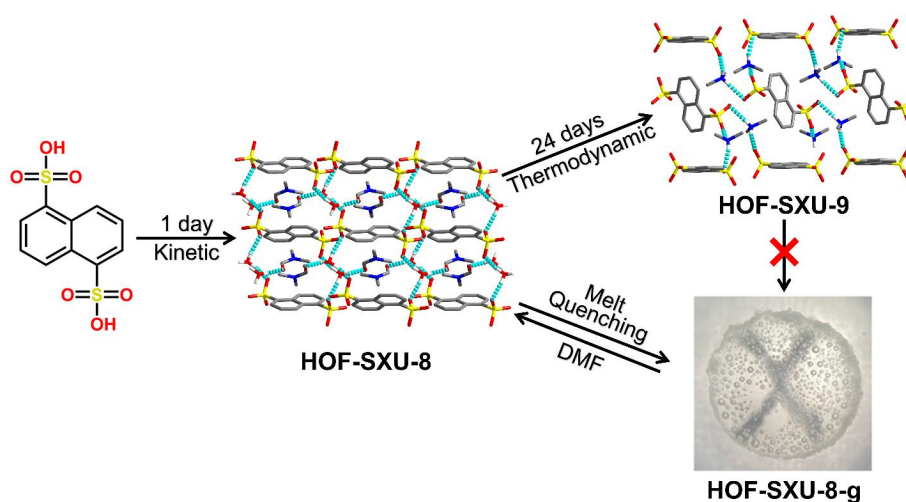
26 **Introduction**

27 In recent years, hydrogen-bonded organic frameworks (HOFs) have attracted much attention as
28 novel multifunctional materials due to their advantages such as easy synthesis, precise single-
29 crystal structure, and low energy consumption in the regeneration process¹⁻⁵. These frameworks
30 offer potential applications in various fields, including catalysis⁶⁻⁸, sensors⁹⁻¹³ and gas
31 adsorption¹⁴⁻¹⁸. Specifically, HOFs are believed to hold great potential in proton conduction due
32 to their rich inherent hydrogen bonding networks. However, a major challenge in harnessing the
33 full potential of HOFs as efficient proton conductors lies in the presence of inter-particle
34 boundaries, which hinders the desired proton transport.

35 Conventionally, most studies have relied on the incorporation of guest molecules, such as water,
36 to enhance the grain contacts and achieve high proton conductivity in HOFs¹⁹⁻²⁶, which is also
37 common in other crystalline materials such as MOFs²⁷⁻³³, COFs³⁴⁻⁴¹. The guest molecules serve
38 as mediators that facilitate proton transfer between the HOF particles. While this approach has
39 shown promising results, it introduces practical limitations, particularly at elevated temperatures
40 where managing humidity becomes cost-intensive and logistically challenging in fuel cell
41 applications. With references to the elimination of grain boundaries in coordination polymers⁴²⁻⁵¹,
42 vitrification is considered as one feasible method. Mason group⁵² has done pioneering work and
43 demonstrated that a desymmetrization strategy can be used to realize glass transitions of an
44 alkylguanidinium sulfonate HOF. However, it is presumably that the entire HOF precursor is in a
45 low energy state and the glassy material obtained does not exhibit significant proton conduction
46 properties.

47 In this work, two single crystals, kinetically stable **HOF-SXU-8** with higher energy and
48 thermodynamically stable **HOF-SXU-9** with lower energy were synthesized through liquid-phase
49 diffusion method. **HOF-SXU-8** could undergo a vitrification transformation to glassy state **HOF-**
50 **SXU-8-g** at around 110°C (Scheme 1). Following this vitrification transformation, the proton

51 conductivity of **HOF-SXU-8-g** without humidity at 100°C reached as high as $5.62 \times 10^{-2} \text{ S cm}^{-1}$,
52 while the conductivity for **HOF-SXU-8** is only $1.31 \times 10^{-7} \text{ S cm}^{-1}$. Remarkably, even at 30°C, the
53 proton conductivity of **HOF-SXU-8-g** maintained a quite high level of $1.2 \times 10^{-2} \text{ S cm}^{-1}$.
54 Furthermore, after a stability test of 16 hours, the performance of **HOF-SXU-8-g** remained
55 almost unchanged. On the other hand, due to its thermodynamic stability and limited free
56 volume, **HOF-SXU-9**, does not undergo such transition. ^1H MAS solid-state NMR revealed that
57 the H^+ mobility was increased significantly after the transformation.



58
59 **Scheme 1:** Schematic representation of the synthesis of **HOF-SXU-8**, **HOF-SXU-9** and the vitrification transformation
60 process to **HOF-SXU-8-g**. (Color code: C atom, gray; O atom, red; S atom, yellow; N atom, blue; H atom, white).

61 Results

62 Crystal structure descriptions of **HOF-SXU-8** and **HOF-SXU-9**

63 Single crystal X-ray analysis reveals that **HOF-SXU-8** crystallizes in the space group $P2_1/n$
64 (please see Table S1 for more details), with 1,5-NSA anion layers separated by DMF molecules.
65 H_3O^+ connects a continuous hydrogen bonding network capable of proton conduction by linking
66 different 1,5-NSA anion layers and the DMF molecules (Fig. S1). Specifically, O(A) on H_3O^+
67 forms one hydrogen bond with O(1A) from the upper layer of 1,5-NSA anion ($\text{O}-\text{H}\cdots\text{O}$, 2.65 Å,

68 166.85°) and another hydrogen bond with O(2A) from the lower layer of 1,5-NSA anion (O-H···O,
69 2.65 Å, 152.37°), and one hydrogen bond with O(3A) on DMF (O-H···O, 2.44 Å, 161.86°) as
70 depicted in Table S3. After soaking **HOF-SXU-8** in the mother solution for 4 days, the **HOF-**
71 **SXU-8** crystals disappeared. Another 20 days later, thermodynamically stable colorless single
72 crystals of **HOF-SXU-9** appeared in the space group $P2_1/c$ (please see Table S2 for more
73 details). Dimethylamine cations in **HOF-SXU-9** link the differently oriented 1,5-NSA anion
74 through two hydrogen bonds (Fig. S2 and Table S4). The purity of bulk crystalline **HOF-SXU-8**
75 and **HOF-SXU-9** were confirmed through good accordance with the powder X-ray diffraction
76 (PXRD) pattern simulated from the crystallographic information files, as shown in Fig. S3.
77 Thermogravimetric (TG) analysis (Fig. S4) revealed that the kinetically stable **HOF-SXU-8** starts
78 to lose weight at 100.26°C with a relatively low thermal stability. In contrast, the
79 thermodynamically stable **HOF-SXU-9** can keep its structure intact to about 310°C. Above
80 310°C, 1,5-NSA anion starts to decompose, leading to the collapse of its skeleton. Overall,
81 thermodynamically stable **HOF-SXU-9** exhibits higher thermal stability.

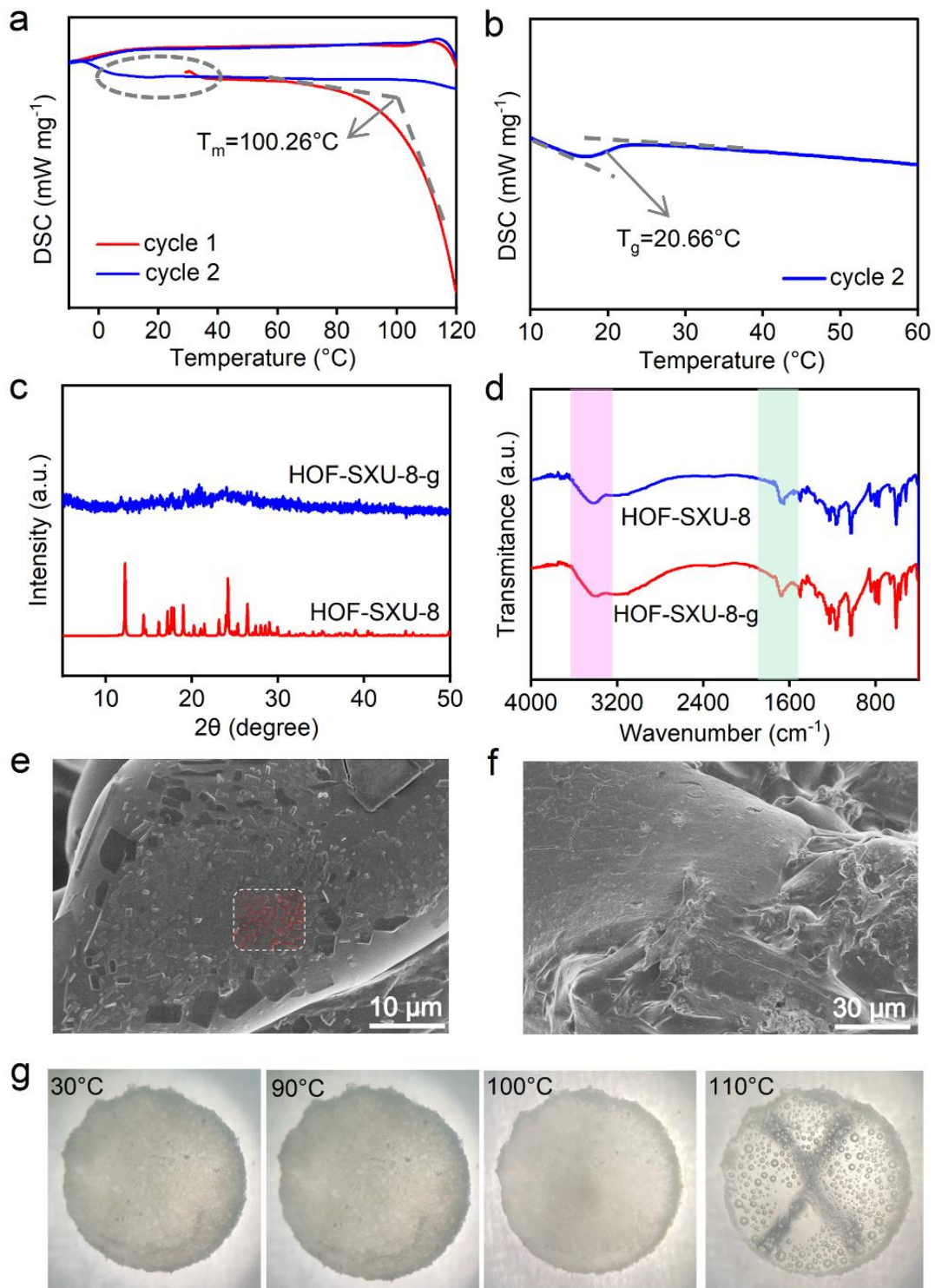
82 Crystal melting and glass formation

83 **HOF-SXU-8** was heated at 110°C for 2h and then quenched rapidly to obtain **HOF-SXU-8-g** in
84 glassy form (details please see supporting information section 2). To confirm the occurrence of
85 vitrification, we used differential scanning calorimeter (DSC) to monitor the enthalpy response of
86 the phase transition process (Fig. 1a). When the first heating cycle was applied, a heat
87 absorption peak appeared at 100.26°C, the appearance of which was attributed to the phase
88 transition of solid-liquid transformation occurring in the crystal. At the second heating cycle, the
89 heat absorption peak disappeared, confirming the occurrence of vitrification, and the glass
90 transition temperature T_g was calculated to be 20.66°C (Fig. 1b). PXRD (Fig. 1c) revealed that
91 the characteristic peaks of crystalline **HOF-SXU-8** disappeared after melt quenching to **HOF-**

92 **SXU-8-g** accompanying by a new broad peak appeared, which is typical of glassy phase⁴⁷⁻⁵⁰.
93 PXRD analyses conducted on the **HOF-SXU-9** sample after heated at 110°C for a duration of 2
94 hours have revealed that the distinctive peaks were still remained (Fig. S5), suggesting that the
95 transformation of **HOF-SXU-9** into an amorphous or glassy state did not occur.

96 Scanning electron microscope (SEM) image clearly shows the presence of grain boundaries in
97 the compressed **HOF-SXU-8** pellet. The morphology of the HOF crystals varies, with well-
98 defined edges and grain boundaries clearly seen in Fig. 1e. While for **HOF-SXU-8-g**, the grain
99 boundaries disappear and the whole material exists in homogeneous form (Fig. 1f). Optical
100 photography is an effective way to observe changes in samples during heat treatment, thus
101 allowing observation of processes such as sintering and melting. **HOF-SXU-8** pellet exhibited
102 different degrees of transparency after processed at different temperatures^{47, 53}. In Fig. 1g,
103 where the sample is placed on the upper side of a printed "X". The "X" is only visible after
104 treated the pellet at 110°C for 2 hours and then cooled. Combining the results of the
105 thermogravimetric analyses, it is clear that the change in transparency is due to the occurrence
106 of vitrification.

107 In Fig. 1d, the C=O (1680 cm⁻¹) stretching vibration comes from the DMF structural unit and the
108 O-H (3413 cm⁻¹) stretching vibration comes from the H₃O⁺ structural unit, which is bound to the
109 O-H bond after the formation of hydrogen bonding and the spectral peak broadens. These
110 characteristic peaks are present in the Fourier transformed infrared (FTIR) spectra of both **HOF-**
111 **SXU-8** and **HOF-SXU-8-g**, proving that the DMF structural unit and the hydrogen bonding
112 network are still present in the structure after the vitrification process (Fig. 1d). Elemental
113 analysis (Table S5) showed that the molecular formulas were (1,5-NSA)_{1.04} (H₃O)_{2.25} (DMF)_{2.2}
114 for **HOF-SXU-8** and (1,5-NSA) (H₃O)_{2.26} (DMF)_{1.56} for **HOF-SXU-8-g**, with the difference related
115 to the loss of approximately 30% DMF after glass formation. **HOF-SXU-8-g** was observed to re-
116 crystallize (Fig. S6) back to **HOF-SXU-8** under heat treatment in DMF atmosphere, which is



117
 118 **Fig. 1:** **a** DSC scans of crystalline **HOF-SXU-8**; **b** Enlarged Figure of the circled part in Fig. 1a; **c** PXRD patterns of
 119 **HOF-SXU-8** and **HOF-SXU-8-g**; **d** FTIR spectra of **HOF-SXU-8** and **HOF-SXU-8-g**; **e** SEM image of the **HOF-SXU-8**
 120 pellet; **f** SEM image of the **HOF-SXU-8-g**; **g** Optical photographs of crystalline **HOF-SXU-8** after different temperature
 121 treatments. The sample was put on the mark "X" to assess the variation in transparency.

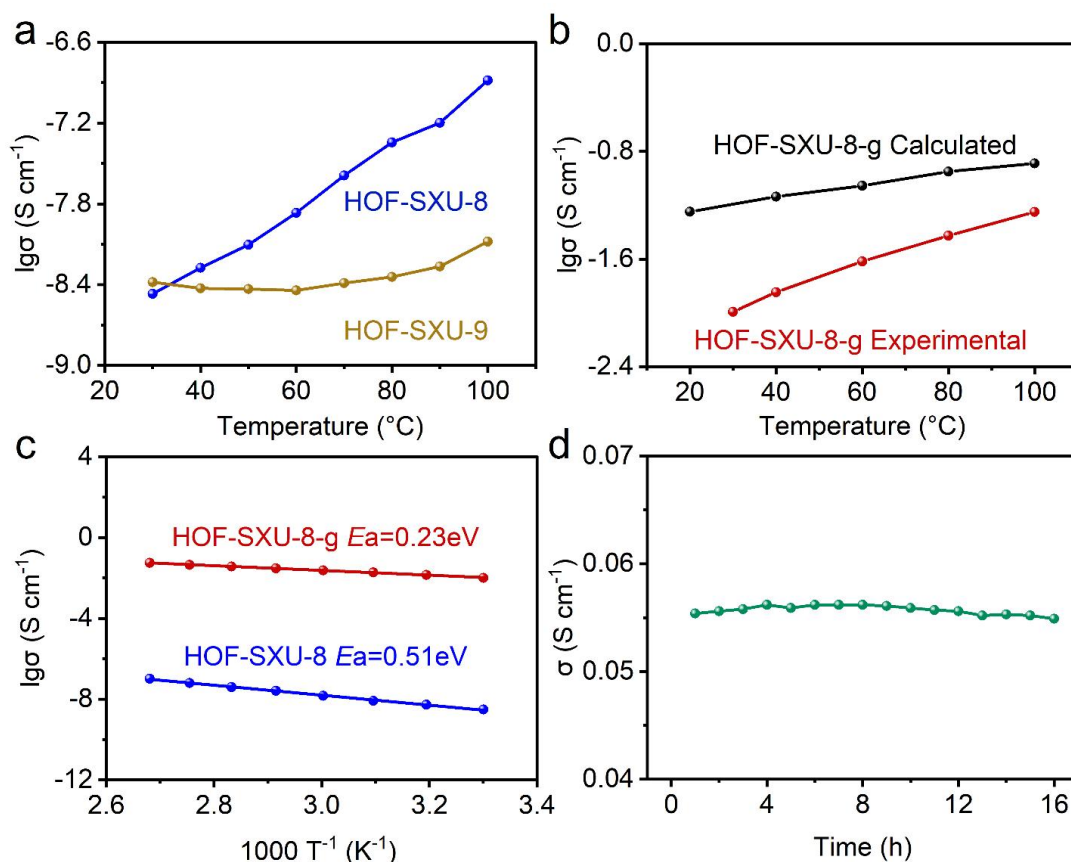
122 proved by the PXRD results (Fig. S7). However, no such re-crystallization was observed in the
123 absence of DMF atmosphere, regardless of the same heat treatment temperature. This proves
124 that the glass transformation is attributed to the partial release of structural DMF molecules,
125 which is different from the complete guest molecules removal as previously reported by Mason's
126 group⁵².

127 Proton conduction properties

128 Electrochemical impedance spectroscopy (EIS) of **HOF-SXU-8**, **HOF-SXU-9** and **HOF-SXU-8-g**
129 were measured without humidity at different temperatures. The conductivity values of **HOF-**
130 **SXU-8** exhibited an increase from $3.40 \times 10^{-9} \text{ S cm}^{-1}$ at 30°C to $1.31 \times 10^{-7} \text{ S cm}^{-1}$ at 100°C (Fig.
131 2a and Fig. S8) with the crystalline phase of **HOF-SXU-8** remained throughout the
132 measurements (Fig. S11). Crystalline structure of **HOF-SXU-9** was also maintained (Fig. S11)
133 and the conductivity slightly increased from $4.15 \times 10^{-9} \text{ S cm}^{-1}$ at 30°C to $8.31 \times 10^{-9} \text{ S cm}^{-1}$ at
134 100°C (Fig. S9), which is ranked at a very low level^{24, 54}. This phenomenon may be related to
135 the strong hydrogen bonding network in the thermodynamically stable **HOF-SXU-9**.

136 Conductivity measurements of **HOF-SXU-8-g** revealed a high value of $1.02 \times 10^{-2} \text{ S cm}^{-1}$ at
137 30°C , and then increased slowly to $5.62 \times 10^{-2} \text{ S cm}^{-1}$ at 100°C (Fig. 2b and Fig. S10). The
138 proton conductivity of **HOF-SXU-8-g** has been improved by about 5 orders of magnitude
139 compared to single-crystalline **HOF-SXU-8**, making **HOF-SXU-8-g** a superior proton conductive
140 material^{20, 22, 23, 25, 55-57}. The Arrhenius plot (Fig. 2c) demonstrates different mechanisms of **HOF-**
141 **SXU-8** and **HOF-SXU-8-g**: the activation energy of **HOF-SXU-8** was calculated to be 0.51 eV
142 for the vehicular mechanism, while the activation energy value of **HOF-SXU-8-g** was 0.23 eV,
143 which is typical for the Grotthuss mechanism⁵⁸. In addition, **HOF-SXU-8-g** can maintain its
144 superior proton conductivity up to 16 h at 100°C (Fig. 2d and Fig. S12) with performance kept

145 almost the same, revealing that the glassy **HOF-SXU-8-g** structure is well retained in this
 146 process.



147
 148 **Fig. 2:** a Proton conductivity of **HOF-SXU-8** and **HOF-SXU-9** with increasing temperature; b Experimental (red) and
 149 calculated (black) proton conductivity of **HOF-SXU-8-g** with increasing temperature; c Arrhenius plot of **HOF-SXU-8**
 150 and **HOF-SXU-8-g**; d Time-dependent proton conductivity of **HOF-SXU-8-g** performed at 100°C .

151 Simulation and Proton conduction mechanism

152 In order to probe the proton conduction mechanism, we performed MD simulations using
 153 LAMMPS 2022 software (details please refer to supporting information). An energy minimization
 154 procedure and an NVT ensemble simulation ($T = 298 \text{ K}$, $p = 1 \text{ bar}$, time = 1 ns) were employed
 155 to obtain equilibrated conformation (Fig. 3a). Further NVT ensemble simulation (15 ns, $T = 298$

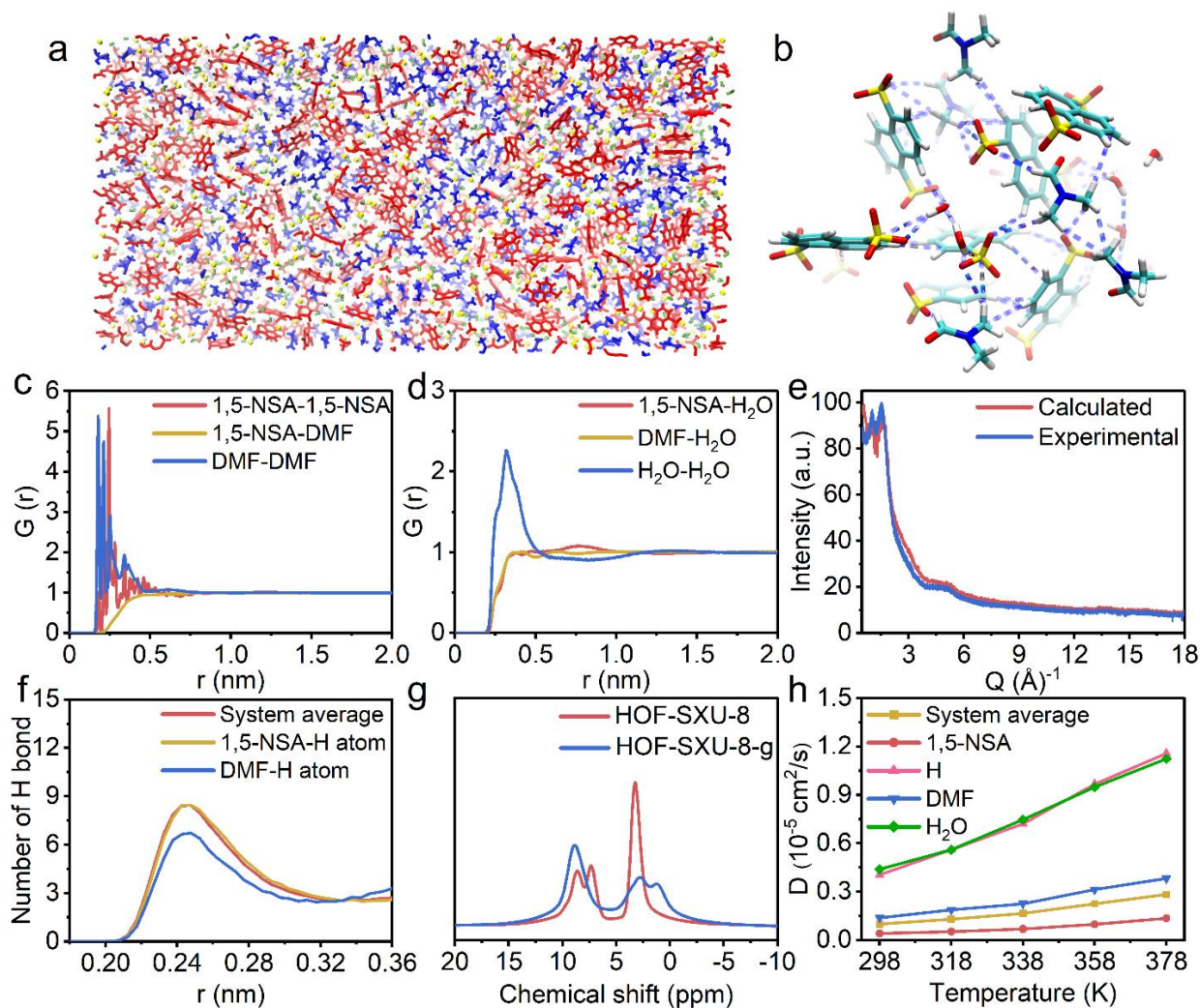
156 K) and annealing procedure (5 ns, 298 K-873 K) were performed to assure that the above
157 obtained system is stable.

158 The radial distribution functions (RDF) of 1,5-NSA anion-1,5-NSA anion, DMF-DMF have similar
159 multiple peaks within 0.6 nm (Fig. 3c), suggesting that both 1,5-NSA anion and DMF
160 spontaneous aggregated to clusters with a radius of 0.5-0.6 nm and exhibited a certain spatial
161 distribution. The RDF of 1,5-NSA anion-DMF is low within 0.6 nm, indicating that the two
162 clusters are mutually exclusive. Beyond 0.8 nm, it gradually converges to 1, indicating that the
163 clusters are more uniformly dispersed with each other. Besides, there are significant interactions
164 between the H₂O molecules to form clusters with radius of 0.4-0.5 nm, while no significant
165 interactions between H₂O and 1,5-NSA anion, or between H₂O and DMF (Fig. 3d). Beyond 1.2
166 nm, the RDFs associated with H₂O converge to 1, indicating that the H₂O clusters are dispersed
167 with other components of **HOF-SXU-8-g** (Fig. 3d). Similarly, the RDF between H⁺ and H₂O (Fig.
168 S14) is exceeded near 0.36 nm, suggesting that an indirect association is formed between H⁺
169 relying on the H₂O molecule, while no significant H⁺-1,5-NSA anion and H⁺-DMF interaction
170 were observed. In summary, the system are summarized by the following features: (1) the
171 substances are divided into three kinds of clusters, i.e., 1,5-NSA anion clusters, DMF molecule
172 clusters, and H⁺-H₂O clusters, with the cluster radius of 0.4-0.6 nm; (2) the clusters were
173 diffused with each other over medium and long distances (>0.8-1.2 nm). We also obtained the
174 X-ray total scattering pattern of **HOF-SXU-8-g** and compared it with the simulated data based
175 on the MD model (Fig. 3e). The experimental X-ray total scattering pattern of **HOF-SXU-8-g**
176 agrees well with that of the simulated results, indicating that the structure of **HOF-SXU-8-g** is
177 consistent with simulated **HOF-SXU-8-g**, further confirming the accuracy of our simulated **HOF-**
178 **SXU-8-g** structure as discussed above.

179 The final equilibrated structure was then subjected to 1 ns equilibrium sampling, with the
180 structure saved every 5 ps to form a trajectory and the interaction energy saved every 1 ps. The

181 electrostatic interaction energies and long-range non-bonding interaction energies between 1,5-
182 NSA anion, DMF molecules, H^+ , and H_2O were decomposed, and Fig. S15 demonstrates the
183 interaction energies of each group as a percentage of the total interaction energy. 1,5-NSA
184 anion-1,5-NSA anion, 1,5-NSA anion- H^+ , DMF-DMF, H^+ - H^+ , and H^+ - H_2O interactions occupied
185 the top four proportion. 1,5-NSA anion- H_2O had a lower interaction energy (4%), whereas 1,5-
186 NSA anion- H^+ (19.7%) and H^+ - H_2O (8.8%) had higher interaction energies, so H^+ acted as a
187 bridge connecting 1,5-NSA anion clusters to H_2O clusters, leading to the two types of clusters
188 mixed and dispersed. Identically, 1,5-NSA anion-DMF has a lower interaction energy (0.9%),
189 while the interaction energies of 1,5-NSA anion- H^+ and H^+ -DMF account for 19.7% and 3.8%,
190 respectively, thus H^+ also serves as a bridge between 1,5-NSA clusters and DMF clusters. In
191 short, among the three original constituents of the **HOF-SXU-8-g** system, H^+ acts as a bridge
192 between the clusters of the whole system.

193 The snapshot of the local structure in the middle period of simulation was drawn in Fig. 3b to
194 display the rich hydrogen bonding. Hydrogen bonding analysis was carried out on the simulated
195 trajectories to elucidate the contribution of different components. The number of hydrogen
196 bonds generated by 1,5-NSA anion, DMF molecules, and the system averaged over different
197 distances was analyzed (Fig. 3f). The hydrogen bonds formed by 1,5-NSA anion and DMF
198 molecules in the system were distributed mainly at the distance range of 0.22-0.28 nm, and the
199 number of hydrogen bonds formed by 1,5-NSA anion was significantly higher than that of DMF
200 molecules. Thus, for this system, more hydrogen bonds can be formed in the same volume due
201 to the release of DMF, which strengthens the hydrogen bonding network in the system and
202 forms a glassy structure, and in turn would greatly enhance the proton conductivity of the glassy
203 **HOF-SXU-8-g**.



204
 205 **Fig. 3:** **a** Structure diagram of simulated **HOF-SXU-8-g** (Cluster Color: 1,5-NSA anion, red; DMF molecule, blue; H⁺,
 206 yellow; H₂O molecule, green); **b** Transient local structure of simulated **HOF-SXU-8-g** in MD simulation process (Color
 207 code: C atom, pale blue; O atom, red; S atom, yellow; N atom, blue; H atom, white); **c** RDFs associated with 1,5-NSA
 208 anion and DMF molecules; **d** RDFs associated with H₂O molecules; **e** X-ray total scattering pattern of **HOF-SXU-8-g**
 209 at room temperature (blue) compared with that calculated from MD simulations (red); **f** Distribution of the number of
 210 hydrogen bonds over distance; **g** ¹H MAS solid-state NMR spectra of **HOF-SXU-8** and **HOF-SXU-8-g** at room
 211 temperature; **h** Self-diffusion coefficients of different components from **HOF-SXU-8-g** at different temperatures.

212 The self-diffusion coefficients of different components in the simulated **HOF-SXU-8-g** increased
 213 with the increasing of temperature (Fig. 3h). The self-diffusion coefficients of H⁺ and H₂O are the
 214 largest two among all species and are always closer at different temperatures, which also
 215 confirms the existence of H⁺-H₂O clusters. The calculated conductivity values of the simulated

216 **HOF-SXU-8-g** system at different temperatures are shown in Fig. 2b. At 308 K and without an
217 applied electric field, the conductivity is $6.41 \times 10^{-2} \text{ S cm}^{-1}$, which is quite close to the
218 experimentally obtained conductivity of $1.20 \times 10^{-2} \text{ S cm}^{-1}$. Due to the huge difference in the
219 self-diffusion coefficients between H^+ and 1,5-NSA anion, the main contribution to the electrical
220 conductivity (90.7%) of the system comes from the diffusion of hydrogen ions (H^+ - H_2O clusters)
221 (Fig. 3h). Since the conductivity is mainly related to H^+ on H_3O^+ , ^1H MAS solid-state NMR was
222 utilized to study H^+ mobility (Fig. 3g). The peaks with chemical shifts of 5-10 ppm are attributed
223 to H^+ on H_3O^+ . The peak of H^+ on H_3O^+ of **HOF-SXU-8** is cleaved because H_3O^+ forms hydrogen
224 bonds with 1,5-NSA between the layers and with the DMF molecule. While H^+ on H_3O^+ of **HOF-**
225 **SXU-8-g** only exhibits one averaged single peak at 5-10 ppm because of its isotropic property
226 after vitrification. It could also be deduced that after vitrification, the ease of movement of H^+
227 increases due to the peak intensity significantly increased (Fig. 3g), which further leads to the
228 substantial increase in the proton conductivity.

229 **Discussion**

230 In summary, we have successfully prepared two hydrogen bonded organic frameworks:
231 kinetically stable **HOF-SXU-8** and thermodynamically stable **HOF-SXU-9**, through controlling
232 the reaction time. **HOF-SXU-8** could be transformed by rapidly melt quenching method at about
233 110°C to glassy state **HOF-SXU-8-g** triggered by partial loss of DMF molecules. **HOF-SXU-8-g**
234 exhibits superior anhydrous proton conduction performance from 30°C to 100°C , with
235 conductivity reaching as high as $5.62 \times 10^{-2} \text{ S cm}^{-1}$ and excellent stability. MD simulations and
236 X-ray total scattering experiments reveal the **HOF-SXU-8-g** system can be divided into three
237 kinds of clusters, i.e., 1,5-NSA anion, DMF molecule, and H^+ - H_2O clusters. In which, the H^+
238 plays important role to bridge these clusters and the conductivity is mainly related to H^+ on H_3O^+ .
239 ^1H MAS solid-state NMR revealed that the H^+ mobility was increased significantly after the glass
240 transformation. We believe that the glassy HOF will be a design well worth exploring in depth

241 and may open new avenues for obtaining safe and high-performance pure solid-state
242 electrolytes.

243 **Methods**

244 **Synthesis**

245 All solvents and materials were purchased commercially and used without further purification.
246 1,5-NSA was purchased from Adamas-Beta. DMF and n-hexane were purchased from Tianjin
247 Damao Chemical Reagent Factory.

248 Synthesis of **HOF-SXU-8**: 1,5-NSA (2.08 mmol, 600 mg) was dissolved in DMF (5 mL) in a 20
249 mL glass vial by ultrasound at 25°C. 1 mL of the above solution was then placed in a 20 mL
250 glass vial and 2 mL hexane was added slowly. After 1 day, transparent **HOF-SXU-8** crystals
251 were precipitated in the bottom of the glass vial. Crystallographic data for **HOF-SXU-8** has been
252 deposited with the Cambridge Crystallographic Data Centre, CCDC, depository number
253 2292290.

254 Synthesis of **HOF-SXU-9**: The above obtained **HOF-SXU-8** was observed to disappear after 4
255 days of its appearance. Another 20 days later, thermodynamically stable crystals of **HOF-SXU-9**
256 appeared in the bottom of the glass vial. Crystallographic data for the **HOF-SXU-9** has been
257 deposited with the Cambridge Crystallographic Data Centre, CCDC, depository number
258 2292291.

259 Synthesis of **HOF-SXU-8-g**: **HOF-SXU-8-g** were obtained by rapid melt quenching of **HOF-**
260 **SXU-8**. **HOF-SXU-8** crystals were ground to powder first and 15 mg of the powder was pressed
261 into a pellet. The pressed **HOF-SXU-8** pellet was further heated in an oven at 110°C for 2h and
262 then quickly quenched to obtain **HOF-SXU-8-g**.

263 **Characterizations**

264 TG was performed on a NETZSCH STA 449C thermal analyzer under nitrogen (N₂) atmosphere
265 from 30 to 900°C at a rate of 10°C/min. T_g was determined by differential scanning calorimetry
266 (DSC) (STA 404 C, Netzsch) and collected with a Netzsch DSC 200 F3 under N₂ atmosphere (Al
267 crucible). The morphology of **HOF-SXU-8** and **HOF-SXU-8-g** were analyzed using SEM (Zeiss
268 EVO 60 S). The changes of the samples after heat treatment were observed by optical and
269 microscope. ¹H MAS Solid-state NMR was measured by the single pulse technique with a
270 relaxation time of 3.00 s at a 8 kHz spinning rate at 28°C (¹H resonance frequency of 400.33
271 MHz). Elemental analyses (C, H, S and N) were performed in-house using an Elementar vario
272 EL elemental analyzer.

273 X-ray characterization

274 The single crystal data were collected with a Rigaku XtaLab P200 diffractometer and a Dectris
275 Pilatus 200K system at 302 K. The system was equipped with A MicroMax007 HF/VariMax
276 rotating anode X-ray generator with confocal monochromatic Mo-K α radiation. The collected
277 data were solved and refined by full matrix least squares using SHELXL 2016/4. The phases of
278 the synthesized crystals were identified using powder X-ray diffraction (PXRD) method. The
279 PXRD experiments were performed on a Rigaku Ultima IV x-ray diffractometer at 40 kV and 40
280 mA with a scan rate of 1°/min and a 2 θ range of 5-50°, and the diffractometer was equipped with
281 a Cu sealed tube ($\lambda = 1.5406 \text{ \AA}$). Based on the single crystal diffraction data, simulated PXRD
282 patterns were generated using Mercury 3.9 software.

283 Proton conductivity measurements

284 Proton conductivity was characterized through a two-probe method using a Solartron SI1260
285 Impedance / Gain-phase analyzer coupled with a Solartron 1287 dielectric interface. Before
286 conductivity measurements, **HOF-SXU-8** and **HOF-SXU-9** samples were placed in a steel mold
287 and pressed into pellets of 6 mm diameter and 1.5 mm thickness using a tablet press. For the

288 measurement of **HOF-SXU-8-g**, 100 mg **HOF-SXU-8** was first pressed into a 6 mm diameter
 289 pellet (pressure 10 MPa). **HOF-SXU-8** pellet was then put into a 2032 steel cell in a glove box
 290 filled with ultra pure Ar (H_2O and $O_2 < 1$ ppm) and heated in an oven at $110^\circ C$ for 2 hours to
 291 obtain **HOF-SXU-8-g** by rapidly melt quenching. All measurements in this work were performed
 292 under a stream of ultra-pure N_2 ($>99.99\%$). The measurements were conducted in the
 293 frequency range from 0.5 Hz to 1 MHz with an input voltage amplitude value of 100 mV. The
 294 impedance values at each temperature were repeatedly measured after 30 minutes of
 295 equilibrium until the measured values keep stable. The resistance values were obtained by
 296 fitting the impedance profile using Zview software. The circuit equivalent used for fitting is as
 297 follows:



298

$$\sigma = l/RS \quad (1)$$

$$\sigma T = A \exp(-E_a/k_B T) \quad (2)$$

299 R1 corresponds to the resistances of wire and electrode, while R2 accounts for the bulk
 300 resistance of the pellet. CPE1 denotes the nonideal capacitance corresponding to the bulk. The
 301 conductivity was calculated using equation (1), where σ is the conductivity ($S\ cm^{-1}$), l is the
 302 pellet thickness (cm), S is the solid electrolyte area (cm^2), and R is the bulk resistance of the
 303 pellet (R_2 in the circuit equivalent) fitted by the equivalent circuit of the semicircle in Nyquist plot
 304 using zview software. The activation energy (E_a) of the material conductivity is estimated from
 305 equation (2), where A is the exponential prefactor, k_B is the Boltzmann constant, and T is
 306 experiment temperature.

307 X-ray total scattering measurements

308 **HOF-SXU-8** sample was packed in a borosilicate capillary with a diameter of 0.7 mm, then in-
309 situ heated at 110°C and quenched rapidly back to room temperature to obtain **HOF-SXU-8-g** in
310 the borosilicate capillary. The X-ray total scattering data was collected on the Mythen-II detector
311 over 115° covering the Q range up to 18 Å⁻¹ (21 keV; λ = 0.5921 Å). The wavelength, zero error
312 and instrument contribution to the peak profile was determined using the line position and line
313 shape standard NIST Si sample with the refined wavelength 0.5921 Å. For the experiment, an
314 empty capillary was used as the background. The collected scattering data was applied
315 absorption, background, and Compton scattering corrections then normalized to give the X-ray
316 total scattering pattern.

317 Computational method

318 Molecular dynamics simulations were employed to study the material. The simulation is based
319 on GAFF force field, and the force field parameters are obtained from quantum chemistry
320 calculation and fitting, using m2seminario method. The long-range electrostatic term in the force
321 field is calculated using Particle-Mesh-Ewald (PME) method, and the van der Waals interaction
322 is calculated using Lennard-Jones potential (cutoff = 1.0 nm).

323 The simulation system contains 1440 naphthalenedisulfonic molecules, 2016 DMF molecules
324 and 2880 water molecules, 62464 atoms in total.

325 Firstly, energy minimization procedure and an NVT ensemble simulation (T = 298 K, p = 1 bar,
326 time = 1 ns) were employed to obtain a equilibrated and compressed conformation. After these
327 processes, the system has a volume of 7.10×10⁵ Å³ and density of 1.26×10³ kg/m³.

328 Secondly, an NVT ensemble simulation (15 ns, T = 298 K) and an annealing procedure (5 ns,
329 298 K-873 K) were performed to assure that the system has been stable. The annealing
330 procedure were performed in the protocol as below: 100 ps for each cycle, heat from 298 K to

331 873 K for 30 ps, keep temperature in 873 K for 30 ps, anneal from 873 K to 298 K in 5 ps, and
332 keep temperature in 298 K for 35 ps, repeat the cycle for 50 times. After the above annealing
333 procedure, the end structure kept the same as the start structure, which indicates that the
334 system has been the stable structure already (Fig. S13).

335 A sampling for 1 ns was performed to the productive simulation, in which the timestep is 1 fs,
336 trajectories were recorded every 5 ps, energies and properties were recorded every 1 ps.

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473

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476 **Ethics declarations**

477 **Competing interests**

478 The authors declare no competing financial interests.

479 **Data Availability**

480 The X-ray crystallographic coordinates for structure reported in this article have been deposited
481 at the Cambridge Crystallographic Data Centre (CCDC), under deposition number CCDC
482 2292290 and 2292291. These data can be obtained free of charge from The Cambridge
483 Crystallographic Data Centre via www.ccdc.cam.ac.uk/data_request/cif.

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