

# Fluid Transport and Storage in the Cascadia Forearc inferred from Magnetotelluric Data

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## Article

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# Fluid Transport and Storage in the Cascadia Forearc inferred from Magnetotelluric Data

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## Abstract

Subduction of hydrated oceanic lithosphere can carry water deep into the Earth, with important consequences for a range of tectonic and magmatic processes. Most fluid is released at relatively shallow depths in the forearc where it is thought to play a critical role in controlling mechanical properties and seismic behavior of the subduction megathrust. Here we present results from three-dimensional inversion of nearly 400 long-period magnetotelluric sites, including 64 offshore, to provide new insights into the distribution of fluids in the forearc of the Cascadia subduction zone. Our amphibious dataset provides new constraints on the geometry of the electrically resistive Siletzia terrane, a thickened section of oceanic crust accreted to North America in the Eocene, and the conductive accretionary complex, which is being underthrust all along the margin. Fluids accumulate, over time-scales likely exceeding 1 My, above the plate interface in metasedimentary units, while the mafic rocks of Siletzia remain dry. Fluids in metasediments tend to peak at fixed slab-depths of 17.5 and 30 km, suggesting control by metamorphic processes, but also concentrate around the edges of Siletzia, suggesting that this mafic block is impermeable, with dehydration fluids escaping up-dip along the megathrust. Our results demonstrate that lithology of the overriding crust can play a critical role in controlling fluid transport and sequestration in a subduction zone, with potentially important implications for mechanical properties.

# 1 Introduction

Cascadia is an end-member subduction zone with a very young and hot subducting plate, capable of producing very large (Mw 9) megathrust earthquakes [1]. The central forearc crust, generically referred to as Siletzia [2], is a thickened section of oceanic crust that accreted to North America in the late Eocene (Fig. 1). The late-Tertiary Klamath [3], and Wrangellia [4] terranes, respectively form the southern and northern domains of the Cascadia forearc. Seismic refraction profiles [5] suggest that Siletzia is thickest in central-northern Oregon and thins to the north. In the Olympic peninsula the full crustal section is formed from accreted Tertiary sediments and metasediments, backstopped in the east by a thick section of Siletzia, which has been heavily deformed and tilted past vertical [6].

Cascadia is an archetype for episodic tremor and slip (ETS), characterized by bursts of non-volcanic seismic tremor accompanied by low-frequency earthquakes and geodetically detected slow slip [7] at or above the plate interface. There is substantial evidence [8–10] that ETS occurs where effective stresses are orders of magnitude less than lithostatic, which is most readily explained by overpressured fluids along or near the megathrust. Indeed, a low velocity layer (LVL) with anomalously high Poisson’s ratios, also attributable to overpressured fluids [11], has been detected within, or slightly above, the subducting oceanic crust in Cascadia [12–14].

Although ETS occurs essentially everywhere in Cascadia, there are significant along-margin variations, with reduced intensity and longer recurrence intervals beneath the thickened central core of Siletzia in north-central Oregon [15]. The same area is almost devoid of crustal and intraplate seismicity [16], and may be more weakly coupled in the seismogenic locked zone [17]. This correlation of forearc crustal geology with patterns of ETS, seismicity, and plate locking has often been ascribed to variations in fluid pressure along the megathrust, due either to lithologically-controlled variations in crustal permeability [15, 17, 18], or variable hydration of the incoming plate [19–23].

Seismic imaging provides numerous indirect constraints on fluids and fluid-mediated processes in subduction zones, allowing estimates of the abundance of hydrous minerals in the subducting

48 plate [20, 24], forearc crust [25, 26], and forearc mantle [27], fluid content and consolidation in the  
49 accretionary prism [22, 28], and volumes of over-pressured fluids near the megathrust [11, 12, 29,  
50 30]. Magnetotelluric (MT) imaging of electrical resistivity can provide more direct constraints on  
51 the present distribution and connectivity of free fluids in the crust. Previous MT studies in Cascadia  
52 have suggested along-margin variations in sediment subduction, high concentrations of fluids near  
53 the ETS zone, and evidence for flux melting beneath the arc [31–35]. Here we present new detailed  
54 images of electrical resistivity in Cascadia, based on three-dimensional (3D) inversion of 393 long-  
55 period MT sites (Fig. 1). The dataset includes 64 new offshore sites, extending from the trench  
56 to the coast, 102 new onshore sites, as well as all available long-period MT sites from legacy  
57 land profiles throughout the region [32], and 144 sites from the EarthScope MT transportable  
58 array [36]. Although the modeling domain extends from spreading ridge to back arc, and from  
59 northern California to southern British Columbia, our focus here is on the forearc in Oregon and  
60 SW Washington, where data density is highest. Details on the dataset and the inversion approach  
61 are provided in Methods and the supplemental material (SM).

## 62 **2 Results**

63 Our preferred 3D inverse solution is shown in Fig. 2, with a more complete set of model sections  
64 provided in Fig.

65 Fig. 2 reveals a conductive body (C1), extending from offshore to an abrupt landward termi-  
66 nation beneath a massive resistive body (R1). We interpret C1 to represent fluid-rich sedimentary  
67 and metasedimentary rocks of the accretionary complex, and R1 to represent comparatively dry  
68 mafic rocks of Siletzia. Indeed, the geometry of R1 is broadly consistent with geological and geo-  
69 physical constraints on Siletzia, including spatial extent [37], greater thickness in Oregon [5], and  
70 significant disruption and unroofing of the mafic section in the Olympic peninsula [6, 38] (R2).  
71 In Oregon and southern Washington Siletzia is very thick in the east, thinning abruptly to  $\sim 10$   
72 km in a layer that extends offshore. The conductive accretionary complex is thrust a significant  
73 distance ( $\sim 50$  km, more in places) under this thin layer, terminating at the nearly vertical edge of  
74 R1. Offshore, conductivity of C1 is highest from 44-47°N; however, this may reflect limited off-  
75 shore data coverage outside of this area (Fig. 1). A deeper, higher amplitude, peak in conductivity

76 (C2) occurs in all sections shown at a slab interface depth of 30-35 km, sometimes shallower. C2  
77 commonly appears as a second peak within a broader conductive layer encompassing C1. Finally,  
78 there are conductive upwellings (C3) that rise towards the surface beneath the Cascade arc, beyond  
79 the eastern edge of Fig. 2. These arc conductive features [32] can be seen more clearly in Fig.

80 Resistive material extends beyond the southern edge of Siletzia into the Klamath terranes but  
81 changes to a relatively thin upper crustal layer (R3). Resistivity variations in this area are complex  
82 (Fig.

83 To provide a more complete view of the geometry of key features, we derive two plan-view  
84 images from the 3D resistivity model. First, we compute the depth extent of the resistors (R1-  
85 R4), defined by the bottom of the 300 ohm-m isosurface. Fig. 3a shows that Siletzia is 30-40 km  
86 thick within a series of discrete blocks (a-e), distributed from Puget Sound to its southern edge,  
87 where a resistive curtain extends to the plate interface (f). This 3D image substantially revises  
88 previous models of Siletzia as a relatively uniform thick mafic block that thins to the north [5]. It  
89 is consistent with tomographic images [26] which show patches of high Vs just above the plate  
90 interface, coincident with thickened blocks of Siletzia. There is a sharp transition from the thin  
91 outboard edge of Siletzia to the much thicker blocks to the east, marking the abrupt termination  
92 of the underthrust accretionary complex. One exception to this picture occurs in central Oregon,  
93 where a resistive body near the coast (g) extends to the plate interface, with thinner Siletzia to the  
94 east underlain by a shallow conductor (R4/C2 in Fig. 2).

95

96 The geometry of Siletzia revealed here is consistent with its hypothesized origin as a chain of  
97 seamounts [39] or an oceanic plateau [40]. Plausibly, the deep resistive patches represent the roots  
98 of individual eruptive centers. Subsequent middle-Eocene basaltic magmatism [2] may also have  
99 contributed to the deep geometry of the resistive body, for example feature g in Fig. 3a. There is  
100 a strong correlation between the imaged thickness of Siletzia and modern tectonic and magmatic  
101 activity. Faults in the forearc crust [18] often trace the edges of the thickest resistive blocks, and  
102 tend to concentrate where Siletzia is thin (Fig. 3a). Variations in the western limit of arc vol-  
103 canism further correlate with the eastern edge of thick resistive material, with the most westerly  
104 eruptive vents always occurring where Siletzia is thin or absent, providing evidence for the role of  
105 pre-existing crustal structure in controlling modern magmatic processes [41].

106

107 A second view of the 3D model is provided by a 2D map of conductance integrated from the  
108 plate interface (Fig. 3b) to 10 km above it. A series of high-conductance features elongated par-  
109 allel to slab contours, and aligned with the outer edge of Siletzia, extend from 45°-48°N (Fig. 3b,  
110 A-C). The northern terminus of this band coincides with the limit of densest data coverage and is  
111 not well constrained. To the south, the band continues (I), albeit with decreased conductance near  
112 45°N. Downdip extents are variable, ranging from 35 km in the Olympic Peninsula, to 20 km in  
113 central Oregon. Except in the Olympic peninsula, these high conductance features are abruptly  
114 truncated by the thickened sections of resistive Siletzia (Fig. 3a), suggesting the downdip extent of  
115 subducted metasediments is primarily controlled by the deep geometry of Siletzia. In central Ore-  
116 gon structure is more complex: conductive feature D is somewhat deeper (35 km slab depth) and  
117 sits in the “cavity” formed by deep resistive features a, b, g. Along the southern edge of Siletzia a  
118 thin band of high conductance (E) is located just updip and north of the resistive curtain (f in Fig.  
119 3a).

120

121 Conductance above the plate interface is systematically enhanced in three zones: offshore (G,  
122 H, I), near the forearc mantle corner (FMC; A, B, C, D, F), and beneath the arc. As illustrated in  
123 Fig. 4a, conductance typically peaks at three nominal slab depths: 17.5 km, 30 km, and 65 km.  
124 The central peak (nearest the FMC) is typically highest and exhibits significant variations in depth.  
125 Depths for the seaward peak are more constant, 15-20 km, with conductance always decreasing  
126 trenchward. The third peak, beneath the arc, is much broader and exhibits significant variations  
127 with latitude. Here we restrict our attention to the two forearc peaks, which are plotted on a map  
128 of ETS density in Fig. 3c. In the Klamath terrane almost all forearc conductance peaks occur at  
129 nominal depths of 17.5 and 30 km, at least north of 42°N where we have good data coverage. In  
130 contrast, the depths of conductance peaks are more variable beneath Siletzia. Between 45°N and  
131 the southern Siletz margin, the seaward edge of resistive block a extends up-dip to 40 km, and  
132 conductance peaks are centered at depths of 17.5 and 35 km, similar to beneath the Klamaths.  
133 However, where resistive blocks b, c and d extend further up-dip, the inboard peak is deflected  
134 seaward to slab depths as shallow as 20 km.

135

136 As in previous works [32] we interpret the imaged conductive features to reflect fluids released  
137 from the subducting plate through sediment compaction, dehydration of clay minerals, pore col-  
138 lapse in meta-basalts, and a series of prograde metamorphic reactions [28, 42, 43]. Although these  
139 processes overlap in P/T space, pulses of fluid release are expected, as mechanical thresholds and  
140 the stability fields for different mineral assemblages are crossed [44, 45]. The observed characteris-  
141 tic depths for enhanced conductivity likely reflect these facies transitions. However, along-margin  
142 variations in lithology will further control the transport and accumulation of fluids, and we argue  
143 that these variations play a key role in determining the spatial pattern of imaged conductive anoma-  
144 lies and hence the distribution of subduction-zone fluids.

145

### 146 **3 Discussion**

147 Conductivity in the forearc crust varies by over three orders of magnitude, often abruptly. We  
148 believe these differences are most reasonably explained by variations in lithology and their corre-  
149 sponding changes in fluid content, especially given substantial seismic and geodynamic evidence  
150 for fluids in the megathrust [12, 29, 30] and in the overriding crust [25, 26]. Assuming resistivity  
151 can be taken as an indicator of lithology, the MT images provide important new constraints on the  
152 structure, creation, and evolution of the forearc crust. Most obviously, our results support previous  
153 suggestions [46, 47] that the entire Cascadia margin is being underthrust by accretionary complex  
154 material, driving uplift of the Coast Ranges [47, 48]. The 3D geometry of resistive Siletzia, seg-  
155 mented into thick, fault-bounded blocks, also provides clues to the genesis and accretion of this  
156 terrane and is deserving of further study.

157

158 The pattern of conductive features near the plate interface places constraints on fluid pathways,  
159 storage, and egress. Fluid inputs at the trench include pore water and hydrous minerals, both in the  
160 sediment layer, and in the underlying mafic oceanic crust. Free water in sediments is lost through  
161 compaction at depths  $< 10$  km [28, 49], and is thus not a likely source for the patches of high con-  
162 ductivity, which are all deeper. Similarly, dehydration of clay minerals should be largely complete  
163 at temperatures of  $\sim 150^\circ\text{C}$ , which in Cascadia is well above the depth of the imaged conductors.

164 That there is a peak in conductance at a relatively constant slab depth centered at  $\sim 17.5$  km sug-  
165 gests a possible fluid source at corresponding P/T conditions. We suggest this peak may at least  
166 in part result from pore collapse in the meta-basalts, which is believed to occur at temperatures of  
167 200-400°C [28, 44]. Our results may thus provide an important constraint on this poorly under-  
168 stood process.

169

170 The highest peaks in above-slab conductance at 20-35 km depth almost certainly reflect fluids  
171 released from meta-basalt dehydration. The subducting oceanic crust is relatively dry in Cascadia,  
172 with at most 5 vol% structural water [20]. We show in the SM that the implied rate of fluid input,  
173 together with estimates of the fluid volume required to explain observed conductance (Fig. 4a),  
174 places a lower bound on fluid residence times of  $\sim 1$  My. Thus, the main conductive features in  
175 Fig. 3b represent areas where fluids are sequestered on long timescales but they are not necessarily  
176 coincident with source regions.

177

178 Within Siletzia the main areas of fluid storage are just up-dip of thickened resistive crust that  
179 extends vertically to near the plate interface. We suggest that the sequestered fluids are sourced  
180 by dehydration reactions occurring at a relatively constant range of slab depths, 35 km and below,  
181 near the FMC [50]. We further suggest that Siletzia is relatively impermeable. Where these rocks  
182 extend to near the plate interface fluids are transported up-dip along the megathrust, ultimately  
183 leaking out along the outer edge where they accumulate in adjacent metasediments. Further sup-  
184 port for this conclusion is provided in Fig. 4b, where we plot above-slab conductance (Fig. 3b)  
185 averaged in non-overlapping strips centered upon the 22.5 and 32.5 km slab contours. An inverse  
186 correlation is evident; at latitudes where conductance is low along the deeper contour (i.e., where  
187 fluids do not accumulate in the overriding crust near the source), conductance (and by inference  
188 fluid accumulation) is increased up-dip. Even where the roots of Siletzia do not extend to the plate  
189 interface, as in the Puget Sound area, or are absent altogether (the Klamath terranes), peak con-  
190 ductance always occurs up-dip of the FMC (Fig. 3c), where tremor intensity is greatest and fluid  
191 pressures are presumably greatest. In fact, meta-basalt dehydration reactions occur over a range of  
192 depths, near and down-dip from the zone of maximum tremor [24, 43, 50]. Fluids escaping into  
193 the crust likely move more rapidly up-dip than vertically, due to strong permeability anisotropy

194 along and near the plate interface. In all cases fluid accumulation must reflect transport history as  
195 an integral over source depths and time. The resulting conductance peaks are thus shifted system-  
196 atically up-dip, relative to the fluid source. This shift is most extreme, and highly variable, where  
197 resistive Siletzia extends landward into the region where dehydration fluids are generated.

198

199 Seismically imaged low-velocity layers (LVLs) near the plate interface in Cascadia [12, 14, 24,  
200 25, 29, 51] and elsewhere [30, 52], have been interpreted either as a thin layer of fluid trapped in the  
201 upper oceanic crust [52], or a thicker zone of fluid-saturated metasediments in the overriding crust  
202 [25, 53, 54]. The MT results clearly support the latter view but are not necessarily inconsistent with  
203 the former. Fluid porosity in the LVL beneath Vancouver Island has been estimated to be 2.7-4%  
204 [29]. Assuming  $\sigma_f = 30$  S/m and an Archie's law exponent  $m = 1.5$ , a 5-km thick layer would  
205 have a conductance of 400-800 S, which if present everywhere, as some have argued [13, 14],  
206 would be resolved by our data. A thinner (e.g., 1 km) less saline layer would have a conductance  
207 of  $\sim 100$  S, and could be present over much of the forearc without violating the MT data (Fig. 3b).

208

209 Seismic reflection images [53, 55] from northern Cascadia may provide insight into the na-  
210 ture of the conductive underthrust rocks. A zone of high reflectivity (the E-layer), interpreted as  
211 a zone of strongly sheared metasediments with sub-horizontal fluid-filled cracks, thickens from 2  
212 km offshore where the megathrust is locked, to a thicker (5-7 km) layer inland where deformation  
213 becomes aseismic [53]. The conductive zones in Fig. 3b are likely similar, representing thick lay-  
214 ered metasedimentary packages that are ductile, strongly sheared, laminated, and filled with fluids.  
215 Indeed, the E-layer beneath Vancouver Island is conductive [54, 56].

216

217 As others have noted [15, 18, 26], there is a significant reduction in ETS frequency between  
218 43-47°N (Fig. 3c), coincident with the thickest parts of Siletzia (Fig. 3a). This has often been  
219 ascribed to reduced fluid pressures in this area, due either to reduced oceanic plate hydration [23]  
220 or permeability [16], or to enhanced fluid leakage into the overriding crust [18]. However, our  
221 results are consistent with a more uniform dehydration fluid source (e.g., note the relative unifor-  
222 mity of total fluid storage along margin implied by Fig. 4b), and a more impermeable forearc crust  
223 in Siletzia. Thus, fluid pressures near the FMC might be expected to be at least as high beneath

224 Siletzia as to the north or south. A source of fluid is a necessary, but not sufficient, condition for  
225 ETS, which most likely results from complex interactions in the stressed rock/fluid system [57],  
226 with cyclic fluctuations in fluid pressure driven by fracturing and fluid movement, followed by  
227 crack sealing by mineral precipitation [58, 59]. Our results suggest that variations in lithology,  
228 with corresponding variations in rheology, rock strength, and mineral dissolution and precipitation  
229 kinetics, may in fact be the primary driver of along-margin variations in ETS intensity in Cascadia.  
230 For example, there is some evidence that localized heterogeneity in the fault zone, with stronger  
231 asperities embedded in a more easily deformed matrix, may be a requirement for ETS [57, 60].  
232 Thus, one possibility is that fault zone heterogeneity is reduced in north-central Oregon. Here Silet-  
233 zia extends to near the plate interface, perhaps resulting in a narrower and more strongly sheared  
234 subduction channel, with fewer large asperities. It is also possible that transverse permeability of  
235 this more focused subduction channel is enhanced, allowing more rapid updip fluid escape, and  
236 reduced fluid pressures in the ETS zone.

237

238 To make further progress on understanding ETS in Cascadia the MT images should be more  
239 fully integrated with other geophysical constraints, especially more detailed seismic imaging and  
240 geodynamic models. In particular, the accretionary complex (mostly, but not completely, metased-  
241 iments) and the competent rocks of Siletzia apparently have vastly different capacities for fluid  
242 storage, allowing MT to effectively map lithology. As these units will also differ greatly in their  
243 rheological properties, the resistivity maps could have important implications for geodynamic  
244 modeling of the subduction zone [61], and may ultimately provide an improved understanding  
245 of variations in seismicity and plate coupling along the margin.

246

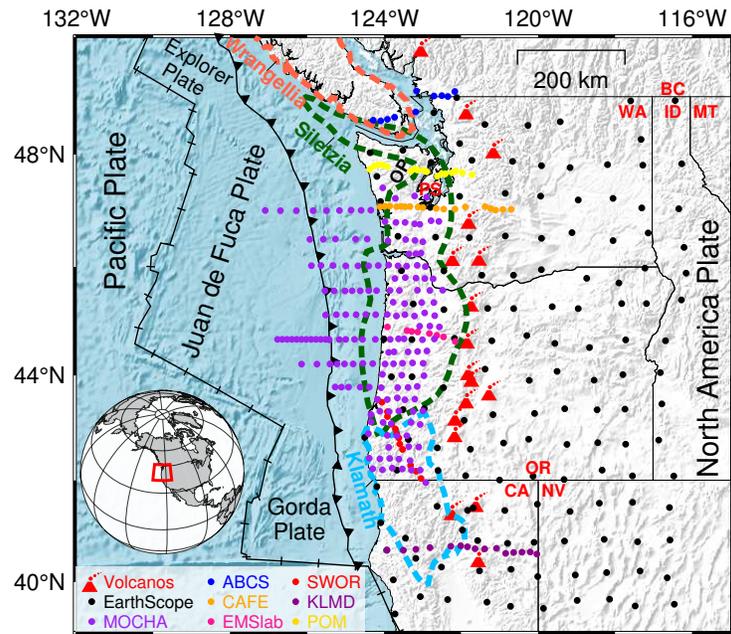


Figure 1: Map of study area, including tectonic plates, arc volcanoes (red triangles), state or provincial boundaries, and topography shown by shaded relief. Inset shows location on west coast of North America. Geologic terranes are outlined with dashed lines and labeled by color. Small filled circles show MT site locations, with different colors for each data source, as indicated in legend. Olympic Peninsula (OP) and Puget Sound (PS) are marked in the northern part of Siletzia.

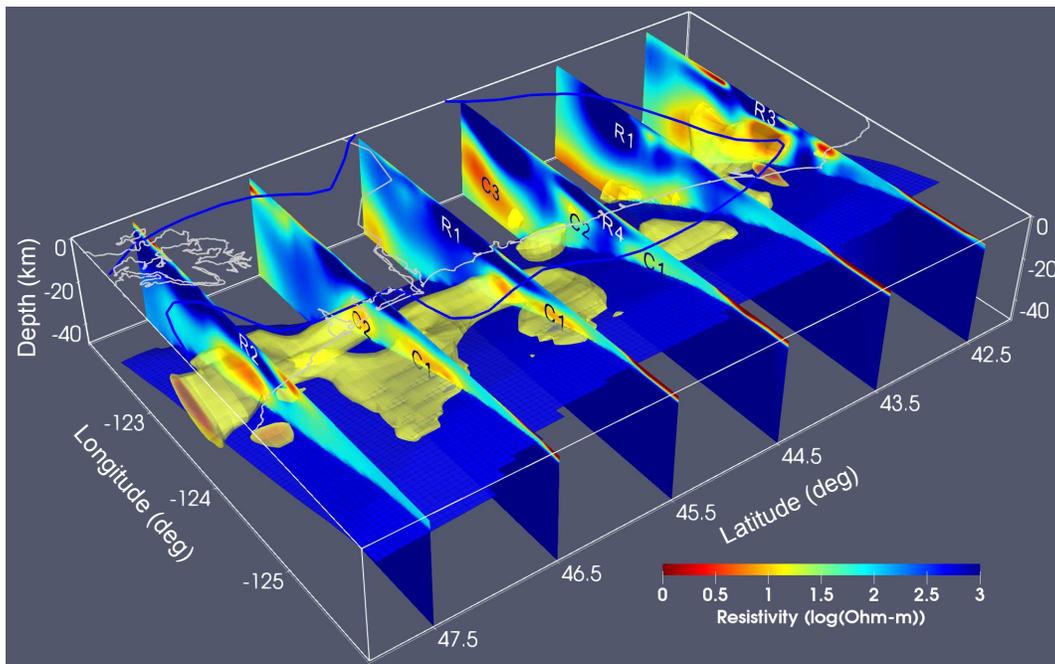


Figure 2: Three-dimensional view of the forearc portion of the model, including cross-sections at 1 degree increments, with superimposed iso-surface for conductive ( $\rho < 16$  ohm-m) material. Curved blue hatchured surface marks the top of subducting plate, from the model of [62]. Coastlines and state borders are shown with white lines, and the outline of Siletzia, as defined in [37] by the heavy blue line, both at  $z = 0$ . Principal conductive and resistive features discussed in text are labelled C1-C3, and R1-R4, respectively.

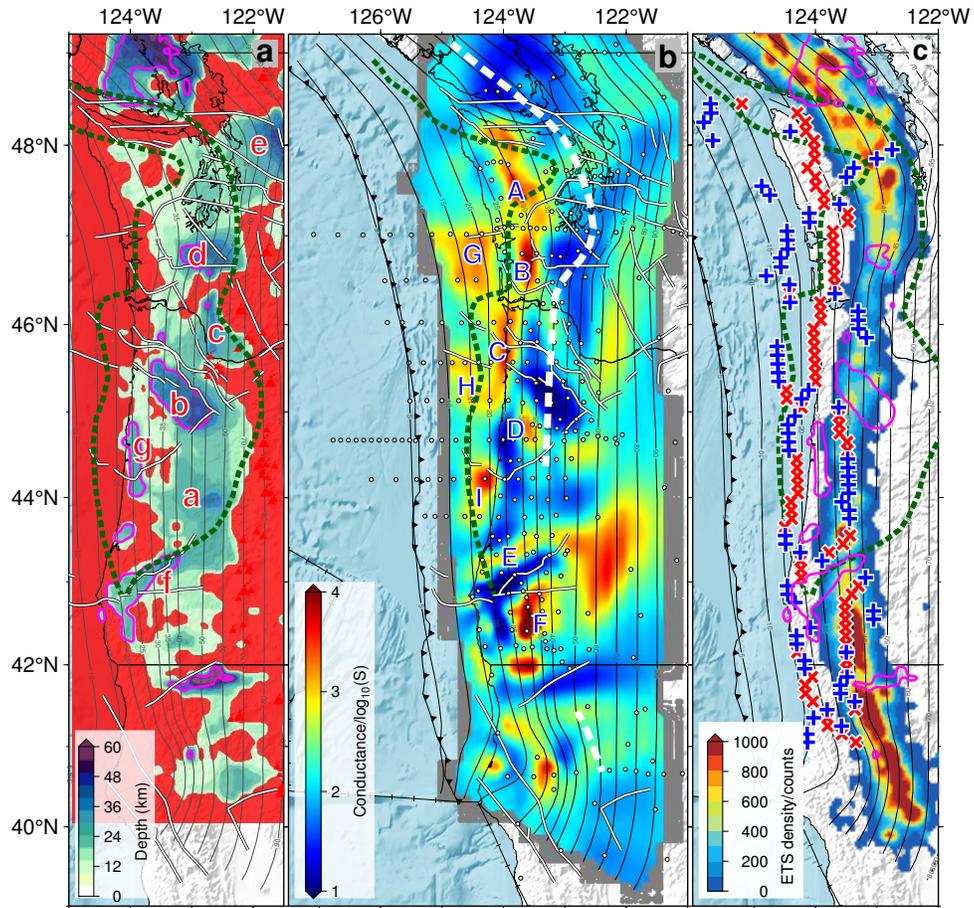


Figure 3: (a) Depth to bottom of the 300 ohm-m isosurface for the preferred model shown in Fig. 2. Examination of cross sections (see Fig.

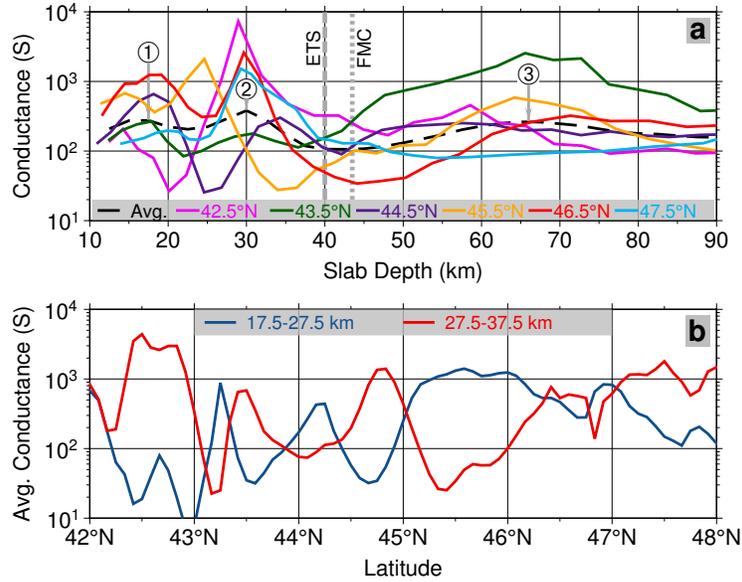


Figure 4: (a) Profiles of conductance as a function of slab depth at latitudes corresponding to sections shown in Fig. 2 (colored lines), and average computed along depth contours (black dashed line). The three nominal peaks (seaward, FMC, and arc) are marked as 1, 2, 3 respectively. Mean slab depth of peak ETS density and FMC position are also marked. (b) Conductance plotted in Fig. 3b, averaged over non-overlapping bands of slab depths (17.5-27.5 and 27.5-37.5 km).

## 247 4 Methods

248 **MT data and inversion.** Our MT dataset includes 165 new sites (including 64 offshore) from  
 249 the “MOCHA” (Magnetotelluric Observations of Cascadia with a Huge Array) project, straddling  
 250 the continental margin in the central part of the Cascadia subduction zone (Fig. 1). These were  
 251 supplemented with 142 EarthScope transportable array MT sites [36], and legacy long-period data  
 252 from six 2D profiles (13 from EMSLAB [31], 18 from CAFE-MT [35], 13 from SWORMT [63],  
 253 15 from KLMD [32], 17 from POM [64], and 10 from ABCS [33]). Most of the profile data  
 254 are discussed in ref. [32]. For the MOCHA onshore sites and the five profiles, site spacing is  
 255 typically 10-20 km, while for the offshore sites, profiles are 40 km apart, with 20 km site spacing,  
 256 except for the central profile at 44.25°N with a denser site spacing of 5 km. In total there are 393  
 257 long-period MT sites, with data covering the period ranging from 7 s to 14,678 s. All onshore sites  
 258 measured vertical magnetic fields; offshore sites did not. Data were processed with standard robust  
 259 processing algorithms [65, 66].

260 For 3D inversion we use the ModEM code [67]. This inversion code allows specification of

261 a prior model, deviations from which are penalized by the regularization. For results reported  
262 here the prior model consists of a layered 1D earth with resistivity 100 ohm-m to a depth of 410  
263 km, 10 ohm-m from there to 660 km, and 3 ohm-m below this. In the first layers of the starting  
264 model we included known bathymetry with the resistivity of seawater set to 0.3 ohm-m [68]. We  
265 also included a realistic ocean sediment distribution [69], with electrical resistivity of the saturated  
266 sediments set to 1 ohm-m.

267 Finally, we included the subducting Juan de Fuca plate in the prior model as a dipping re-  
268 sistive layer, as in [35]. There is substantial evidence that oceanic lithosphere is generally quite  
269 resistive [70]. While resistive bodies are observed in MT inversions over subducting plates [32,  
270 71] without imposing this structure, MT data cannot vertically resolve resistive bodies. Imposing  
271 the seismically-constrained geometry of the plate in the prior model can improve resolution of  
272 conductive features above the plate. We take depth to the top of the shallower parts of the plate  
273 (above 90 km) from ref. [62], and extend to greater depths using seismic tomographic models [72].  
274 Thickness of the plate was set to 50 km based on published models of resistive ocean lithosphere  
275 [73]. The assumed prior model is summarized graphically in Fig.

276 The model grid has a resolution of  $11 \times 9.5$  km horizontally, with vertical layers logarithmically  
277 spaced, starting from very fine layers (20 meters) to discretize the bathymetry and the seafloor sed-  
278 iments. The resulting core grid is  $120 \times 108 \times 57$  cells in latitude, longitude and depth respectively.  
279 Edges were padded with 7 logarithmically increasing grid cells on all sides, to extend the bound-  
280 ary of the study area  $\sim 600$  km in all directions. We ran the inversion more than 30 times, to test  
281 effects of various parameters (e.g., thickness of the slab, effect of offshore sites, effects of seafloor  
282 sediments, covariance parameters) and inversion strategies.

283 The seafloor data presented some challenges to the 3D inversion. Lack of short period data  
284 makes it difficult to constrain shallow structure such as the seafloor sediments. Inversion tests  
285 demonstrated that including *a priori* a realistic distribution of conductive ocean sediments was  
286 essential; anomalous phases of the nominal TE mode (coast-parallel electric field) impedances,  
287 both for onshore and offshore sites near the coast could not be fit without imposing this structure.  
288 Note that including a sediment layer greatly reduces the resistivity contrast at the seafloor, likely an  
289 important factor in improving performance of the inversion. Even with the sediment layer imposed  
290 fitting the offshore TE mode data was a challenge, which we found could be mitigated with a two-

291 stage inversion strategy: we first run the inversion with large error floors on the nominal TM mode  
292 (10% of the off-diagonal impedance), forcing the inversion to concentrate on fitting TE mode data,  
293 for which we used a smaller (5%) error floor. In the second stage the inversion was restarted, using  
294 5% error floors for both modes. The final normalized root-mean-square (nRMS) misfit with 5%  
295 error floors was 2.74, with the largest misfits for TE mode data localized on the seafloor and near  
296 the coast (see Fig.

297 Responses for the full dataset (observed and computed) are shown as maps of apparent resis-  
298 tivity and phase in the SM in Figs.

## 299 **5 Data availability**

300 The data that support the findings of this study are publicly available online at <http://ds.iris.edu/ds/products/emtf/>.

## 301 **6 Code availability**

302 ModEM is freely available online at <https://sites.google.com/site/modularem/> to the academic  
303 community.

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477

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479 D.L., K.K., A.S. and P.B. collected the new MT data. A.K. processed the new land data and K.K.  
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481 interpretation, and wrote the manuscript, with input from all other authors.

482 **Competing Interests:**

483 The authors declare no competing interests.

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