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# Unveiling shrouded oceans on temperate sub-Neptunes via transit signatures of solubility equilibria vs. gas thermochemistry

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- The recent discovery and initial characterization of sub-Neptune-sized exoplanets that receive stellar irra-
- <sup>2</sup> diance of approximately Earth's<sup>1-7</sup> raised the prospect of finding habitable planets in the coming decade.
- 3 Some of these temperate planets may support liquid water oceans, if they do not have massive H<sub>2</sub>/He en-
- 4 velopes and are thus not too hot at the bottom of the envelopes<sup>8-11</sup>. For planets larger than Earth, and
- $_{5}$  especially planets in the  $1.7-3.5~R_{\oplus}$  population <sup>12</sup>, the mass of the  $H_{2}/He$  envelope is typically not suffi-
- 6 ciently constrained to assess the potential habitability 13-16. Here we show that the solubility equilibria vs.
- 7 thermochemistry of carbon and nitrogen gases results in observable discriminators between small H<sub>2</sub> at-
- 8 mospheres vs. massive ones. On temperate sub-Neptunes, the condition to form a liquid-water ocean and
- 9 that to achieve the thermochemical equilibrium are mutually exclusive. The dominant carbon and nitrogen
- gases are typically CH<sub>4</sub> and NH<sub>3</sub> due to thermochemical recycling in a massive atmosphere of a temperate

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planet, and those in a small atmosphere overlying a liquid-water ocean are most likely  $CO_2$  and  $N_2$ , followed by CO and  $CH_4$  produced photochemically.  $NH_3$  is depleted in the small atmosphere by dissolution into the liquid-water ocean<sup>17</sup>. These gases lead to distinctive features in the planet's transmission spectrum, and a moderate number of repeated transit observations with the James Webb Space Telescope should readily tell apart a small atmosphere vs. a massive one via these spectral features on planets like K2-18 b. This method thus provides a way to use near-term facilities to constrain the atmospheric mass and habitability of temperate sub-Neptune exoplanets.

The exoplanet community already has ways to detect an  $H_2$  atmosphere by transmission spectroscopy via its pressure scale height one order of magnitude larger than that of an  $N_2$  or  $CO_2$  atmosphere  $^{18}$ . However, the mass of the  $H_2$  atmosphere – the parameter that controls the temperature at the bottom of the atmosphere and thus the possibility for liquid water  $^{8-10}$  – is not directly measurable from the transmission spectrum. Also, a planet's mass and radius typically allow multiple models of the interior structure  $^{13,14}$ . It is unclear whether the planets in the  $^{12}$   $^{12}$   $^{12}$  population are mostly rocky planets with massive  $^{12}$   $^{12}$   $^{12}$   $^{12}$  or planets with a massive water layer ( $^{12}$   $^{12}$   $^{12}$   $^{12}$   $^{12}$   $^{12}$   $^{12}$   $^{12}$   $^{12}$   $^{12}$   $^{12}$   $^{12}$   $^{13}$   $^{14}$   $^{14}$   $^{15}$   $^{16}$   $^{12}$   $^{15}$   $^{16}$   $^{15}$ 

The temperate sub-Neptune K2-18 b is a harbinger of a class of planets that might be habitable and exemplifies the need for a near-term method to measure the size of an H<sub>2</sub> atmosphere. The planet of  $8.6~M_{\oplus}$  and  $2.6~R_{\oplus}$ is in the habitable zone of an M dwarf star, and has a transmission spectrum (obtained by *Hubble* at  $1.1 - 1.7~\mu m$ ) with confirmed spectral features, which indicates that the planet should host an atmosphere dominated by H<sub>2</sub><sup>6,7</sup>. Interior structure models showed that the planet can have a massive (> $\sim 1000$  bar)  $\rm H_2$  atmosphere overlaying a rocky/Fe core and a supercritical water layer, or a smaller (< 100 bar)  $\rm H_2$  atmosphere with a water-dominated interior<sup>11,16</sup>. For K2-18 b, specifically, a  $\sim 10-100$  bar  $\rm H_2$  atmosphere overlaying a water layer would cause > 200 bar of water to evaporate into the atmosphere, resulting in a hot steam atmosphere inconsistent with the observed transmission spectrum<sup>26</sup>. An even smaller,  $\sim 1$  bar  $\rm H_2$  atmosphere would prevent this steam atmosphere and produce a liquid-water ocean (see Methods), but requires a very small rocky/Fe core and may be disfavored from the planet formation standpoint<sup>27</sup>. However, a planet slightly more massive or smaller than K2-18 b – such as those at the center of the 1.7-3.5  $R_{\oplus}$  planet population<sup>15</sup> – does not have this small-core difficulty to have a small atmosphere, and many such planets and planet candidates have been detected and will soon be available for transmission spectroscopy (Fig. 1).

How can we distinguish a planet with a massive  $H_2$  atmosphere versus an ocean planet with a small  $H_2$  atmosphere? Here we propose, for temperate planets, these two scenarios can be distinguished by detecting signature gases of the gas-phase thermochemical equilibrium vs. solubility equilibria. The key difference between the two scenarios is that the gas-phase thermochemical equilibrium would be achieved in the deep and hot part of the massive atmosphere, and in contrast, it would not be achieved in a small atmosphere overlying a liquid-water ocean. Instead,  $NH_3$  and sulfur species would be sequestered by the ocean  $^{17,28}$  and the abundance of  $CO_2$  would be set by the ocean chemistry (Fig. 2, with the cosmochemical and geological constraints detailed in Methods). This fundamental difference, coupled with photochemical processes, leads to distinctive gas abundances in the observable part ( $<\sim 0.1$  bar) of the atmosphere.

If the planet has a massive  $H_2$  atmosphere, thermochemical reactions in the deep atmosphere recycle O, C, N, S species into  $H_2O$ ,  $CH_4$ ,  $NH_3$ , and  $H_2S^{29-32}$ .  $H_2O$  can form a cloud and the above-cloud  $H_2O$  may be partially depleted as a result<sup>33-35</sup>. Recent calculations have shown that the photodissociation of  $NH_3$  in the presence of  $CH_4$  leads to the formation of HCN, and that CO and  $CO_2$  are produced by the photodissociation of  $CH_4$  together with  $H_2O^{35}$ . The photodissociation of  $H_2S$  leads to the formation of elemental sulfur haze<sup>36,37</sup>, but the haze would

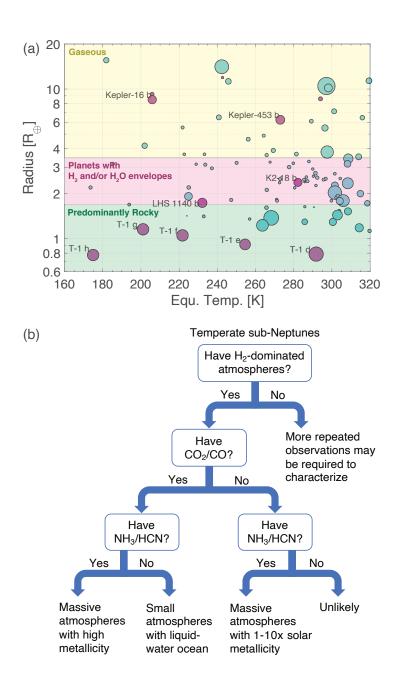


Figure 1: Temperate exoplanets amenable for atmospheric characterization via transmission spectroscopy. (a) Purple dots are confirmed planets with measured masses, and blue dots are planets with unknown masses or planet candidates. Data are taken from the NASA Exoplanet Archive and and the TESS Objects of Interest Catalog. The marker sizes are scaled with the expected S/N of the spectral features of an  $H_2$  atmosphere observed by JWST at 2  $\mu m$ . Most of the temperate planets and planet candidates suitable for atmospheric characterization are larger than Earth and thus more likely to have  $H_2$  atmospheres. (b) A roadmap to characterize the mass of the atmospheres and the habitability of temperate sub-Neptunes by detecting signature gases. See text for details.

be close to the cloud deck and would not mute transmission spectral features<sup>35</sup>. These photochemical products are transported to the deep atmosphere and recycled back to  $CH_4$ ,  $NH_3$ , and  $H_2S$ . An exception is the  $100 \times$  solar metallicity atmosphere, where the thermochemical-equilibrium mixing ratios of  $N_2$  and  $NH_3$  may be similar and those of CO and  $CO_2$  can be substantial compared with the mixing ratios produced by the photodissociation of  $CH_4$  and  $H_2O$  (see Methods).

If the planet instead has a small atmosphere and a liquid-water ocean, this thermochemical recycling cannot occur. Instead, CO<sub>2</sub> is the preferred form of carbon in equilibrium with a massive amount of H<sub>2</sub>O<sup>32,38</sup>, and NH<sub>3</sub> is dissolved in the ocean and largely depleted from the atmosphere<sup>17</sup>. The abundance of atmospheric CO<sub>2</sub> is controlled by the oceanic pH<sup>39-42</sup> and that of N<sub>2</sub> is probably a combined result of the initial endowment and atmospheric escape. A reasonable lower bound of the total mass of CO<sub>2</sub> in the H<sub>2</sub> and H<sub>2</sub>O layers (Fig. 2) can be derived from the cosmochemical constraints of planetary building blocks and the partitioning between the iron core, the silicate mantle, and the water layer (see Methods). Also, the "seafloor" of this thin-atmosphere, H<sub>2</sub>O-rich sub-Neptune will not be not a sharp interface in density and composition, but instead have a finite thickness<sup>43</sup>. The interface will be compositionally stratified with denser material underlying less dense material, and material transport across this "fuzzy layer" is inhibited due to the stratification. Thus, any carbon or nitrogen added to the H<sub>2</sub> and H<sub>2</sub>O envelope by planetesimal accretion late in planet growth will remain in the envelope, and will not be stirred down into the silicate layer. Meanwhile, transit observations can straightforwardly identify H<sub>2</sub>-dominated atmospheres and rule out CO<sub>2</sub> or N<sub>2</sub>-dominated ones only from the size of spectral features<sup>18</sup>.

We have used an atmospheric photochemical model<sup>44</sup> coupled with a radiative-convective model<sup>26</sup> to determine the equilibrium abundances of CO, CH<sub>4</sub>, and C<sub>2</sub>H<sub>6</sub> in small and temperate H<sub>2</sub> atmospheres, for a cosmochemically and geologically plausible range of CO<sub>2</sub> abundance, and compared the compositions and transmission spectra with massive H<sub>2</sub> atmospheres<sup>35</sup>. The input parameters and atmospheric chemistry results of the small-atmosphere models are summarized in Methods. We have used the planetary parameters of K2-18 b<sup>7,45</sup> and the UV spectrum of the M dwarf star GJ 176<sup>46</sup> (similar to K2-18<sup>47</sup>) in these models, even though K2-18 b probably

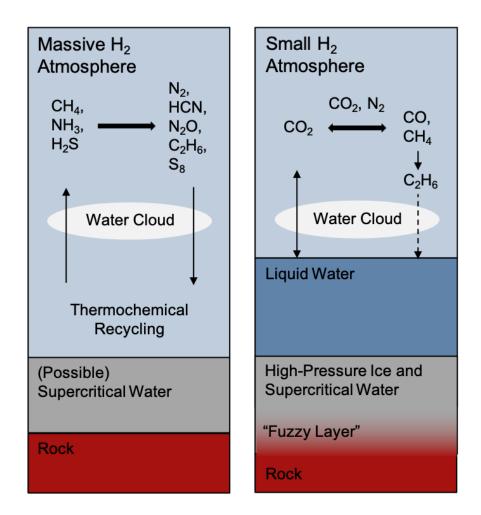


Figure 2: Interior structures of temperate H-rich exoplanets and the associated ranges of atmospheric composition. If the planet has a massive  $H_2$  atmosphere, the deep atmosphere would be hot – enabling thermochemical recycling – but a liquid-water surface would not be possible. If the planet has a small  $H_2$  atmosphere, a liquid-water surface may be possible. On these planets, the equilibrium abundance of atmospheric  $CO_2$  is set by the oceanic chemistry and that of  $N_2$  by atmospheric evolution.

does not have a small atmosphere. The results should be broadly applicable to temperate sub-Neptunes around M dwarf stars, because the differences in atmospheric composition and transmission spectrum are driven by gas-phase thermochemical equilibrium vs. solubility equilibria, and not by basic planetary parameters such as surface gravity. Fig. 3 compares the expected spectra for the massive-atmosphere scenarios and the small-atmosphere scenarios. For K2-18 b, both the massive-atmosphere models with  $1 - 100 \times$  solar metallicity and the small-atmosphere models with a low mixing ratio of  $CO_2$  (400 ppm) provide good fits to the transmission spectrum measured by *Hubble*.

Measuring the transmission spectra in an expanded wavelength range of  $1-5~\mu m$  will distinguish the small atmospheres from massive ones. Using K2-18 b as an example for temperate sub-Neptunes, we see that the massive-atmosphere models and the small-atmosphere models, while having differences within each group, can be distinguished using the spectral regions of 1.7-2.1, 2.6-3.1, and  $4.1-4.9~\mu m$  (the shaded areas in Fig. 3). Both the massive-atmosphere and small-atmosphere models show spectral features of  $H_2O$  and  $CH_4$ , and so observing these two gases alone is unlikely to separate the scenarios. At 1.7-2.1 and  $2.6-3.1~\mu m$ , the transmission spectra show  $NH_3$  and HCN absorption in massive atmospheres but not in small atmospheres. At  $\sim 2.8~\mu m$ , the transmission spectra of small atmospheres show  $CO_2$  absorption. Between 4.1 and 4.9  $\mu m$ , the transmission spectra of small atmospheres (the low- $CO_2$  cases) have prominent features of  $CO_2$  and  $CO_3$  and these features are much weaker in the spectra of massive atmospheres.

The massive atmosphere with  $100\times$  solar metallicity can have CO and CO<sub>2</sub> produced from CH<sub>4</sub> and H<sub>2</sub>O photolysis<sup>35</sup>, and also from gas-phase thermochemical equilibrium (see Methods). The mixing ratios of CO and CO<sub>2</sub> can be comparable with the small atmosphere models. Will these effects create a "false positive" where a massive atmosphere might be mistaken as a small atmosphere, since they could have similar mixing ratios of CO and CO<sub>2</sub> at the  $\sim 0.1$  bar level? Fig. 3 indicates that this mistake will be unlikely, because (1) the spectrum of the massive atmosphere also features NH<sub>3</sub> and HCN absorption but the small atmosphere spectrum does not; (2) the massive atmosphere with  $100\times$  solar metallicity has a greater mean molecular weight ( $\sim 5$ ) than the small

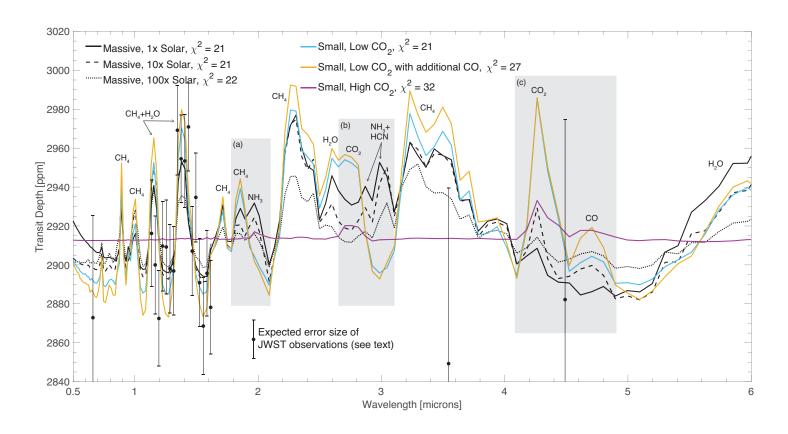


Figure 3: Expected transmission spectrum of temperate sub-Neptune planets of M dwarf stars, using K2-18 b as an example and comparing with the planet's transit depth observed by K2, Hubble and  $Spitzer^7$ . The massive-H<sub>2</sub>-atmosphere models and the small-H<sub>2</sub>-atmosphere models differ in three spectral regions: in (a) and (b), the massive-atmosphere models have absorption features of NH<sub>3</sub> and HCN, while the small-atmosphere models do not; in (c), the small-atmosphere models with a low mixing ratio of  $CO_2$  (400 ppm) have prominent features of  $CO_2$  and CO, while the massive-atmosphere models only have small features of  $CO_2$ . The small-atmosphere models with a high mixing ratio of  $CO_2$  (10%) has a high mean molecular weight ( $\sim$  7) and a cloud top at  $\sim 10^{-3}$  bar (as opposed to  $\sim 3 \times 10^{-2}$  bar the low- $CO_2$  cases) and thus muted spectral features.

atmosphere ( $\sim 2$  in the low-CO<sub>2</sub> cases) and thus smaller transmission features; and (3) the massive atmosphere will have CO<sub>2</sub>/CO $<\sim 0.1^{35}$ , and the small atmosphere typically has CO<sub>2</sub>/CO> 1 (Table 1). Furthermore, a massive atmosphere with  $\gg 100\times$  solar metallicity will have a mean molecular weight much higher than that of an H<sub>2</sub> atmosphere and is thus also distinguishable by transmission spectroscopy.

With moderate time investment (i.e., < 100 hours), JWST will provide the sensitivity to detect the signature 111 gases aforementioned and distinguish massive versus small atmospheres on planets like K2-18 b. As an example, we have used PandExo<sup>48</sup> to simulate the expected photometric precision using JWST's NIRSpec instrument. If 113 ombining two transit observations with NIRSpec's G235H grating and four transits with the G395H grating, the 114 overall photometric precision would be  $\sim 20$  ppm per spectral element at a resolution of R=100 in both chan-115 nels that cover a wavelength range of  $1.7 - 5.2 \mu m$ . These observations would distinguish the small-atmosphere 116 scenarios versus the massive-atmosphere scenarios at  $> 4.5\sigma$ . Additionally, spectral retrievals using Tau-REx<sup>49</sup> 117 indicate that the abundance of NH<sub>3</sub> and HCN in massive atmospheres and that of CO<sub>2</sub> and CO in small atmo-118 spheres would be usefully constrained (Fig. S1). The dedicated exoplanet atmosphere characterization mission 119 ARIEL could also provide the sensitivity to detect these gases with more repeated transit observations<sup>50</sup>. 120

Taken together, the results presented here identify a near-term path to detect small H<sub>2</sub> atmospheres that can 121 be consistent with liquid-water oceans on temperate exoplanets. H<sub>2</sub> atmospheres are probably the only type of 122 temperate atmospheres readily within the reach of JWST and ARIEL for detailed studies, since to characterize a 123 heavier N<sub>2</sub> or CO<sub>2</sub> atmosphere will require co-adding a few tens transits – something not impossible but probably very hard<sup>51–55</sup>. The mass of the H<sub>2</sub> atmospheres – a parameter that is not directly measured by transits but critical 125 for habitability – can be inferred from transmission spectra via the signature gases that indicate solubility equilibria 126 versus gas-phase thermochemical recycling. The biggest uncertainty is probably the temperature at the  $\sim 100$ -bar 127 pressure level in the massive-atmosphere scenarios, which may be affected by ad hot heating mechanisms such 128 as tidal heating. Detailed models of the interior temperature and mixing may further constrain this uncertainty<sup>56</sup>. 129 Based on the considerable range of the parameter space explored, we suggest that the sensitivity of multiple gases

provided by future observatories' expanded wavelength coverage over *Hubble* would enable broad categorization of small versus massive atmospheres, as summarized in Fig. 1. If some of the temperate planets and planet candidates have small H<sub>2</sub> atmospheres, their relative ease for transit observations would significantly enhance the prospect of detecting and characterizing potentially habitable exoplanets within the next decade.

#### **Methods**

Mutual exclusivity of habitability and thermochemical equilibrium On temperate sub-Neptunes, the condition to form a liquid-water ocean and that to achieve the thermochemical equilibrium of carbon and nitrogen
molecules are mutually exclusive. The conversion rate between  $CO/CO_2$  and  $CH_4$ , and that between  $N_2$  and  $NH_3$ , are primarily a function of the temperature and to a lesser extent the pressure<sup>57,58</sup>. To effectively recycle a
molecule, its chemical lifetime must be shorter than the vertical mixing timescale. The vertical mixing timescale,
in turn, is related to the internal heat flux of the planet in the convective layer of the atmosphere<sup>59</sup>. To have a
chemical lifetime of CO that equals to or is shorter than the vertical mixing timescale, the temperature is typically  $NH_3$ 000 K.  $N_2$ 1 is expected to reach thermochemical equilibrium at even higher temperatures.

Fig. 4 shows the "quench point" of carbon and nitrogen species on a temperate sub-Neptune like K2-18 b with massive  $H_2$  atmospheres, based on the pressure-temperature profiles adopted in the massive-atmosphere models<sup>35</sup>. The "quench point" is defined as the pressure level where the chemical lifetime of a gas equals to the vertical mixing timescale<sup>57,58</sup>. The gas is close to thermochemical equilibrium at or below the quench point, and its mixing ratio is carried to the atmosphere above the quench point by vertical mixing. On temperate sub-Neptunes, carbon species should reach thermochemical equilibrium at the temperature of 1000 - 1200 K and the pressure of > 100 bar. Nitrogen species reach thermochemical equilibrium at the higher temperature of  $\sim 1400$  K. These temperatures are substantially higher than the critical point of water. As such, we do not expect thermochemical equilibrium of carbon and nitrogen molecules on a temperate planet with a liquid-water ocean. Conversely, a temperate planet with a massive  $H_2$  atmosphere that achieves thermochemical equilibrium at depths

is not habitable.

One might also consider the intermediate situation between massive atmospheres with thermochemical equi-155 librium and small atmospheres with liquid-water oceans, e.g., the atmospheres with a surface pressure of 1-100156 bars (Fig. 4). For many sub-Neptunes, this intermediate-atmosphere scenario would still require a massive water 157 layer underneath to explain their mass and radius. If water is in the liquid form at the interface with the atmo-158 sphere, the evaporation of this ocean will make the atmosphere H<sub>2</sub>O-dominated<sup>26</sup>. If water is supercritical, any 159 H<sub>2</sub> layer of 1 – 100 bars should be well mixed with the water layer. Therefore, such an intermediate endowment 160 of H<sub>2</sub> would most likely result in a non-H<sub>2</sub>-dominated atmosphere, which is distinguishable with transmission 161 spectroscopy<sup>18</sup>. 162

Massive atmosphere models. The photochemical models of massive,  $H_2$ -dominated atmospheres are described in a companion paper<sup>35</sup>. In this work, we adopt the massive atmosphere models of K2-18 b and compare them with the small atmosphere models. The massive atmosphere models explore the atmospheric metallicity of  $1 - 100 \times$  solar abundance.

The photochemical models of the massive atmospheres assume that the dominant O, C, N, and S species are 167 H<sub>2</sub>O, CH<sub>4</sub>, NH<sub>3</sub>, and H<sub>2</sub>S at the base of the photochemical domain, which was defined as the pressure level 10-168 fold greater than the tropopause pressure (i.e., typically around 1 bar)<sup>35</sup>. Gases produced in the modeled domain 169 can be transported through the lower boundary. This setup assumes that thermochemical recycling in the deeper 170 atmosphere effectively recycles the photochemical products into H<sub>2</sub>O, CH<sub>4</sub>, NH<sub>3</sub>, and H<sub>2</sub>S. Here, based on the 171 quench points shown in Fig. 4, we can further assess how realistic this assumption is. Fig. 4 shows that, at the 172 quench point of CO-CH<sub>4</sub>, the temperature is likely deep in the CH<sub>4</sub>-dominated regime. The same is also true for 173 the  $N_2$ -NH<sub>3</sub>, except for the  $100 \times$  solar-metallicity case in which the quench point is close to the equal-abundance 174 boundary between N<sub>2</sub> and NH<sub>3</sub>. A quantitative description of this behavior is provided below. Note that the 175 quench point depends on the strength of vertical mixing and the deep-atmosphere temperature depends on the 176 thermal history of the planet, and so the detail behavior may be complex<sup>56</sup>.

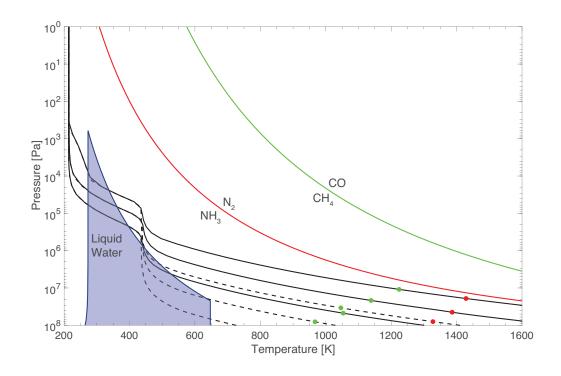


Figure 4: Modeled pressure-temperature profiles of a sub-Neptune exoplanet like K2-18 b with massive  $H_2$  atmospheres that receive stellar irradiance of Earth's, for an internal heat flux of  $T_{\rm int}=60$  K (similar to Neptune, black solid lines) and  $T_{\rm int}=30$  K (similar to Earth, black dashed lines). The pressure-temperature profiles assume  $1\times$ ,  $10\times$ , and  $100\times$  solar metallicity. The shaded area shows the pressure and temperature range where pure liquid water is stable. The green and red lines show the equal-abundance boundaries for major carbon and nitrogen gases in a solar-metallicity gas in thermochemical equilibrium, and the green and red dots show the expected quench point for CO and that for  $N_2$  respectively. The quench points are estimated by equating the chemical lifetime to the vertical mixing timescale  $^{57,58,60}$ , with the eddy diffusion coefficient estimated using a mixing-length theory  $^{59}$  ( $\sim 10^4$  m $^2$  s $^{-1}$  at the pressure of  $10^6-10^8$  Pa). The quench point for CO $_2$  follows that of CO, and similarly, that of HCN occurs at a similar pressure and temperature as  $N_2$   $^{57,58}$ .

Applying a thermochemical equilibrium model<sup>38</sup>, we find that for  $1-10\times$  solar metallicity and at the quench pressure, the mixing ratio of CO is  $10^{-10}-10^{-8}$  and that of CO<sub>2</sub> is  $< 10^{-10}$ . These quantities are substantially smaller than the respective mixing ratios found by the photochemical models at the lower boundary. The mixing ratio of N<sub>2</sub> is  $10^{-6}-10^{-4}$ , and that of HCN is minimal. These estimates indicate that the  $1-10\times$  solar massive atmospheres on temperate planets are safely in the CH<sub>4</sub>- and NH<sub>3</sub>-dominated regime, and the assumption that the dominant C and N species are CH<sub>4</sub> and NH<sub>3</sub> at the base of the photochemical domain is valid.

For the  $100 \times$  solar metallicity case, the mixing ratios of CO at the quench pressure can reach  $10^{-5} - 10^{-2}$  and 184 that of  $CO_2$  can reach  $10^{-6} - 10^{-3}$ , with the pressure-temperature profiles shown in Fig. 4, which can exceed the 185 mixing ratios found by the photochemical model at  $\sim 1$  bar level<sup>35</sup>. While CH<sub>4</sub> is still the dominant carbon-bearing 186 molecule, a higher deep-atmosphere temperature would result in greater mixing ratios of CO and CO<sub>2</sub>. This 187 comparison indicates that the mixing ratio of CO and CO<sub>2</sub> in the part of the atmosphere probed by transmission 188 spectroscopy (typically  $10^2 - 10^4$  Pa) can indeed reach or exceed the  $10^{-6} - 10^{-2}$  level. The mixing ratio of 189  $N_2$  at the quench pressure is comparable to that of NH<sub>3</sub> and that of HCN is  $\sim 10^{-6}$ . Again, a higher deepatmosphere temperature favors N2 and further reduces the equilibrium abundance of NH3. The input mixing ratio 191 of NH<sub>3</sub> in the photochemical model for the  $100 \times$  solar massive atmosphere may thus be overestimated, while the photochemical production of HCN should still dominate in the observable part of the atmosphere. Because the 193 mixing ratio of NH<sub>3</sub> expected in small atmospheres is  $< 10^{-12}$  (Fig. 6), and because the spectral features of 194  $NH_3$  and HCN are already prominent at the  $10^{-6} - 10^{-4}$  level<sup>35</sup>, even reducing the mixing ratio of  $NH_3$  in the 195 massive atmosphere by a few orders of magnitude would not make its transmission spectrum similar to that of a small atmosphere. For example, the mixing ratio of NH<sub>3</sub> would be  $\sim 10^{-2}$  if all nitrogen is in NH<sub>3</sub> for the  $100\times$ 197 solar metallicity, and a reduction from that by three orders of magnitude would still result in a mixing ratio of 198  $10^{-5}$ . Therefore, and as discussed in the main text, we do not expect the  $100\times$  solar or higher metallicity massive 199 atmosphere would likely become a false-positive imitation of small atmospheres. 200

Small atmosphere models. We use an atmospheric photochemical model<sup>44</sup> to simulate the atmospheric chem-201 istry in small H<sub>2</sub> atmospheres and trace the steady-state abundances of CO, CH<sub>4</sub>, C<sub>2</sub>H<sub>6</sub>, and other potential gases 202 of interest. The photochemical model includes a comprehensive reaction network for O, H, C, N, and S species 203 (including sulfur aerosols, hydrocarbons, and the reactions important in H<sub>2</sub> atmospheres), and it has been used to 204 study the lifetime and equilibrium abundance of potential biosignature gases in H<sub>2</sub> atmospheres<sup>61</sup>. The pressure-205 temperature profiles (Fig. 5) used as the basis for the photochemical model are calculated with the climate module 206 of 1D-TERRA<sup>26</sup>. The module uses a correlated-k approach with the random overlap method to include molecular 207 absorption, collision-induced opacities, and the continuum of water vapor to calculate the radiative equilibrium, 208 and the appropriate (moist or dry) adiabatic lapse rate to apply the convection adjustment. The module has been 209 tested against the cases of Earth, Venus, and Mars, as well as with other radiative-convective and 3D climate 210 models for modeling steam atmospheres<sup>26</sup>. 211

As examples, we study H<sub>2</sub> atmospheres of 1 bar on a sub-Neptune planet that has a stellar irradiance similar 212 to Earth and orbits around an early M star similar to K2-18. A 1-bar H<sub>2</sub> atmosphere on such a planet would likely 213 have a surface temperature consistent with a liquid-water ocean (Fig. 5). For this planet, a 10-bar endowment of H<sub>2</sub> 214 would result in a massive steam atmosphere, making it uninhabitable<sup>26</sup>. We adopt the "water-world" interpretation 215 of the  $1.7-3.5~R_{\oplus}$  planet population that centers at  $10~M_{\oplus}$ , and  $2.5~R_{\oplus}^{15,21}$ , and assume 50% of water by mass in this study. In this interpretation, many sub-Neptunes may be ocean planets with deep oceans that do not require a 217 massive H<sub>2</sub> envelope to explain their radius, and can conceivably have moderate-size H<sub>2</sub> atmospheres. This may 218 not be directly applicable for K2-18 b, which resides on the low-density side of the  $1.7-3.5~R_{\oplus}$  population. The specific choices of these parameters are however unimportant, because atmospheric chemistry is not sensitive to 220 moderate changes in the surface gravity. 221

The mixing ratio of  $N_2$  on such a planet is probably set by atmospheric evolution (as opposed to the solubility equilibrium or geological recycling) and is assumed here to be 1%. As  $N_2$  only minimally participates in the chemical cycles and does not have strong spectral features in the infrared, its exact abundance is not our main

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concern. The photochemical model indicates that the NH<sub>3</sub> produced by photodissociation of N<sub>2</sub> in H<sub>2</sub> atmospheres has negligible mixing ratios ( $< 10^{-12}$ , see below).

CO<sub>2</sub> is the main form of carbon in thermochemical equilibrium with  $H_2O^{32,38}$ . If a liquid-water ocean exists, the partial pressure of  $CO_2$  is set by atmosphere-ocean partitioning, which in turn is mainly controlled by the oceanic pH<sup>39–42</sup>. The pH is affected by the abundance of cations in the ocean, which come from complex water-rock reactions and dissolution of the seafloor. The rates of the processes involved are uncertain; therefore, we explore the mixing ratio of  $CO_2$  from 400 ppm to 10%, corresponding to the pCO<sub>2</sub> range from the present-day Earth to early Earth<sup>62</sup> and including the predicted range for ocean planets<sup>41</sup> that is still consistent with an  $H_2$ -dominated atmosphere.

Is the 400-ppm  $CO_2$ , or  $4 \times 10^{-4}$  bar partial pressure in a 1-bar atmosphere, a reasonable lower bound of the  $CO_2$  partial pressure on a water world? We consider this question from a cosmochemical and geochemical perspective. Assuming equilibrium (during planet formation) between a Fe-core, a silicate mantle, and a wellmixed supercritical volatile envelope, the partitioning of C mass between reservoirs is described by

$$C_{\text{total}} = C_{\text{core}} + C_{\text{silicate}} + C_{\text{envelope}},$$
 (1)

where all reservoir masses are in kg, and

$$C_{core}/M_{core} = D_{C}(C_{silicate}/M_{silicate}),$$
 (2)

where  $D_C$  is a dimensionless partition coefficient,  $M_{core}$  (kg) is the mass of the Fe-dominated core, and  $M_{silicate}$  (kg) is the mass of the silicate mantle (molten during planet formation). For the partitioning between the envelope and the silicate mantle,

$$C_{\text{envelope}}k(g_{\text{esi}}/A_{\text{esi}})(\mu_{\text{avg}}/\mu_{\text{C}})s_{\text{C}} = C_{\text{silicate}}/M_{\text{silicate}},$$
 (3)

where k is a stochiometric correction from C mass to the mass of the C-bearing species in the envelope (i.e.,  $44/12 \sim 3.7$  for CO<sub>2</sub>),  $g_{esi}$  is gravitational acceleration at the envelope-silicate interface in m s<sup>-2</sup>,  $A_{esi}$  is the area of the envelope-silicate interface in m<sup>2</sup>,  $\mu_{avg}$  is the average molecular weight of the envelope (in Da),  $\mu_{C}$  is the

molecular weight of the C-bearing species (in Da), and  $s_{\rm C}$  is the solubility of the C-bearing species (in Pa<sup>-1</sup>).

Here we have assumed that the molten silicate layer is well-stirred.

Supposing  $M_{\rm core}/M_{\rm silicate} \sim 0.5$  (like Earth) and  $D_{\rm C} \sim 10^3$  f<sup>3</sup>, then  $C_{\rm core}/C_{\rm silicate} \sim 500$ . If  $C_{\rm silicate}/M_{\rm silicate} \sim 50$ 247 ppm then  $C_{\rm core}/M_{\rm core} \sim 2.5$  wt%, or  $C_{\rm total}/(M_{\rm core}+M_{\rm silicate}) \sim 1$  wt%, which is a reasonable lower bound for 248 the primordial carbon endowment (see below). For  $s_C = 0.55$  ppm/Mpa<sup>64</sup>, the envelope partial pressure of the C 249 species  $(=C_{\rm envelope}k(g_{\rm esi}/A_{\rm esi})(\mu_{\rm avg}/\mu_{\rm C}))$  is  $\sim 10^3$  bars. For a  $5-M_{\oplus}$  and  $1.5-R_{\oplus}$  core+mantle<sup>15</sup> that defines 250 the envelope-silicate boundary, and  $\mu_{\rm avg}/\mu_{\rm C}=0.4$  (appropriate for  ${\rm CO_2}$  in a  ${\rm H_2O}$ -dominated supercritical layer 251 during planet formation), the CO<sub>2</sub> mass in the envelope is 0.2% of an Earth mass. This estimate shows that even 252 though most C is in the core, still-significant reservoirs of C exist both in the silicate and in the envelope<sup>64–67</sup>. 253 Recent indications that the partition coefficient  $D_C$  is  $\ll 10^3$  at the pressures and temperatures that are relevant 254 for assembly of sub-Neptunes<sup>68</sup> would imply even more envelope C enrichment. 255

Following the formation of the liquid-water ocean, almost all of the  $CO_2$  will be dissolved in the ocean. For a  $5-M_{\oplus}$  water layer, the  $CO_2$  mass in the envelope estimated above corresponds to a concentration of  $\sim 0.01$  mol/L of dissolved  $CO_2$ . Here we have also assumed that the ocean is well-stirred. A higher oceanic pH leads to more effective dissolution and less  $CO_2$  in the atmosphere. As an extreme, if cations are leached from the silicate and not charge-balanced by chloride ions, then an ocean composition with a pH of 9-10 ("a soda lake") will result<sup>69</sup>. Using the equilibrium constant of carbonate and bicarbonate dissociation<sup>70</sup>, the  $CO_2$  partial pressure in equilibrium with this ocean would be  $5\times 10^{-5}\sim 7\times 10^{-4}$  bar, which is consistent with the assumed lower bound.

The partition coefficient gives the ratios of concentration of a species in the Fe-dominated core to the concentration of the same species in the silicate mantle. Therefore doubling the total amount of C in the core+mantle
will double the concentration in the magma. What is the whole-planet C content? In principle, a planet can form
without accreting volatiles. However, a thin-atmosphere sub-Neptune must have a thick volatile (H<sub>2</sub>O) layer in
order to match density data. It is very likely that a world that forms with 10s of wt% H<sub>2</sub>O will also accrete
abundant C. We develop this point in more detail in the following paragraph.

At  $T_{\rm eff} \sim 300$  K, the minimum liquid water content to explain most sub-Neptune masses and radii is  $> \sim 50$ 269 wt% even if there is no Fe-metal core<sup>16</sup>. This is more H<sub>2</sub>O than can possibly be produced by hydrogen-magma 270 reactions<sup>71</sup>, and instead implies a contribution of planet building blocks from the temperature range beyond the 271 water ice snowline. This is a zone where (in the Solar System), abundant refractory carbon is found. Specifically, 272 the carbon content of primitive chondrite meteorites (CI and CM type) is 2-6 wt%<sup>72</sup>. Although we do not fully understand the origin of this refractory carbon, proposed mechanisms for forming this refractory carbon would 274 also apply to exoplanetary systems<sup>73</sup>. Therefore we assume a planet bulk composition of  $(2-6 \text{ wt}\%) \times (1-x)$ 275 carbon, where x is the H<sub>2</sub>O mass fraction, and the remainder of the planet's building blocks are assumed to have 276 C content similar to that of primitive chondrites. This is a conservative lower limit on bulk C content for a 277 thin-H<sub>2</sub>-atmosphere sub-Neptune, for the following two reasons. (i) It considers only refractory C, not C ices 278 (e.g., CO<sub>2</sub> ice) which could be important in the case of whole-planet migration. (ii) Some primitive bodies in 279 the Solar System appear to be more C-rich than the most primitive chondrite meteorites; for example, the surface 280 of the dwarf planet Ceres may contain 20 wt% C<sup>74</sup>. These large bulk C contents map to substantial envelope C 281 contents (Equations 1-3). As such, the  $4 \times 10^{-4}$  bar partial pressure of CO<sub>2</sub>, while not the absolute lower limit, is 282 a cosmochemically and geologically reasonable lower bound of the CO<sub>2</sub> partial pressure on a water world. 283

The pressure at the water-rock boundary of a  $10-M_{\oplus}$  and  $2.5-R_{\oplus}$  planet is  $\sim 500$  GPa<sup>75,76</sup>, and this overloading pressure should suppress volcanism completely<sup>41,77,78</sup>. Therefore we do not include any volcanic outgassing in the standard models. As variant models, we consider the possibility of minor and intermittent sources of CO into the atmosphere. Evaporation of meteorites may provide a source of CO and  $CO_2^{79}$ , and water-rock reactions at the temperature relevant to the "fuzzy layer" may produce CO (and not CH<sub>4</sub> as it is thermochemically disfavored at high temperatures). The rates of these processes are unknown, but numerical experiments with the photochemical model indicate that an additional CO source of  $10^{10}$  molecule cm<sup>-2</sup> s<sup>-1</sup> would lead to a steady-state abundance of CO greater than that of H<sub>2</sub>, effectively resulting in a CO-dominated atmosphere. A CO source of  $10^{9}$  molecule cm<sup>-2</sup> s<sup>-1</sup> would produce the CO-dominated atmosphere in the 10%-CO<sub>2</sub> case but not in the 400ppm-CO<sub>2</sub> case. We therefore include a low-CO<sub>2</sub> case with the CO source of  $10^{9}$ 

molecule  $cm^{-2} s^{-1}$  as a variant model.

Table 1 summarizes the input parameters and results of the photochemical models, and Fig. 6 shows the 295 vertical profiles of main gases and photochemical products. CO is produced from the photodissociation of CO<sub>2</sub> 296 and can build up to the  $10^{-5}$  and  $10^{-2}$  mixing ratio level for the low-CO<sub>2</sub> and the high-CO<sub>2</sub> cases. OH from 297 the photodissociation of H<sub>2</sub>O destroys CO and maintains its steady-state mixing ratio. CH<sub>4</sub> is also produced 298 photochemically and can build up to a substantial mixing ratio ( $\sim 10^{-3}$ ) in the low-CO<sub>2</sub> case. The abundance 299 of CH<sub>4</sub> is much smaller in the high-CO<sub>2</sub> case, because of a more oxidized atmosphere. Together with the high 300  $CH_4$  mixing ratio,  $C_2H_6$  is produced in the low- $CO_2$  case and can accumulate to a mixing ratio of  $\sim 10^{-7}$ . 301 C<sub>2</sub>H<sub>2</sub>, as expected, is short-lived and only has significant mixing ratios in the upper atmosphere. Here we have 302 applied a deposition velocity of  $10^{-5}$  cm s<sup>-1</sup> to account for the loss of carbon due to organic haze formation and 303 deposition<sup>44</sup>; removing this sink does not substantially change the results shown in Fig. 6. The additional source 304 of CO would result in moderately more CO, CH<sub>4</sub>, and C<sub>2</sub>H<sub>6</sub> in the atmosphere (Model 1a in Table 1 and Fig. 6). 305 We see that CO, and in the low-CO2 case, CH<sub>4</sub>, can build up to the mixing ratio levels that cause significant features in the planet's transmission spectrum (Fig. 3). 307

Before closing this section, we address whether NH<sub>3</sub> can be produced substantially by water-rock reactions 308 and then emitted into the atmosphere. Hydrothermal systems on early Earth may produce NH<sub>3</sub> from the reduction 309 of nitrite and nitrate<sup>80,81</sup>. On a planet with an H<sub>2</sub>-dominated atmosphere, however, atmospheric production of the 310 oxidized nitrogen including nitrite and nitrate should be very limited. Moreover, the storage capability of NH<sub>3</sub> by 311 the ocean is vast and limits the emission into the atmosphere. At the pH value of 8 (a lower pH would further favor 312 the partitioning of NH<sub>3</sub> in the ocean),  $10^{-6}$  bar of atmospheric NH<sub>3</sub> requires a dissolved ammonium concentration 313 of  $10^{-3}$  mol/L in equilibrium<sup>70</sup>. The mass of NH<sub>3</sub> in the atmosphere and ocean is then  $\sim 10^{-5}$  of the planetary mass. This would only be possible if much of the planet's rocky core begins with a volatile composition similar to 315 carbonaceous chondrites, and most of this nitrogen is partitioned into the atmosphere and ocean as NH<sub>3</sub><sup>82</sup>, which 316 is highly unlikely as N<sub>2</sub> is thermochemically favored. Therefore, the concentration of dissolved NH<sub>3</sub> should be small and so is the atmospheric NH<sub>3</sub> on a planet with a massive ocean.

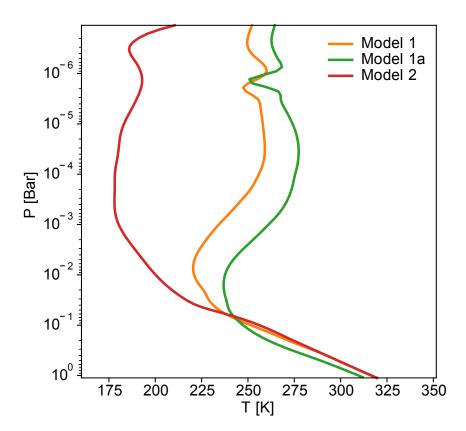


Figure 5: Modeled pressure-temperature profiles in a cold sub-Neptune like K2-18 b that has a small  $H_2$  atmosphere. For the stellar irradiance, we use  $S = S_{\rm Earth} * (1 - A_{\rm B})$ , with a Bond albedo of  $A_{\rm B} = 0.3$  (similar to Earth), to account for the radiative effects clouds would have in the otherwise cloud-free climate model. The surface albedo reflects a dark ocean (0.06) and convective adjustment follows a moist adiabat above the ocean. We show the results for the low- $CO_2$  case (Model 1 in Table 1), the low- $CO_2$  case with additional CO sources (Model 1a), and the high- $CO_2$  case (Model 2). The surface temperatures in these models are consistent with a liquid-water ocean.

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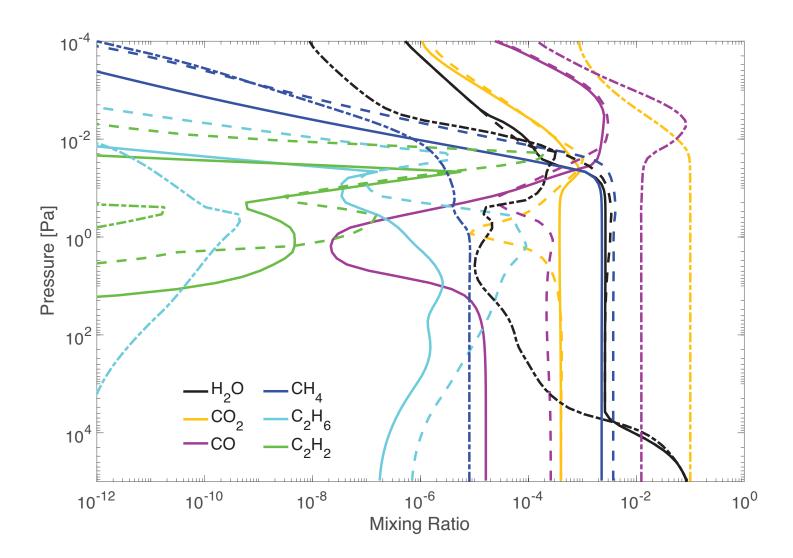


Figure 6: Modeled abundances of main gases and photochemical products in a cold sub-Neptune like K2-18 b that has a small  $H_2$  atmosphere. Solid, dashed, and dash-dot lines show the results for the low- $CO_2$  case (Model 1 in Table 1), the low- $CO_2$  case with additional CO sources (Model 1a), and the high- $CO_2$  case (Model 2). The photochemical abundance of nitrogen molecules such as  $NH_3$  and HCN is  $< 10^{-12}$ .

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- Author Contribution RH conceived and designed the study, simulated the photochemical models, interpreted the results,
  and wrote the manuscript. MD performed the JWST observation simulations and atmospheric retrievals. MS computed
  the pressure-temperature profiles. EK derived the cosmochemical and geological lower bounds for the carbon content. SS
  contributed interior structure models and insights. HR oversaw the development of the radiative-convective model used in
  the study. All authors commented on the overall narrative of the paper.
- 493 **Competing Interests** The authors declare that they have no competing financial interests.
- 494 **Correspondence** Correspondence and requests for materials should be addressed to Renyu Hu (email: renyu.hu@jpl.nasa.gov).

- Data Availability The raw data that are used to generate the figures in this paper are available from the corresponding
- author upon reasonable request.
- 497 Code Availability The photochemical code used in this study is written in C and will be released for unlimited re-use at
- 498 GitHub. The radiative-convective model used in the study is proprietary to Deutsches Zentrum für Luft- und Raumfahrt and
- requests for collaboration should be addressed to Heike Rauer (Heike.Rauer@dlr.de).
- 500 **Supplementary Information** Supplementary Information is available for this paper.

#### 501 1 Supplementary Information A: XXX

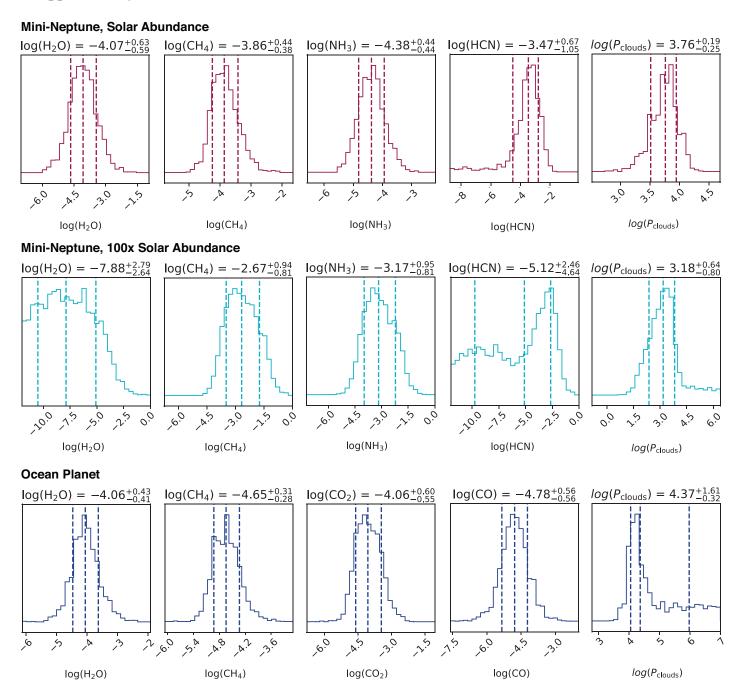


Figure S1: Retrieved posterior distributions of the abundances of the main chemical compounds and the cloud pressure in example massive-atmosphere and small-atmosphere scenarios. The input transmission spectra are shown in Fig. 3 and the expected uncertainties are calculated using PandExo<sup>48</sup>, assuming to observe two transits of K2-18 b with NIRSpec/G235H and four transits with NIRSpec/G395H. A detailed characterization of the atmosphere of K2-18 b, including distinguishing a small atmosphere versus a massive one and measuring the abundances of H<sub>2</sub>O, CH<sub>4</sub>, NH<sub>3</sub>, HCN, CO<sub>2</sub>, and CO, will be achievable with moderate time investment of JWST.

Model	Name	$CO_2$	CO flux	СО	CH <sub>4</sub>	$C_2H_6$
			$\mathrm{cm}^{-2}~\mathrm{s}^{-1}$			
1	Low-CO <sub>2</sub>	$4 \times 10^{-4}$	0	$1.6 \times 10^{-5}$	$2.3 \times 10^{-3}$	$2.2 \times 10^{-7}$
1a	Low-CO <sub>2</sub> Variant	$4 \times 10^{-4}$	$1.0 \times 10^{9}$	$2.6 \times 10^{-4}$	$3.7 \times 10^{-3}$	$9.0\times10^{-7}$
2	High-CO <sub>2</sub>	0.1	0	$1.2 \times 10^{-2}$	$8.0 \times 10^{-6}$	$< 10^{-12}$

Table 1: Summary of the photochemical model parameters and results. The volume mixing ratios of  $CO_2$  (as inputs), CO,  $CH_4$ , and  $C_2H_6$  (as results) are column-averaged and dimensionless.

### **Figures**

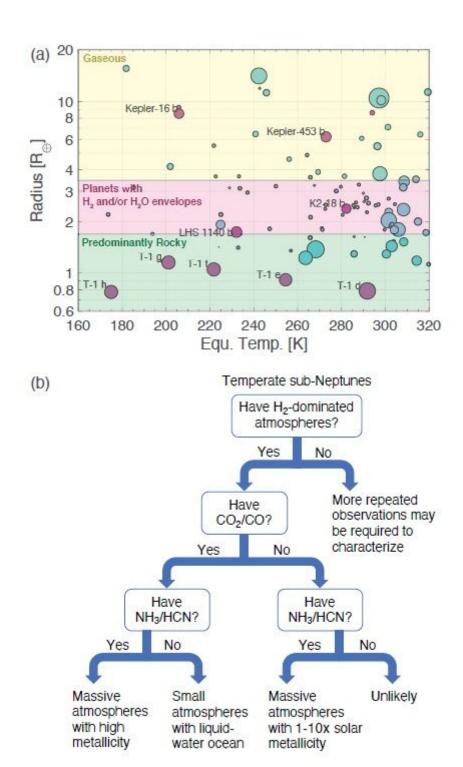


Figure 1

Temperate exoplanets amenable for atmospheric characterization via transmission spectroscopy. (a) Purple dots are confirmed planets with measured masses, and blue dots are planets with unknown masses or planet candidates. Data are taken from the NASA Exoplanet Archive and and the TESS Objects

of Interest Catalog. The marker sizes are scaled with the expected S/N of the spectral features of an H2 atmosphere observed by JWST at 2  $\mu$ m. Most of the temperate planets and planet candidates suitable for atmospheric characterization are larger than Earth and thus more likely to have H2 atmospheres. (b) A roadmap to characterize the mass of the atmospheres and the habitability of temperate sub-Neptunes by detecting signature gases. See text for details.

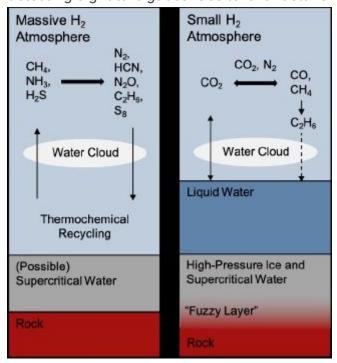


Figure 2

Interior structures of temperate H-rich exoplanets and the associated ranges of atmospheric composition. If the planet has a massive H2 atmosphere, the deep atmosphere would be hot – enabling thermochemical recycling – but a liquid-water surface would not be possible. If the planet has a small H2 atmosphere, a liquid-water surface may be possible. On these planets, the equilibrium abundance of atmospheric CO2 is set by the oceanic chemistry and that of N2 by atmospheric evolution.

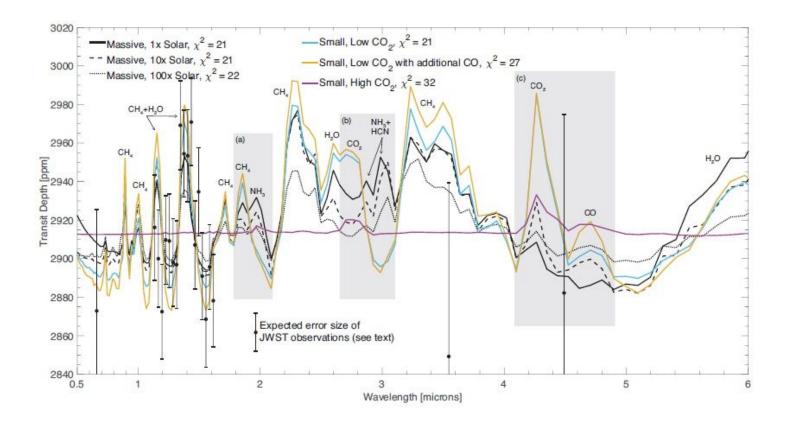


Figure 3

Expected transmission spectrum of temperate sub-Neptune planets of M dwarf stars, using K2-18 b as an example and comparing with the planet's transit depth observed by K2, Hubble and Spitzer7. The massive-H2-atmosphere models and the small-H2-atmosphere models differ in three spectral regions: in (a) and (b), the massive-atmosphere models have absorption features of NH3 and HCN, while the small-atmosphere models do not; in (c), the small-atmosphere models with a low mixing ratio of CO2 (400 ppm) have prominent features of CO2 and CO, while the massive-atmosphere models only have small features of CO2. The small-atmosphere models with a high mixing ratio of CO2 (10%) has a high mean molecular weight ( $\upalpha$  7) and a cloud top at  $\upalpha$  10–3 bar (as opposed to  $\upalpha$  3 × 10–2 bar the low-CO2 cases) and thus muted spectral features.

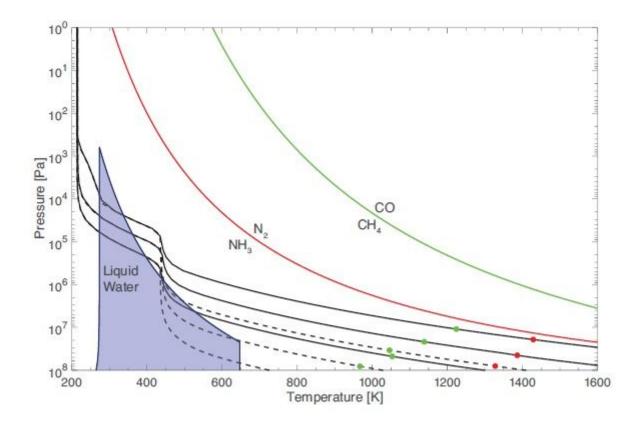


Figure 4

Modeled pressure-temperature profiles of a sub-Neptune exoplanet like K2-18 b with massive H2 atmospheres that receive stellar irradiance of Earth's, for an internal heat flux of Tint = 60 K (similar to Neptune, black solid lines) and Tint = 30 K (similar to Earth, black dashed lines). The pressure-temperature profiles assume 1×, 10×, and 100× solar metallicity. The shaded area shows the pressure and temperature range where pure liquid water is stable. The green and red lines show the equal-abundance boundaries for major carbon and nitrogen gases in a solar-metallicity gas in thermochemical equilibrium, and the green and red dots show the expected quench point for C0 and that for N2 respectively. The quench points are estimated by equating the chemical lifetime to the vertical mixing timescale57, 58, 60, with the eddy diffusion coefficient estimated using a mixing-length theory59 (№ 104 m2 s−1 at the pressure of 106 − 108 Pa). The quench point for C02 follows that of C0, and similarly, that of HCN occurs at a similar pressure and temperature as N257, 58.

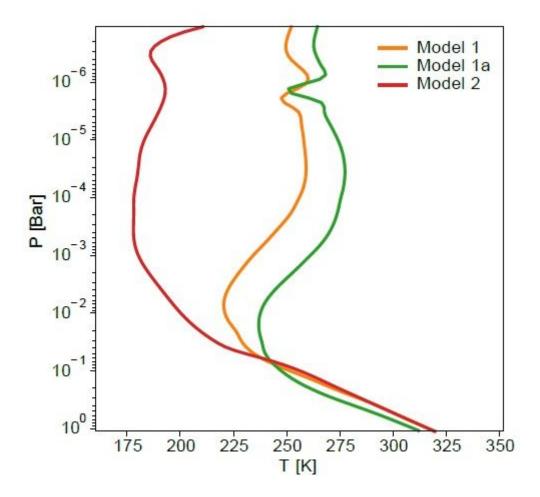


Figure 5

Modeled pressure-temperature profiles in a cold sub-Neptune like K2-18 b that has a small H2 atmosphere. For the stellar irradiance, we use  $S = SEarth \ \mathbb{I} \ (1 - AB)$ , with a Bond albedo of AB = 0.3 (similar to Earth), to account for the radiative effects clouds would have in the otherwise cloud-free climate model. The surface albedo reflects a dark ocean (0.06) and convective adjustment follows a moist adiabat above the ocean. We show the results for the low-C02 case (Model 1 in Table 1), the low-C02 case with additional C0 sources (Model 1a), and the high-C02 case (Model 2). The surface temperatures in these models are consistent with a liquid-water ocean.

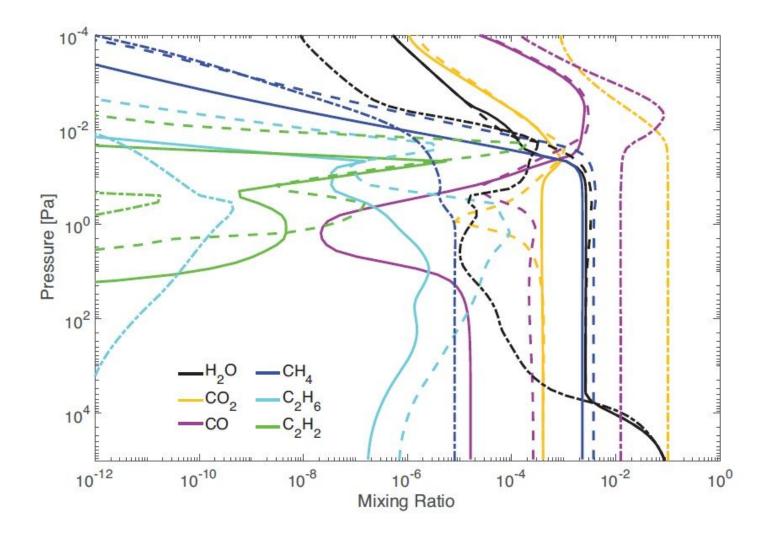


Figure 6

Modeled abundances of main gases and photochemical products in a cold sub-Neptune like K2-18 b that has a small H2 atmosphere. Solid, dashed, and dash-dot lines show the results for the low-C02 case (Model 1 in Table 1), the low-C02 case with additional C0 sources (Model 1a), and the high-C02 case (Model 2). The photochemical abundance of nitrogen molecules such as NH3 and HCN is < 10–12.