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Asthenospheric zircon below Galápagos dates plume activity

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Mantle plumes are active for long periods of time^{1,2}, however dating the onset of their activity is difficult. The magmatic products of the Galápagos plume, for example, have been subducted and fragmentarily accreted to the Caribbean and South American plates^{3,4}. Based on submarine and terrestrial exposures it is inferred that the plume has been operating for ~90 Myrs⁵ or perhaps even longer (e.g., ~139 Myrs⁶). Here we show that the activity of the plume dates back to ~170

31 **Ma. Evidence for this comes from 0 to 168 Ma zircon with isotopic plume signature**
32 **(Galápagos Plume Array; GPA) recovered from lavas and sediments from ten**
33 **islands of the archipelago. Given lithospheric plate motion, this result implies that**
34 **GPA zircon predating the Galápagos lithosphere (i.e., >14 Ma) formed at**
35 **asthenospheric depths. Thermo-mechanical numerical experiments of plume-**
36 **lithosphere interaction show that old zircon grains can be stored within local**
37 **asthenospheric stable domains to be later captured by subsequent rising plume**
38 **magmas. These results open new avenues for research on mantle plume dynamics in**
39 **similar tectonic settings.**

40 Global tomography and numerical models suggest that mantle plume occurrences are
41 closely linked to the margins of large low-shear velocity provinces (LLSVPs^{7,8}). In these
42 marginal zones, the ascent of material connects deep mantle dynamics with surface
43 processes through mantle plume activity. This will eventually form large igneous
44 provinces (LIPs), hotspot tracks and volcanoes⁹, like the modern Galápagos Archipelago
45 (Fig. 1), Hawai'i and Easter Islands. Recent studies suggest that despite striking
46 differences in the surficial expression of the Galápagos, Eastern and Hawai'i plumes, they
47 share a common generation mechanism originating at the Pacific LLSVP^{10,11}. Mantle
48 plume upwelling in the Pacific has been active since at least the mid-Jurassic, as recorded
49 in the Pigafetta Basin, which contains the oldest oceanic crust of the Pacific plate (~170–
50 160 Ma^{12,13}).

51 The Pacific plates exposed offshore Central America, Colombia-Ecuador, the Caribbean
52 and associated accreted onshore rock exposures contain a Mesozoic to recent record of
53 Pacific mantle upwelling events that resulted in new oceanic lithosphere, LIPs and
54 hotspot tracks¹⁴. The early plume products range from early Cretaceous to the last Pacific
55 LIP event: the Ecuadorian-Colombian-Caribbean LIP (ECCLIP) that formed mostly at
56 ~90 Ma, with potential additional events in the range 139–74 Ma^{2,5,6,15}. Though the
57 ECCLIP is generally considered to be a product of the Galápagos plume, other authors
58 however, based on paleomagnetic reconstructions, suggest that the ECCLIP originated
59 2000 km east of the Galápagos hotspot, and may thus not be derived from the same mantle
60 plume¹⁶.

61 The deep-rooted Galápagos mantle plume has generated several oceanic islands (<4 Ma
62 old¹⁷ and this study) on top of a young oceanic lithosphere (10 Ma in the northern part of the
63 Archipelago, 14 Ma in the southern part¹⁸ created at the nearby Galápagos Spreading

64 Center (GSC; Fig. 1). Galápagos volcanoes are fed by a mix of plume- and asthenosphere-
65 derived melts that provide important insight into heterogeneities of the mantle sources of
66 ocean island volcanism¹⁹. Isotopic and trace element compositions of basaltic lavas in the
67 Galápagos Archipelago indicate melting of several distinct mantle sources that include
68 components from recycled oceanic and continental crust materials^{20,21}. Palaeomagnetic
69 and geochemical data record a complex interaction between the hot spot and the GSC²².
70 The interplay between the mantle plume and the GSC dates back to Oligocene times (~23
71 Ma), when the Farallon plate tore into the Cocos and Nazca plates and the aseismic Cocos,
72 Carnegie and Malpelo ridges formed^{23,24}. Upon drifting away from the GSC, the oldest
73 parts of these hotspot tracks have been subducted at the Central America and Nazca
74 subduction zones (the oldest present-day ages of ridges are ~11–14 Ma^{25,26}).

75 We conducted an extensive sampling of basaltic and pumice rocks, inland deposits
76 collected on the floor of a lava tube and, of sands from uphill stream beds and beaches
77 (Figs. S1, S2, Fig. S3, Table S1) in 10 Galápagos islands (Figs. S1, S2; Table S1) covering
78 most of the main area of the Archipelago. Thirty-seven (37) zircon-bearing samples were
79 retrieved from the Central (Isabela 3 samples, Rábida 2, Santa Fe 5), Southern (Española
80 1, Floreana, 12), and Eastern (Baltra 2, Genovesa 1, Pinzón 2, San Cristobal 4, Santa Cruz
81 5) isotopic zones defined by²⁰. From these samples we analysed 238 zircon grains for U-
82 Pb dating and Hf and O isotopes compositions (see Supplementary Material for methods
83 and description of the zircon grains). Many samples (including magmatic rock samples;
84 Fig. 2a) contain a significant number of zircon crystals, most of which are <4 Ma old,
85 indicating direct crystallization from magmas erupted in the islands. Twenty of these
86 samples, however, contain zircon grains that significantly predate the spreading ridge-
87 and plume-related magmatic evolution of the Archipelago and associated lithosphere
88 (>14 Ma; Tables S2.1, S2.2).

89 The results of this study are displayed in Fig. 2 and show that, except for one grain (with
90 a 21.0 Ma core and a 18.5 Ma rim; Fig. S4.1-2), all zircon grains with ages up to 168 Ma
91 exhibit high positive $\epsilon\text{Hf}(t)$ (6–14; Tables S3.1, S3.2) and $\delta^{18}\text{O}_{(\text{zircon})}$ values well within
92 the range of mantle zircon²⁷ (4–6‰; Table S4). The zircon grains within this ~0–168 Ma
93 range are distributed continuously without age gaps, defining a “Galápagos plume array”
94 (GPA; Fig. 2a). Zircon younger than 0.2 Ma is rare (steep slope in Fig. 2b), given that
95 most recent lavas are scarcely exposed to erosion. On the contrary, zircon in the range
96 0.2–4 Ma is the most abundant (shallow slope in Fig. 2b), indicating that lavas of that age

97 have the highest exposure to erosion in the different islands. GPA zircon that has pre-
98 Galápagos ages in the range ~4–168 Ma is scarce (steep slope in Fig. 2b), but spreads
99 evenly and is isotopically indistinguishable from the younger zircon, suggesting the same
100 plume-related mantle origin.

101 The GPA trend is interrupted at 168 Ma (Fig. 2a) by the appearance of low $\epsilon\text{Hf}_{(t)}$ and high
102 $\delta^{18}\text{O}_{(\text{zircon})}$ values in zircon of Triassic age (213 Ma) and older. The non-GPA zircon form
103 a distinct array extending in ages from 213 to 3055 Ma with heterogenous $\epsilon\text{Hf}_{(t)}$ and
104 $\delta^{18}\text{O}_{(\text{zircon})}$. Whereas zircon in the range ~213–835 Ma (plus the outlier zircon grain with
105 a 21.0 Ma core and a 18.5 Ma rim) shows mostly low $\epsilon\text{Hf}_{(t)}$ (-27.7–1.8, only two show
106 higher values at 4.1 and 6.7) and generally high $\delta^{18}\text{O}_{(\text{zircon})}$ (4.7–10.8‰), typical of
107 continental crust, zircons in the range ~835–3055 Ma has a variety of $\epsilon\text{Hf}_{(t)}$ (-9.1–8.2)
108 and $\delta^{18}\text{O}_{(\text{zircon})}$ (4.7–11.3‰) values consistent with both juvenile and continental crust
109 signatures.

110 The non-GPA zircon grains indicate old (>213 Ma) to recent (~20 Ma, for the outlier)
111 external sources not related to plume activity. The absence of continental basement below
112 the Galápagos Archipelago^{28,29} rules out the possibility of continental crust provenance,
113 as could apply to other oceanic environments^{30,31,32}. Given the uncertainty about the
114 source of the exotic non-GPA zircons and because they are not direct magmatic products
115 of plume activity, their provenance exceeds the scope of this study and is not further
116 discussed here. Further discussions about transport mechanisms are nevertheless explored
117 in the Supplementary Material.

118 The most ground-breaking finding of our extensive zircon study is the group of GPA
119 zircons that pre-date the Galapagos lithosphere and with clear Hf and O isotopic mantle
120 signatures. Given that the age of the Galápagos Islands lavas exposed to erosion is <4
121 Ma (Fig. 2b), that the Galápagos lithosphere is as young as 10–14 Ma¹⁸ and plate motion
122 has removed any older lithosphere from above the plume head, any juvenile GPA zircon
123 older than 14 Ma must have formed in the asthenosphere and have been later picked up
124 by rising hot-spot magmas at asthenospheric depths (i.e., > ~50 km^{28,33}).

125 Two pre-Galápagos GPA zircon grains with ages of ~18 and 22 Ma are slightly younger
126 than the time when the Farallon plate was split by the GSC (just above the plume head)
127 and the Cocos and Carnegie ridges began to form (23 Ma³⁴). This suggests that ridge-
128 forming magmas did not fully escape from the plume head and crystallized zircon at

129 asthenospheric depths. The same reasoning can be extended to the other GPA zircon
130 grains older than 23 Ma. During this time, a number of magmatic events took place,
131 including the eruption of the Ecuadorian-Colombian-Caribbean LIP (ECCLIP) with a
132 major phase of LIP construction at ~ 90 Ma⁶. Notably, we sampled two zircon grains (93
133 and 94 Ma) formed close to this major event (Fig. 2b). The GPA zircon also includes ages
134 younger and older than the major ECCLIP event, clustering at early Tertiary (53, 55 and
135 65 Ma) and Jurassic (159, 164 and 168 Ma) times (Fig. 2b). The latter would allow
136 expanding 30 Myrs back in time the magmatic history of the Galápagos plume recorded
137 in its accreted dispersed fragments⁶. These GPA zircon data indicate that a) zircon is
138 magmatic and juvenile, b) zircon formed much earlier than the recent lavas that brought
139 them to the surface and formed the present-day Galápagos islands, and c) zircon formed
140 in the asthenosphere and was stored at depth while staying unaffected by other plume-
141 related magmatic events throughout the last ~ 170 Ma. These juvenile GPA zircon grains
142 hence offer a unique opportunity to date the evolution of the mantle plume and to evaluate
143 plume dynamics and asthenospheric flow.

144 A first-order observation is that our data places the onset of the Galápagos plume to at
145 least ~ 170 Ma, much earlier than previously thought. We can, however, discard
146 contamination. Even if most GPA zircon was sampled from surficial detritus, one of the
147 oldest GPA zircon grains (i.e., 164 Ma), was sampled from a basaltic lava at the Alcedo
148 Volcano on the uninhabited Isabela Island (Figs. 2b and S1). We can then rule out that
149 this and other GPA zircons sampled from detrital material were delivered to the islands
150 by surficial or anthropogenic processes. In addition, the presence of post- and pre-14 Ma
151 GPA zircons in sands from almost virgin beaches and uphill streams (in the Baltra,
152 Floreana, San Cristobal, Santa Cruz and Santa Fé islands) and inland lava tube deposits
153 (Santa Cruz) clearly point to a local provenance from erosion of exposed volcanic rocks.
154 Furthermore, a provenance study carried out on beaches from eleven islands of the
155 archipelago shows that mineral grains and rock fragments derive from locally exposed
156 basaltic rocks and excludes external sources³⁵. All lines of evidence thus point to the
157 crystallisation of 14–168 Ma GPA zircons in the sub-lithospheric source of Galápagos
158 lavas.

159 Zircon forms after significant fractional crystallisation of primitive Zr-subsaturated
160 basaltic magmas e.g.,^{36,37} and is stable down to ~ 300 km depths, below which it transforms
161 into reidite³⁸. Experimental work shows that zircon can survive in the presence of mafic

162 melt for long periods of time as long as the volume of melt that interacts with a zircon
163 crystal is small³⁹. It is possible then that GPA zircon crystallised from near-solidus Zr-
164 saturated evolved basaltic liquids at plume-head regions with limited melt fraction. Once
165 formed, GPA zircon survived in hot mantle, even more so if the crystals are shielded
166 within other mineral grains. If shielded within a Pb-free mineral (e.g., olivine), zircon
167 grains can retain their U-Pb crystallisation ages even at 1500 °C, independently of their
168 residence time in the mantle⁴⁰. Eventually, rapidly ascending magmas may pick up these
169 zircon grains or zircon-bearing mineral or rock fragments. At this stage, dissolution of
170 zircon occurred if not shielded and/or the magma resided long in a magma chamber.
171 Ultimately, however, some asthenospheric zircon grains indeed survived and reached the
172 surface in the crystallising magmas that, in turn, eventually reached Zr-saturation and
173 formed new zircon, as demonstrated by many young zircon grains that date the onset of
174 the present “IOB” Galapagos vulcanism at 4 Ma (Fig. 2b).

175 Our finding of old asthenospheric mantle zircon grains challenges current ideas about
176 asthenosphere convection and plume/lithosphere interaction. Contrary to expectations
177 from current understanding of asthenospheric motion, our finding implies that zircon was
178 not dispersed by convective flow even after more than 100 Mys residence time in the
179 asthenosphere. To explore whether this exceptional behaviour of the asthenosphere and
180 associated processes are physically feasible, we carried out numerical thermo-mechanical
181 simulations. Our reference scenario (Fig. 3) features a plume with a radius of 150 km and
182 an excess temperature of 250 °C with respect to the local mantle potential temperature,
183 consistent with recent estimates from seismic tomography⁴¹. The plate has a constant
184 velocity of 5 cm/yr, and an initial thermal age of 30 Ma (see Methods section for further
185 details). After an initial rising stage, the plume head is sheared along with the moving
186 plate, and splits into smaller plumes that produce partial melt beneath the lithosphere
187 (Fig.3a). Zircons will form in these partially molten regions, which are coloured green or
188 violet (depending on whether they consist of plume or ambient mantle material) to track
189 their subsequent location. We initially introduce 90,000 passive tracers that represent
190 zircon. The tracers are placed in the upper mantle and we track their motion once the
191 mantle partially melts for the first time. Their subsequent path is tracked until they arrive
192 in a partially molten region for the second time, when they are assumed to be extracted
193 to the surface along with freshly erupting lavas. The zircon age is accordingly the age
194 between formation and eruption (in Ma). Zircon that is transported to depths >300 km has

195 its U-Pb age reset upon reidite formation^{e.g.,42} and is no longer considered in the
196 interpretation.

197 The model results show that zircon can stay in the shallow upper mantle for extended
198 periods of time (Fig. 3). Counterintuitively, not all zircon is dragged along with the
199 moving lithosphere, but much instead initially moves in the opposite direction. This is
200 because the plume's rising velocity is larger than the plate motion, which induces small
201 scale convective cells (Fig. 3b-c). Some of this zircon returns to the plume area, whilst
202 the rest is mixed in with the asthenospheric upper mantle. The dynamics of the plume is
203 cyclic, with periods of slow and steady activity interrupted by more active phases. This
204 results in discontinuous magmatic activity and sub-lithospheric circulation patterns that
205 are mostly confined to the uppermost asthenospheric mantle, allowing the preservation
206 of zircons. Most passive tracers are erupted to the surface within 50 million years of
207 plume formation and contain young zircon (Fig. 3d and 3e). Yet, old zircons can be
208 dragged into these lavas and erupt as well, in accordance with our observations in the
209 Galápagos Archipelago.

210 In order to test the sensitivity of our results to changes in the model parameters, we
211 performed over 20 simulations that demonstrate that the features described above are
212 reproducible. The systematic analysis shows, however, that differences in the model
213 results arise as a function of plume radius, temperature and oceanic plate age. Old oceanic
214 plates (40 Ma) reduces the amount of melt produced, but can preserve very old zircons
215 (~130 Ma), while young plates (15 Ma) generally give rise to a smaller amount of old
216 zircons (see Fig. S5, S6 and Fig. S7). A smaller plume radius (<100 km) produces less
217 melt, while higher plume temperatures result in a fully molten layer below the lithosphere
218 with corresponding young zircon ages. Plate velocity and the manner in which a plume is
219 introduced in the models are of second order importance.

220 Our results thus suggest that, once formed in a plume head, zircon crystals can remain
221 within the asthenospheric mantle for extended periods of time. Following these results,
222 the recorded asthenospheric zircon ages hence allow dating the Galápagos plume back to
223 Jurassic times, a much older age than previously reported. Similar zircon observations
224 and models of asthenospheric flow below ocean islands could apply to other plume-
225 related hot spots. Therefore, a systematic analysis of zircon from ocean islands will allow
226 monitoring temporal ranges of plume activity and dynamics over much longer periods

227 than those implied by the ages of the erupted lavas, hotspot tracks, plateaus and,
228 eventually, plume-related terranes accreted to active continental margins.

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371

372 **Methods**

373 To avoid potential contamination from mineral separation equipment, the sand and soil
374 samples were panned to yield heavy mineral concentrates on Galápagos beaches^{43,44}. The
375 basalt samples were processed at Mainz University, where extreme care was taken to
376 avoid mineral contamination. Approximately 1–5 kg of sample was crushed to ~ 250 µm
377 grain size using a jaw crusher and a roller mill.

378 The heavy mineral fraction was obtained using the panning technique and a Frantz
379 magnetic separator at Mainz University followed by final panning with alcohol in the
380 Beijing SHRIMP Centre, China. Zircons for isotopic analysis were handpicked under a
381 binocular microscope and mounted in epoxy resin. About 100 g of the homogenized rock
382 material were pulverized in a tungsten carbide mill (Siebtechnik) for chemical and whole-
383 rock isotopic analysis. Zircon was analysed in situ for U-Pb and O isotopes using the
384 SHRIMP II at the Beijing SHRIMP Centre, Chinese Academy of Geological Sciences,
385 and for Hf isotopes using the LA-MC-ICP-MS at Hong Kong and Frankfurt Universities
386 (Supplementary Tables S2.1, S2.2, S3.1, S3.2 and S4). The analyses were guided by
387 optical and cathodoluminescence (CL) images (Supplementary Figures S4.1, S4.2).
388 Analytical procedures are presented in detail in the Supplementary Material.

389 **Geodynamic modelling approach**

390 The numerical experiments herein showed have been performed using LaMEM^{45,46}, a
 391 visco-elasto-plastic finite difference code. LaMEM solves the fundamental conservation
 392 equation of mass (1), momentum (2) and energy (3) assuming the rocks to be
 393 incompressible:

$$\frac{\partial v_i}{\partial x_i} = 0 \quad \underline{1}$$

$$\frac{\partial \tau_{ij}}{\partial x_j} - \frac{\partial P}{\partial x_i} + \rho g = 0 \quad \underline{2}$$

$$\rho C_p \frac{DT}{Dt} = \frac{\partial}{\partial x_i} \left(k \frac{\partial T}{\partial x_i} \right) + H_a + H_s + H_r \quad \underline{3}$$

394 Where v_i are the component of the velocities along x_i direction, τ_{ij} are the component of
 395 the deviatoric stress tensor, P is the pressure, g is the gravity acceleration and ρ is the
 396 density. DT/Dt is the substantial time derivative of the temperature, C_p is the heat
 397 capacity, k thermal conductivity, H_a , H_s and H_r are the adiabatic, shear and radiogenic
 398 heat sources.

399 The conservation equations are solved in a fixed Eulerian frame of reference, using a
 400 finite difference staggered grid scheme, while the advection is explicit with time and
 401 performed using lagrangian particles and a second order Runga Kutta scheme. The
 402 lagrangian particles carries all the historical information needed to solve the equations.

403 **Numerical Design:** The main goal of our simulations is to understand if plume-
 404 lithosphere interaction is able to generate chemical distinct mantle domains that preserves
 405 geochronological information for a sufficient long period of time (i.e. >80 Myrs).
 406 Furthermore, we want to assess if some of these chemical heterogeneities can be captured
 407 again by the mantle plume flux during later model stage. For simplification, we do not
 408 consider the effects of radiogenic, adiabatic, and shear heating and assuming only a
 409 thermal dependent density and viscous constitutive models. We employ lateral inflow-
 410 outflow boundaries condition to simulate the motion of the oceanic plate, as well as a
 411 plume inflow boundary condition. To identify the potential domains that undergoes to
 412 chemical refinement, we interpolate the volumetric melt fraction (ϕ) from a precomputed
 413 mantle petrological phase diagram and track them using passive tracers.

414 Density depends on temperature and volumetric melt fraction:

$$\rho_{solid} = \rho_0(1 - \alpha(T - T_{ref})) \quad 4$$

$$\rho_{eff} = \rho_{solid}(1 - \phi) + \rho_{melt}\phi \quad 5$$

415 Where ρ_{eff} is the effective density, ρ_{solid} is the solid density, ρ_{melt} is the melt density, α is
 416 the thermal expansion, and T is the actual temperature and T_{ref} is the reference state
 417 temperature. We assume that melt feedbacks on density and viscosity reach their
 418 maximum effects at $\phi=0.08$ (see eq. 5) as we do not explicitly account for melt extraction
 419 processes. The mantle phase diagram has been computed using `Perple_X`⁴⁷ using the
 420 pyrolite composition from⁴⁸, and using the solution model of⁴⁹ – for further references
 421 see⁴⁶. All the relevant properties of each of the compositional phase are listed in
 422 Supplementary Table S6.

423 We employ a purely viscous constitutive model in our numerical simulation, using both
 424 linear (diffusion-creep) and non-linear (dislocation-creep) relations. The viscosity of the
 425 material has a lower and upper cutoff and the effective viscosity is the harmonic average
 426 between diffusion and dislocation creep and the cut-offs:

$$\eta_{tot} = \frac{1}{\eta_{diff}} + \frac{1}{\eta_{disl}} + \frac{1}{\eta_{upper}} \quad 6$$

427

428 Where η_{tot} , η_{diff} , η_{disl} , and η_{upper} are the total effective, diffusion creep, dislocation creep,
 429 and upper cut-off viscosities respectively, whereas the minimum cut-off is introduced by
 430 adding a minimum stress in parallel to the viscoelastoplastic stress rheology. Diffusion
 431 and dislocation creep are computed using the law described in⁵⁰:

$$\eta_{diff} = \frac{1}{2} B_{diff} \exp(-\chi\phi) \exp\left(-\frac{E_{act} + PV_{act}}{RT}\right) \quad 7$$

$$\eta_{disl} = \frac{1}{2} B_{disl}^{\frac{1}{n}} \varepsilon_{II}^{\frac{1}{n}-1} \exp\left(-\frac{\chi\phi}{n}\right) \exp\left(-\frac{E_{act} + PV_{act}}{nRT}\right) \quad 8$$

432

433 B_{diff} and B_{disl} are the pre-exponential factor of diffusion and dislocation creep respectively.
 434 α is the pre-exponential factor to simulate the melt viscous weakening, n is the stress
 435 exponent, ε_{II} is the second invariant of the strain rate tensor, E_{act} and V_{act} are the activation
 436 energy and volume respectively (see Supplementary Table S6).

437 ***Initial setup & Boundary condition:*** We perform 2D numerical experiments using a
438 domain that extends 4000 km x 1000 km along x and z direction and with a grid resolution
439 of 256 x 128 elements along x and z direction. This strategy allows to perform long-term
440 experiments able to cover the time evolution recorded by the geochemical data here
441 presented. Moreover, employing such large numerical domains prevents inflow boundary
442 condition to interfere with the processes that we simulate. The mantle is initially
443 isothermal (1350 °C) using the half-space cooling model to describe the initial
444 lithospheric thermal structure. The uppermost portion of the compositional field is
445 composed by a thin oceanic crust (10 km), and 90 km of lithospheric mantle, while the
446 rest of the domain is filled with mantle phases (i.e. upper and lower mantle). The thickness
447 of the lithospheric mantle is self-adjusting during the evolution of the simulation as a
448 function of the 1200 °C isotherm position. A depth dependent post-spinel phase transition
449 at 660 km is introduced to adjust the inflow velocity of the plume allowing it to have a
450 smooth temperature profile (the density jump associated with this phase transition is $\Delta\rho$
451 = 300 kg/m³ see Supplementary Table S6). The initial plume is located at the bottom of
452 the numerical domain and in its centre. In most numerical experiments (Supplementary
453 Table S7), we introduce a rectangular thermal perturbation at the bottom to simulate an
454 initial plume conduit to trigger the upwelling as soon as the simulation starts, which has
455 the same phases as the plume and an excess temperature of 250 °C. Its width is equal to
456 the inflow diameter (i.e. 300 km) and covers almost all the lower mantle with its height.
457 Since, the zircon that have been collected from Galápagos are older than the Nazca plate,
458 we assume that the pacific plate retains its integrity for the whole duration of the
459 simulation. The initial age of the plate, and the temperature of the inflow boundary
460 condition is varied from 15 Myrs to 40 Myrs and the plate velocity is varied from 1-10
461 cm/yrs.

462 We employ a free slip boundary condition at the upper boundary with a constant surface
463 temperature of 0 °C. Lateral boundaries are no-heat flux and free slip boundary condition
464 except for a narrow inflow-outflow window. The inflow-outflow window extends from
465 the top of the domain to a minimum depth of 100 km to a maximum of 350 km. In all
466 numerical experiments the velocity between 0 to 100 km along z -direction is constant,
467 and it directed from left to the right side of the numerical domain. Then in most of the
468 numerical experiments we introduce a buffer inflow-outflow window in which the
469 velocity is linearly decreased to 0 km (the relax distance is varied from 0 to 250 km). The

470 inflow plate has a constant thermal age that increases towards the right as a function of
471 the plate velocity. The bottom boundary is permeable with no tangential velocity
472 components (following⁵¹), and the temperature at the boundary is constant and equal to
473 the ambient mantle temperature (1350 °C) with a gaussian thermal perturbation with a
474 radius of 150 km and a $\Delta T = 250$ °C. Within the plume inflow window (-150, 150 km
475 along x direction), particles with plume phase are injected with a temperature equal to the
476 one of the bottom boundary condition. Outside this inflow window the temperature of the
477 material is assumed to be equal to the ambient mantle temperature and has the same phase
478 of the normal mantle.

479 ***Mantle chemical heterogeneities:*** In order to track the mantle chemical heterogeneities,
480 we use two strategies: first, we highlight areas of the mantle or plume that reaches $\phi=0.05$
481 at least once (i.e. we change the visualization phase, see **Fig. 3**); second, we activate
482 passive tracers ('passive tracers zircon' in **Fig. 3**) that are associated with a chemical
483 heterogeneous mantle domain and start tracking their position, temperature, pressure and
484 melt quantity. The initial position of the passive tracers is defined by a refined grid
485 spanning from -1200 to 1200 km and from 100 to 200 km along x and z direction
486 respectively. As soon as they are activated, we record the age of the melting event,
487 assuming that this portion of mantle could bear geo-chronological information (see Fig.
488 3). These mantle domains have to be interpreted as potential portions of the mantle that
489 can bear zircon. Once the passive tracers melt again (when $\phi>0.05$), we assume that the
490 rising melts will bring the zircons to the surface and remove them from the model domain
491 ('pre-eruption zircons' in **Fig. 3**). When the passive tracer is erupted, we collect
492 information about the zircon age, and the timing at which the eruption event occurs (see
493 **Fig.3**). The age of the potential zircon generated during the partially melting of the mantle
494 is measured in Ma, while the actual simulation time is expressed in Myrs (as for the initial
495 thermal age of the lithosphere). The chemical heterogeneous mantle domains that are
496 tracked give an estimation of the age of the event that may lead to the generation of zircon
497 population, and represents an estimation of the maximum age retained by them.

498 **Reference Methods**

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534 **Author contributions** Y.R-A and A.G-C developed the original idea for the study, Y.R-
535 A, I.W. H.X, A.G, J.W, S.B, dated the samples, Y.R-A and P.M. performed the data
536 reduction of young zircon grains, B.K and A.P performed the geodynamic simulations.
537 All authors including Th.T contributed to discussions and writing of the manuscript.

538 **Competing interests** The authors declare no competing interests.

539 **Supplementary information** the online version contains supplementary material
540 available at

541 **Figure captions**

542 **Figure 1: Plate tectonic configuration** | Plate-tectonic configuration of the Pacific and
543 the Caribbean region and location of uplifted and accreted ECCLIP (Ecuadorian-
544 Colombian-Caribbean-Large-Igneous-Province) fragments. See supplementary Figure
545 S1 for sample locations containing zircon grains.

546 **Figure 2: Isotopic composition ($\epsilon\text{Hf}_{(t)}$ and $\delta^{18}\text{O}_{(\text{zircon})}$) and age of Galápagos zircons** |

547 A) U-Pb age vs $\epsilon\text{Hf}_{(t)}$ (blue) and $\delta^{18}\text{O}_{(\text{zircon})}$ (red) of analysed zircons. Both high $\epsilon\text{Hf}_{(t)}$
548 and low $\delta^{18}\text{O}_{(\text{zircon})}$ define the Galápagos Plume Array (GPA, blue and red rectangles,
549 respectively), which extends from 0 Ma to ca. 170 Ma (note significant scatter at >170
550 Ma). B) Age of analysed zircons sorted by age of spot. The distribution shows four
551 sectors separated by slope breaks, including 1: age range of zircon of less abundant
552 (most recent) volcanic rocks exposed to erosion (<0.2 Ma); 2: age range of zircon of
553 most abundant volcanic rocks exposed to erosion (0.2-4 Ma); 3 and 4: age ranges of
554 pre-Galápagos Islands zircon, comprising 3: zircons belonging to the Galápagos Plume
555 Array that extend from 4 to 168 Ma well beyond the oldest age of the exposed lavas
556 and the age of basement oceanic lithosphere (10-14 Ma), and 4: exotic zircons older
557 than 168 Ma with scattered $\epsilon\text{Hf}_{(t)}$ and $\delta^{18}\text{O}$, including continental crust signature. For
558 reference, 23 Ma (estimated split of Farallon plate) and 90 Ma (major ECCLIP event)
559 are indicated in B. In A and B, lighter coloured circles for $\epsilon\text{Hf}_{(t)}$ and age data,
560 respectively, correspond to zircon grains separated from rock samples.

561 **Figure 3: Numerical models of plume-lithosphere interaction** | Snapshots **a-c** show
562 the dynamics of a plume (with 1600 °C) interacting with a 30 Myrs old oceanic
563 lithosphere that moves with 5 cm/yr to the right. Partially molten regions (indicated
564 by purple shaded areas produce chemically distinct mantle domains (green/violet). We

565 additionally highlight the chemical heterogeneous mantle domain that undergoes the
566 zircon-reidite phase transition. The pathway of zircons in the mantle is tracked by
567 passive tracer zircons (circles) which are coloured by age until a zircon arrives at a
568 partially molten region for the second time when they are assumed are removed from
569 the numerical domain (indicated by stars). Both the age (**d**) and amount (**e**) of erupted
570 passive tracer zircons are tracked throughout the simulation and show that old zircons
571 can be preserved in the shallow mantle for extended periods of time.

572

Figures

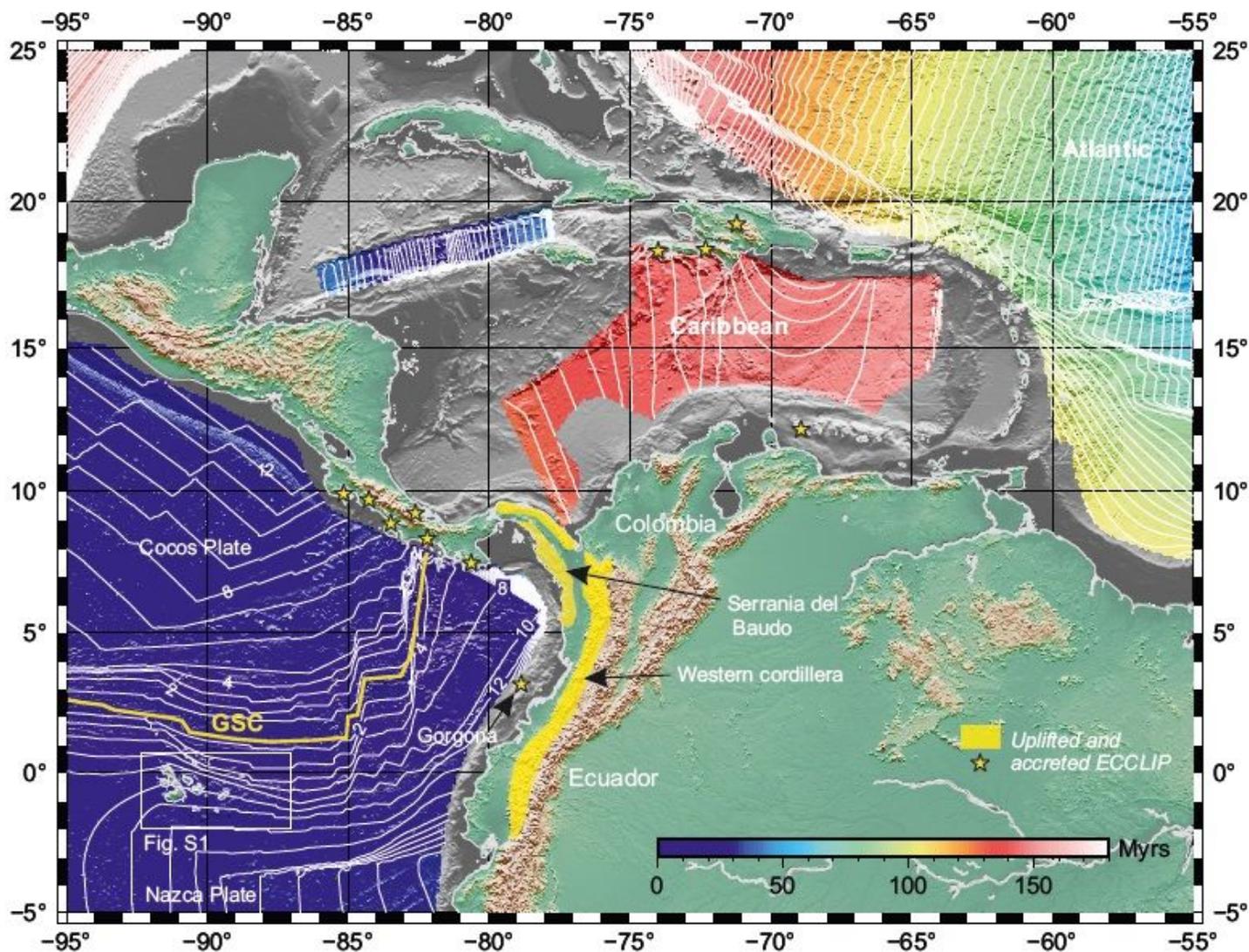


Figure 1

Plate tectonic configuration | Plate-tectonic configuration of the Pacific and the Caribbean region and location of uplifted and accreted ECCLIP (Ecuadorian- Colombian-Caribbean-Large-Igneous-Province) fragments. See supplementary Figure S1 for sample locations containing zircon grains.

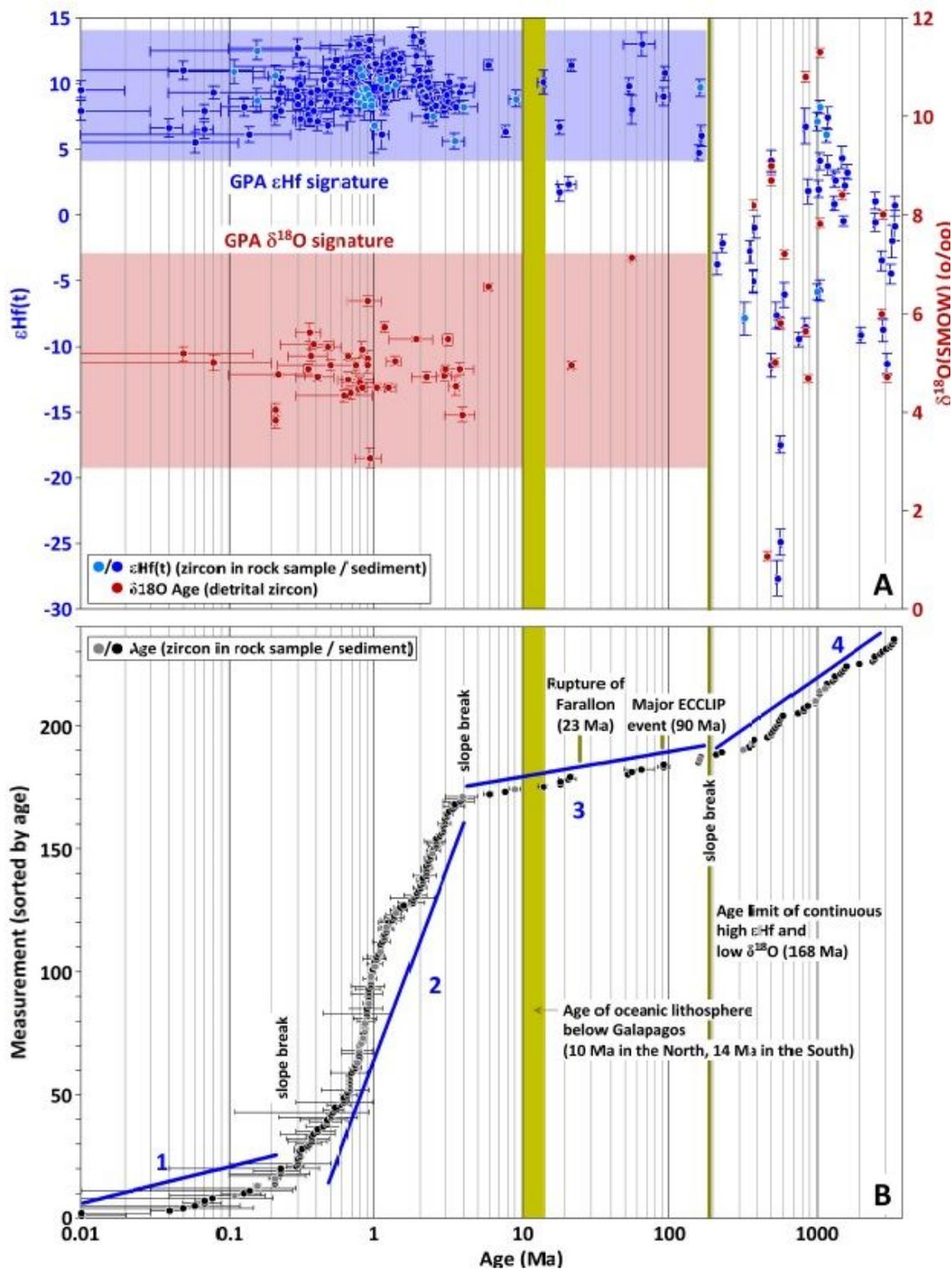


Figure 2

Isotopic composition ($\epsilon\text{Hf}(t)$ and $\delta^{18}\text{O}(\text{zircon})$) and age of Galapagos zircons | A) U-Pb age vs $\epsilon\text{Hf}(t)$ (blue) and $\delta^{18}\text{O}(\text{zircon})$ (red) of analysed zircons. Both high $\epsilon\text{Hf}(t)$ and low $\delta^{18}\text{O}(\text{zircon})$ define the Galapagos Plume Array (GPA, blue and red rectangles, respectively), which extends from 0 Ma to ca. 170 Ma (note significant scatter at >170 Ma). B) Age of analysed zircons sorted by age of spot. The distribution shows four sectors separated by slope breaks, including 1: age range of zircon of less

abundant (most recent) volcanic rocks exposed to erosion (<0.2 Ma); 2: age range of zircon of most abundant volcanic rocks exposed to erosion (0.2-4 Ma); 3 and 4: age ranges of pre-Galápagos Islands zircon, comprising 3: zircons belonging to the Galápagos Plume Array that extend from 4 to 168 Ma well beyond the oldest age of the exposed lavas and the age of basement oceanic lithosphere (10-14 Ma), and 4: exotic zircons older than 168 Ma with scattered $\epsilon\text{Hf}(t)$ and $\delta^{18}\text{O}$, including continental crust signature. For reference, 23 Ma (estimated split of Farallon plate) and 90 Ma (major ECCLIP event) are indicated in B. In A and B, lighter coloured circles for $\epsilon\text{Hf}(t)$ and age data, respectively, correspond to zircon grains separated from rock samples.

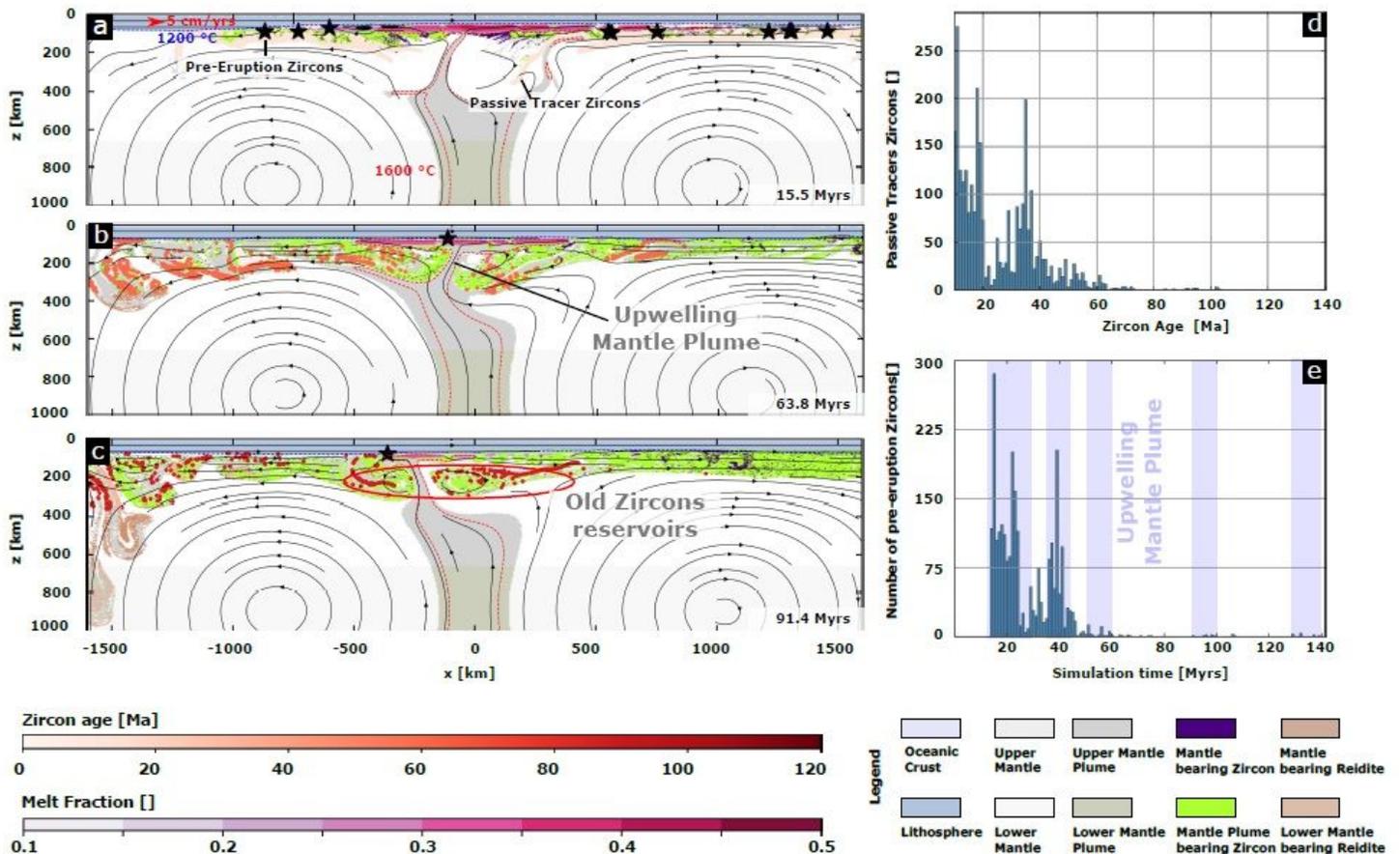


Figure 3

Numerical models of plume-lithosphere interaction | Snapshots a-c show the dynamics of a plume (with 1600 °C) interacting with a 30 Myrs old oceanic lithosphere that moves with 5 cm/yr to the right. Partially molten regions (indicated by purple shaded areas) produce chemically distinct mantle domains (green/violet). We additionally highlight the chemical heterogeneous mantle domain that undergoes the zircon-reidite phase transition. The pathway of zircons in the mantle is tracked by passive tracer zircons (circles) which are coloured by age until a zircon arrives at a partially molten region for the second time when they are assumed are removed from the numerical domain (indicated by stars). Both the age (d) and amount (e) of erupted passive tracer zircons are tracked throughout the simulation and show that old zircons can be preserved in the shallow mantle for extended periods of time.

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