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## Research Article

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# Nonlinear dynamic analysis of variable lead preloaded single nut ball screw considering the variation of working parameters

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## Abstract:

In this study, a 5-degrees-of-freedom dynamic model of variable lead preloaded single nut ball screw is proposed. The proposed model considers the effect of excitation amplitude, deflection angle, the number of balls, and the preload into consideration. Moreover, to study the effect of working parameters on nonlinear dynamics, bifurcation diagram, 3-D frequency spectrum, phase diagram and Poincaré section with different system parameters are shown in discussion section. The numerical analysis indicates that the system can exhibit different motion states, and the continuous frequency component can be observed. Moreover, a series of experiments are conducted to estimate the dynamic parameters and validate the proposed dynamic model.

**Keywords:** Single nut ball screw; Nonlinear dynamics; Working parameters;

## 1. Introduction

Ball screws have been used in missiles and aircraft to move control surface [1]. It is also used in machine tools, robots, and many other precision assembly equipment due to its high efficiency, great stiffness, and long service life [2]. The positioning accuracy of ball screw feed system is usually influenced by preload [3], thermal deformation [4], and health condition of each kinematic joint [5]. During the positioning process, the precision of the feed system can be strongly affected by working condition, manufacturing accuracy, and assembly quality. Due to the nonlinear Hertzian contact stress, piecewise restoring force, and deflection angle, the system can exhibit wide variety of nonlinear dynamics. However, the influence of assembly error on nonlinear dynamic characteristic has been discussed by few papers. In addition, the preload can be changed with the increase of service life. Therefore, it is necessary to study the nonlinear behavior of ball screw feed system to avoid the reduction of positioning accuracy under the complex working conditions. The dynamic characteristics of ball screw has been studied by many researchers under different conditions.

Liu et al. [6] proposed a 9-degrees-of-freedom dynamic model of ball screw feed system, the effect of assembly error on the nonlinear dynamics of the system was discussed, and the screw nut in the proposed model had 3-degrees-of-freedom. However, the screw nut was also subjected to torsional moment, the moment should not be ignored. Guo and Yi [7] investigated the relationship between the detected vibration signals and the variation of preload, and developed a lumped dynamic model to study the preload variation of the ball screw feed system. Bo et al. [8] developed a mathematical model to predict the structural nature frequency influenced by moving component. Considering the influences of dimension errors of balls in screw nut, Ni and Qi [9] proposed a calculation model to analyze the effect of ball's dimension error on mechanical properties of a ball screw drive system, and the fatigue life influenced by ball's machining accuracy was studied. Chen and Tang [10] found that the force of

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each ball was not same, and developed a stiffness transition matrices for the ball-screw mechanism to study the influence of the contact stiffness. Due to the large friction force, the contact state of each component could change, and hence influenced the dynamic characteristics of the ball screw feed system. Using hybrid element method, Zhang et al. [11] theoretically investigated the relationship between ball screw feed system natural frequency and feed rates, and pointed that the resonance may occur when the natural frequency of the feed system was near the torque ripple harmonic frequency of the motor. Dong and Tang [12] established a hybrid model to study the structural dynamics of the ball screw drive, and analyzed the axial, torsional and flexural dynamic behaviors of ball screw feed system by using the method of continuous beam. Wang et al. [13] developed a time-varying dynamic model to investigate the vibration characteristics of multi-degree-of-freedom feed system, the restoring force of screw nut was considered as a piece-wise function, the coupling effects of screw-nut position and excitation amplitude were investigated by simulation. Xu et al. [14] presented an analytical restoring force function to study the effect of sliding platform position, external excitation, and nominal contact angles. Moreover, the dynamic response of the system is discussed. As high acceleration could cause the large inertial force of the moving components, which lead the change of the contact stiffness, and bring influence to the dynamic behaviors of the ball screw feed system. Considering the effect of acceleration, Zhang et al. [15] proposed an equivalent dynamic model using lumped parameter method, and derived the equivalent axial stiffness of screw nut based on the contact state which is influenced by the variation of inertial force. Moreover, it pointed that the return tube of a ball screw driven mechanism has been designed to provide the path for balls rolling in screw nut grooves. Due to the complexity of working environment, the return tube might damage which is caused by high excitation frequency and excitation amplitude. Huang et al. [16] developed an impact-contact formula for the feed system, and investigated the transient behavior using finite element method. Due to the complexity of working condition, the contact angle of angular contact ball bearing was not a constant value. Liao and Lin [17] established a three-dimensional expression of balls based on the geometry parameters at different position angles, and discussed the effect of radial and axial deformation on the contact deformation, normal force and contact angle. Due to the machining error, assembly error and complicated working condition, the distribution of the load distribution of ball screws was unbalanced. Liu et al. [18] established an appropriate transformation coordinate system to analyze the static load distribution of ball nut. Moreover, the effect of initial contact angle and axial load were discussed in details. In order to analyze the contact load of balls in screw nut when there exist a turning torque caused by assembly error, Zhao et al. [19] establish a load distribution model. Furthermore, the effect of the elastic creep on transmission accuracy was studied. As the sliding wear of ball nut was the main reason which caused the decrease of preload. Wei et al. [20] proposed a two body abrasion model to study the variations of wear depths in the axial direction which directly related to the preload of screw nut. Okwudire and Altintas [21] analytically and experimentally proved that the lateral dynamics of the ball screw feed system could affect the positioning accuracy under the condition of high acceleration. Considering the influence of the deformation of the screw nut and geometry error, Zhou et al. [22] proposed a modified load distribution model. In addition, an experiment was conducted to validate the proposed model. As the increasing of velocity and acceleration of the ball screw drive system, the resonant of the system may degrade the positioning accuracy. Diego et al. [23] developed a high-frequency dynamic model of a ball screw drive feed system. Simultaneously, the frequency variation with different worktable

position was studied. Kolar et al. [24] investigated the effect of frame properties on the dynamic characteristics, and conducted a cutting experiment to validate the simulation result. To reduce the impact of thermal error and control vibration, lots of theoretical analysis and experimental research had been done by many researchers[25-28]. Wang et al. [29] established a quasi-static model with consideration of thermal effect, and proposed a calculation strategy considering the variation of boundary conditions. Using the proposed method, the temperature field and thermal error were predicted accurately. To obtain the key point transient temperatures, Li et al. [30] presented a random thermal network model. Based on the proposed model, the positioning accuracy could be accurately predict. Beyond that, using a finite element model, Mi et al. [31] study the dynamic characteristics of the spindle nose of a horizontal machining center during the cutting process.

From literature reviews, the physical model based methods could obtain the vibration response of the system with consideration of different working parameters. The effect of deflect angle has been rarely studied. Due to the incipient fault and periodic maintenance, the working parameters of screw nut such as geometric error and preload can be changed over time. Furthermore, due to the complexity of working condition, the system can be sensitive to the changed parameters. This can lead the occurrence of unstable vibration, and decrease the stability of the system. However, few papers pay attention to the influence of the variation of working parameters on dynamic characteristics of single nut ball screw with 5 degrees of freedom. The outline of this paper is as follows: Section 2 proposed the 5 degrees-of-freedom nonlinear dynamic model and introduced the relationship between deflection angle and contact deformation. Section 3 revealed the effect of excitation amplitude, deflection angle, the number of balls, and preload by the bifurcation diagram and 3-D frequency spectrum. Section 4 measured the dynamic parameters of system, and validated the proposed model by vibration response and amplitude frequency curve. Section 5 lists the conclusions.

## **2. Dynamic model and governing equations of motion**

As shown in Fig. 1, the single nut consists of ball nut, balls and screw shaft. The single nut is preloaded by variable lead, the value of preload force can be shown in Table.1. Due to the existence of preload and the difference of contact force, the balls in screw nut can be divided into two parts: left section and right section, namely. Therefore, the axial restoring force of a single nut can be deduced into a piecewise function which can be seen in Eqs. 15. In this paper, the feed rate is set to 0. Due to the existence of the initial deflection angle, the deformation of each ball in the screw nut can be different, it depends on the deflection angle  $\varphi$ , the distance  $L_i$  and  $L_j$  which can be shown in Fig.1. Therefore, some balls in screw nut can be deformed due to the unique stress condition, and some balls are not. Because of the nonlinear Hertzian contact force and piecewise restoring function, the vibration system can exhibit wide variety of dynamics.

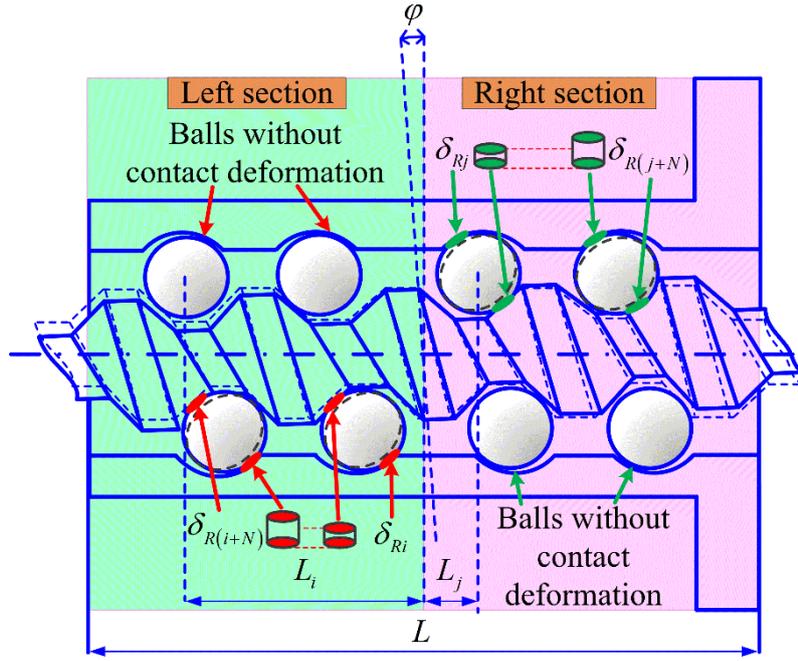


Fig.1 schematic of screw nut

**Table. 1**

Parameters for ball screw

Parameters	Value
Pitch circle diameter $d_m$	42mm
Lead of ball screw $P$	16mm
Ball diameter $D$	7.144mm
Initial contact angle $\alpha_0$	45°
Overall length of screw nut $L$	144mm
Number of loaded circle $N_c$	2.5
Number of loaded balls $N$	23
Mass of screw shaft $m$	5.84kg
Preload $F_p$ [13]	1690N

According to Ref. [32], the expression of the position angle of  $i^{\text{th}}$  ball in the left section of ball nut, and the position angle of the  $j^{\text{th}}$  ball in the right section of ball nut can be given by

$$\theta_i = \frac{2\pi}{N}(i-1) \quad 1$$

$$\theta_j = \frac{2\pi}{N}(j-1) \quad 2$$

where  $N$  is the number of balls per circle in ball nut. As shown in Fig. 1,  $L_i$  is the axial distance between  $i^{\text{th}}$  ball and the center of screw nut in the left section, and  $L_j$  is the axial distance between  $j^{\text{th}}$  ball and the center of screw nut in the right section. The expression of  $L_i$  and  $L_j$  can be given by

$$L_i = \left( N_c - \frac{\theta_i}{2\pi} \right) P \quad 3$$

$$L_j = \frac{\theta_j}{2\pi} P \quad 4$$

$P$  is the lead of screw shaft, and  $N_c$  represents the number of loaded circle. According to Ref. [33-35], the elastic deformation along the radial direction of  $i^{\text{th}}$  ball for the left section of ball nut can be formulated as:

$$\delta_{rLi} = (x + u_i L_i \varphi_y) \cos \theta_i + (y - u_i L_i \varphi_x) \sin \theta_i \quad 5$$

and for the right section, the elastic deformation along the radial direction of  $j^{\text{th}}$  ball of ball nut can be formulated as:

$$\delta_{rRj} = (x + u_i L_j \varphi_y) \cos \theta_j + (y - u_i L_j \varphi_x) \sin \theta_j \quad 6$$

where  $x$  and  $y$  represent the displacement of screw shaft, and  $\varphi_x$  and  $\varphi_y$  are the rocking motions about  $x$  and  $y$  axes Coefficient  $u_i$  is a dimensionless constant which can be determined by [33]:

$$u_i = \begin{cases} -1 & \text{left section} \\ 1 & \text{right section} \end{cases} \quad 7$$

the elastic deformation along the axial direction of  $i^{\text{th}}$  ball for the left section of ball nut can be formulated as [33]:

$$\delta_{zLi} = z + \frac{d_m}{2} (\varphi_x \sin \theta_i - \varphi_y \cos \theta_i) + v_i (A \sin \alpha_0 + z_p) \quad 8$$

Similarly, the elastic deformation along the axial direction of  $j^{\text{th}}$  ball for the right section of ball nut can be formulated as [33]

$$\delta_{zRj} = z + \frac{d_m}{2} (\varphi_x \sin \theta_j - \varphi_y \cos \theta_j) + v_i (A \sin \alpha_0 + z_p) \quad 9$$

where  $d_m$  represents the pitch circle diameter. The dimensionless constant  $v_i$  is dominated on the configuration of rolling elements which can be expressed as follows [33]:

$$v_i = \begin{cases} 1 & \text{left section} \\ -1 & \text{right section} \end{cases} \quad 10$$

Therefore, the total contact deformation along the contact angle direction of the ball in the left section of screw nut can be formulated as:

$$\delta_{Li} = \sqrt{\delta_{rLi}^2 + \delta_{zLi}^2} - A \quad 11$$

$A$  is the distance between screw shaft raceway curvature center and screw nut raceway curvature center. The total contact deformation along the contact angle direction of the ball in the right section of screw nut can be formulated as:

$$\delta_{Rj} = \sqrt{\delta_{rRj}^2 + \delta_{zRj}^2} - A \quad 12$$

In addition, for the left section of screw nut, the contact angle of the  $i^{\text{th}}$  ball can be written as:

$$\alpha_{Li} = \arctan \left( \frac{\delta_{zLi}}{\delta_{rLi}} \right) \quad 13$$

for the right section of screw nut, the contact angle of the  $j^{\text{th}}$  ball in the right section of screw nut can be written as:

$$\alpha_{Rj} = \arctan \left( \frac{\delta_{zRj}}{\delta_{rRj}} \right) \quad 14$$

According to Ref. [13, 36, 37], the axial piecewise restoring force function can be formulated as:

$$F_{NZ}(x, y, z, \varphi_x, \varphi_y, t) = \begin{cases} K \sum_{i=1}^N \delta_{iL}^{3/2} \sin \alpha_{Li} - K \sum_{j=1}^N \delta_{jR}^{3/2} \sin \alpha_{Rj} & -z_p \leq z \leq z_p \\ K \sum_{i=1}^N \delta_{Li}^{3/2} \sin \alpha_{Li} & z > z_p \\ -K \sum_{j=1}^N \delta_{Rj}^{3/2} \sin \alpha_{Rj} & z < -z_p \end{cases} \quad 15$$

where  $K$  is the Hertzian contact stiffness between balls and races, which can be calculated by the method according to Ref. [38]. Furthermore, the restoring force along  $x$  and  $y$  axis can be formulated as:

$$F_{Nx}(x, y, z, \varphi_x, \varphi_y, t) = K \sum_{i=1}^N \delta_{Li}^{3/2} \cos \alpha_{Li} \cos \theta_i + K \sum_{j=1}^N \delta_{Rj}^{3/2} \cos \alpha_{Rj} \cos \theta_j \quad 16$$

$$F_{Ny}(x, y, z, \varphi_x, \varphi_y, t) = K \sum_{i=1}^N \delta_{Li}^{3/2} \cos \alpha_{Li} \sin \theta_i + K \sum_{j=1}^N \delta_{Rj}^{3/2} \cos \alpha_{Rj} \sin \theta_j \quad 17$$

The restoring moments about  $x$  and  $y$  axis caused by the deformation of balls in screw nut can be expressed as follows:

$$M_{Nx}(x, y, z, \varphi_x, \varphi_y, t) = K \sum_{i=1}^N \delta_{Li}^{3/2} \left( \frac{d_m}{2} \sin \alpha_{Li} - L_i \cos \alpha_{Li} \right) \sin \theta_i \quad 18$$

$$+ K \sum_{j=1}^N \delta_{Rj}^{3/2} \left( \frac{d_m}{2} \sin \alpha_{Lj} + L_j \cos \alpha_{Lj} \right) \sin \theta_j$$

$$M_{Ny}(x, y, z, \varphi_x, \varphi_y, t) = K \sum_{i=1}^N \delta_{Li}^{3/2} \left( -\frac{d_m}{2} \sin \alpha_{Li} - L_i \cos \alpha_{Li} \right) \cos \theta_i \quad 19$$

$$+ K \sum_{j=1}^N \delta_{Rj}^{3/2} \left( -\frac{d_m}{2} \sin \alpha_{Lj} + L_j \cos \alpha_{Lj} \right) \cos \theta_j$$

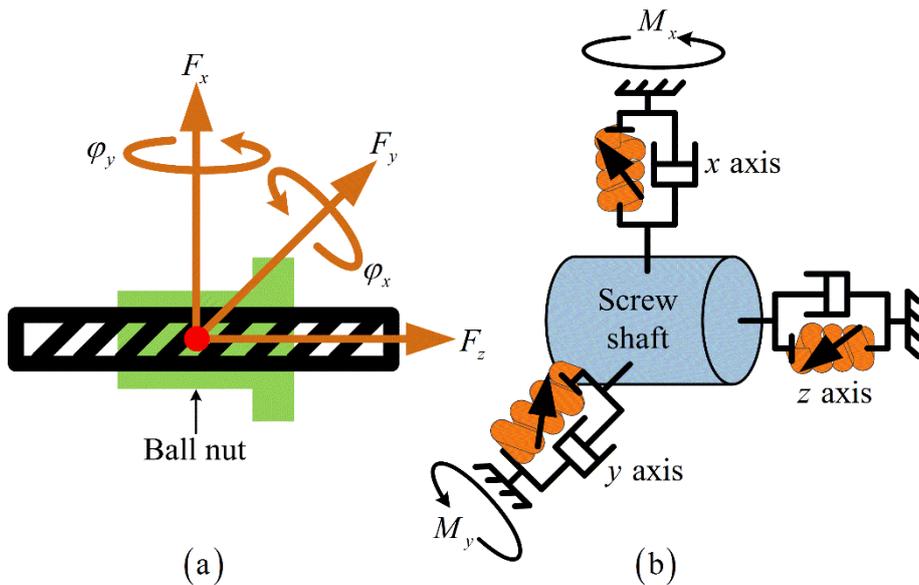


Fig.2 A lumped spring mass model schematic of ball screw. Considering the nonlinear contact restoring force and piecewise restoring force,

the equivalent dynamic model of screw nut can be simplified as a 5- degrees-of-freedom mass-spring-damping model which can be shown in Fig.2. According to the analysis of restoring force and restoring moment along  $x$ ,  $y$ ,  $z$ ,  $\varphi_x$ , and  $\varphi_y$  directions, the governing equations of motion for screw nut with 5 degrees-of-freedom can be written as follows:

$$m\ddot{x} + c_y\dot{x} + F_{NLx}(x, y, z, \varphi_x, \varphi_y, t) = mg + F_x(t) \quad 20$$

$$m\ddot{y} + c_y\dot{y} + F_{NLy}(x, y, z, \varphi_x, \varphi_y, t) = F_y(t) \quad 21$$

$$m\ddot{z} + c_z\dot{z} + F_{NZ}(x, y, z, \varphi_x, \varphi_y, t) = F_z(t) \quad 22$$

$$I_x\ddot{\varphi}_x + c_y\dot{\varphi}_x + M_{Nx}(x, y, z, \varphi_x, \varphi_y, t) = M_x \quad 23$$

$$I_y\ddot{\varphi}_y + c_y\dot{\varphi}_y + M_{Ny}(x, y, z, \varphi_x, \varphi_y, t) = M_y \quad 24$$

where an overdot describes time differentiation;  $m$  represents the mass of screw shaft;  $c_y$  and  $c_z$  represent the viscous damping coefficient along radial and axial direction. The influence of damping originate from the viscous damping is suitable for the structural characteristics of the proposed model [39], the value of  $c_y$  and  $c_z$  can be identified by the test of vibration response, which can be shown in Section.4.  $F_x$ ,  $F_y$ , and  $F_z$  represent the external load. In this study, the expression of external force satisfy  $F_x=F_y=F_z=F_0\sin\omega t$ , where  $F_0$  is the excitation amplitude, and  $\omega$  represents the excitation frequency.  $M_x$  and  $M_y$  are the external moment, and satisfy  $M_x=M_y=lF_0\sin\omega t$ ,  $l$  is the arm of force and satisfy  $l=0.5L$ .

### 3. Results and discussion

In this section, the numerical solution to the governing equations of the proposed dynamic model is archived by fourth order Runge Kutta method. The parameters of studied ball screw is listed in Table.1. In order to accelerate the convergence of calculation, the initial condition for the differential equations is set to  $[\dot{x}, x, \dot{y}, y, \dot{z}, z] = [-3.5741 \times 10^{-4}, -3.4863 \times 10^{-4}, -9.6292 \times 10^{-4}, -3.5542 \times 10^{-4}, -2.0984 \times 10^{-4}, -1.0713 \times 10^{-4}, -3.8483 \times 10^{-4}, -8.5445 \times 10^{-4}, 4.5331 \times 10^{-4}, 1.9655 \times 10^{-4}]$ , the values is determined by the mean value of steady state solution. In this study, a steady state for a differential equation is a solution where the nature of the motion does not change over time. Considering the time consumption and the required computational accuracy, the time step is set to  $5 \times 10^{-6}$ , and the relative tolerance are set to  $10^{-5}$  and  $10^{-5}$ . In order to reflects the energy of vibration signal, the root-mean-square value (RMS) is employed to express the amplitude instead of conventional amplitude in amplitude frequency curve. The definition of RMS value can be given by [40, 41]

$$\text{RMS} = \sqrt{\frac{\sum_{i=1}^N (x_i - \bar{x})^2}{N}} \quad 25$$

where  $x_i$  represents the  $i^{\text{th}}$  value of the signal,  $N$  is the dimension of the signal, and  $\bar{x}$  denotes the mean value of the signal.

#### 3.1 Effect of excitation amplitude

Excitation amplitude is one of the most significant parameters that influence the dynamic behavior of vibration system. In this subsection, the dynamic behaviors are investigated by bifurcation diagram and 3D-spectrum with excitation amplitude and excitation frequency as control parameters. The waveform, spectrum, phase diagram and Poincaré section are used to analyze the complex dynamic behavior at some specific points.

The corresponding bifurcation diagrams of the vibration system with excitation frequency as control parameter under different excitation amplitude ( $F_0=3000\text{N}$ ,

$F_0=6000\text{N}$ ,  $F_0=9000\text{N}$ ,  $F_0=12000\text{N}$ ,  $F_0=15000\text{N}$ ) can be shown in Fig.3. With the increase of excitation amplitude, the vibration system exhibits wide variety of dynamic behaviors including chaotic motion, quasi-periodic motion and periodic- $n$  motion. In the range of  $\omega \in [100, 610]\text{rad/s}$ , with the increase of excitation amplitude, more jump discontinuous phenomena appear, and the system exhibits chaotic motion in this interval. In the range of  $[610, 760]\text{rad/s}$  which is near resonance frequency, the extent of chaotic motion increases markedly, the chaotic characteristics strengthen. This indicates that the excitation amplitude have impacts on the dynamic behavior of vibration system. Increasing the excitation amplitude, in the range of  $[960, 1210]\text{rad/s}$ , the influence range of chaotic motion relative to the whole excitation frequency broadens. Further increase the excitation frequency, the motion of the system changed from simple periodic-1 motion to quasi-periodic motion and chaotic motion in the range of  $[1350, 1730]\text{rad/s}$ .

The bifurcation diagram of the system at  $F_0=15000\text{N}$  is exhibited in Fig. 4. The system jump from one to the other at relatively lower excitation frequency of interval  $T_1$ . In the range of  $T_1$ , the system exhibits a narrow quasi-periodic motion and chaotic motion. As excitation frequency is increased, the system shows chaotic motion in the range of  $T_2$ . Increasing the excitation frequency, the system exhibits periodic- $n$  and chaotic motion in the interval  $T_3$ . By increasing excitation frequency further, after experiencing chaotic motion in the range of  $T_4$ , the motion of the system from simple periodic-1 motion in the interval  $T_5$ , enters to periodic- $n$  and quasi-periodic motion in the range of  $T_6$ , although at some region the system exhibits chaotic motion, finally to periodic-1 motion at  $T_7$ .

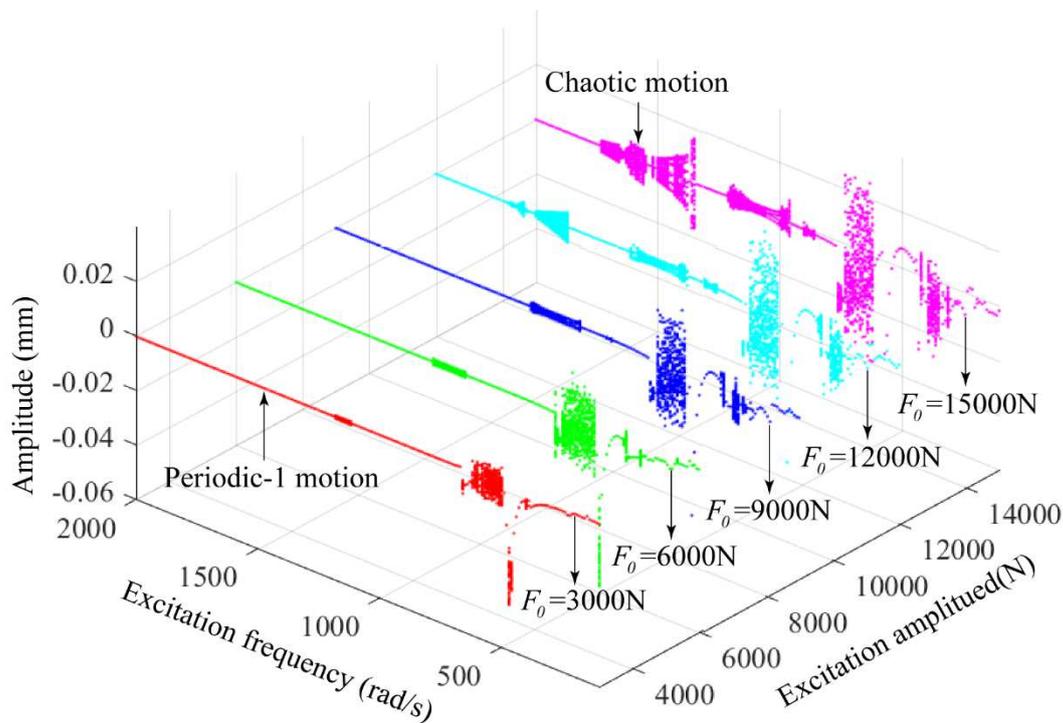


Fig.3 Bifurcation diagrams with excitation frequency  $\omega$  as the control parameter at  $\varphi=0^\circ$ ,  $N=23$ , and Preload= $F_p$  for different excitation amplitudes  $F_0$ .

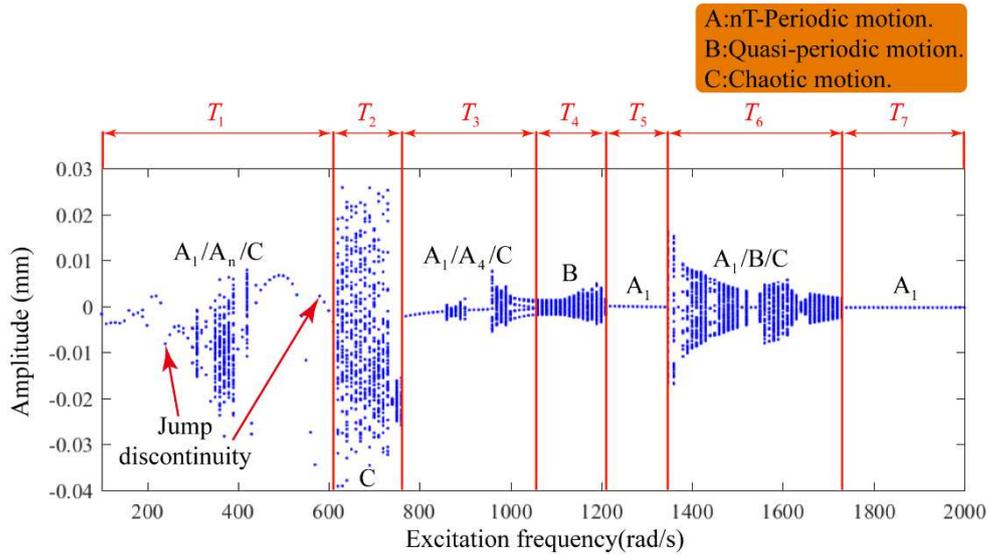


Fig.4 Bifurcation diagrams with excitation frequency  $\omega$  as the control parameter at  $\varphi=0^\circ$ ,  $N=23$ , and Preload= $F_p$  for  $F_0=15000N$ .

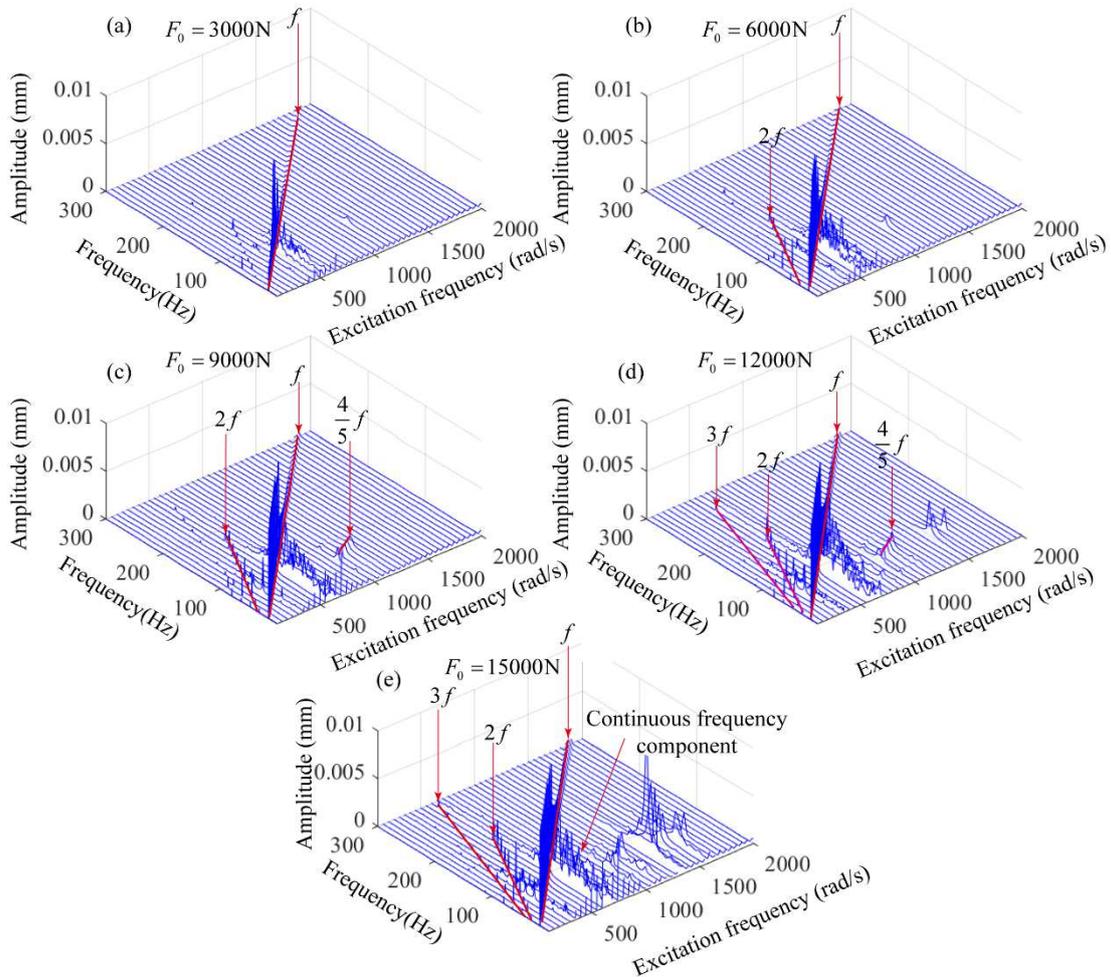


Fig.5 3-D frequency spectrums for different excitation amplitude with excitation frequency  $\omega$  as the control parameter at  $\varphi=0^\circ$ ,  $N=23$  and Preload= $F_p$  for (a)  $F_0=3000N$ , (b)  $F_0=6000N$ , (c)  $F_0=9000N$ , (d)  $F_0=12000N$ , (e)  $F_0=15000N$ .

The 3-D frequency spectrum with respect to excitation frequency for different excitation amplitude is illustrated in Fig.5. At low excitation amplitude, the

fundamental frequency  $f$  is the dominant frequency component. As the excitation amplitude is increased, the continuous frequency component appears in the interval  $[610, 760]$ rad/s and frequency multiplication ( $2f, 3f$ ) become more obvious. Furthermore, frequency demultiplication ( $4f/5$ ) appears when  $F_0=9000$ N, and the main frequency changes from the fundamental frequency  $f$  to frequency demultiplication ( $4f/5$ ) in the range of  $[1350, 1700]$ rad/s when  $F_0>9000$ N.

To further investigate the effect of excitation amplitude on dynamic behavior of the system, the dynamic response are illustrated with excitation amplitude  $F_0$  as control parameter for different excitation amplitude in Fig. 6. In Fig. 6(a1), the motion of the system exhibits simple periodic-1 and periodic-3 motion in interval  $T_1$ , and the main frequency is the fundamental frequency  $f$  in Fig. 6(b1). With the increase of excitation amplitude, the system shows quasi-periodic and chaotic motion in the range of  $T_2$ , and the vibration system appears continuous frequency component. As shown in Fig. 6(a2) and Fig. 6(a3), the system exhibits simple periodic-1 motion at  $T_1$ , and quasi-periodic motion in interval  $T_2$ . Beyond that, as shown in Fig. 6(b3),  $3f/10$  becomes the main frequency as excitation amplitude is  $F_0>8600$ N.

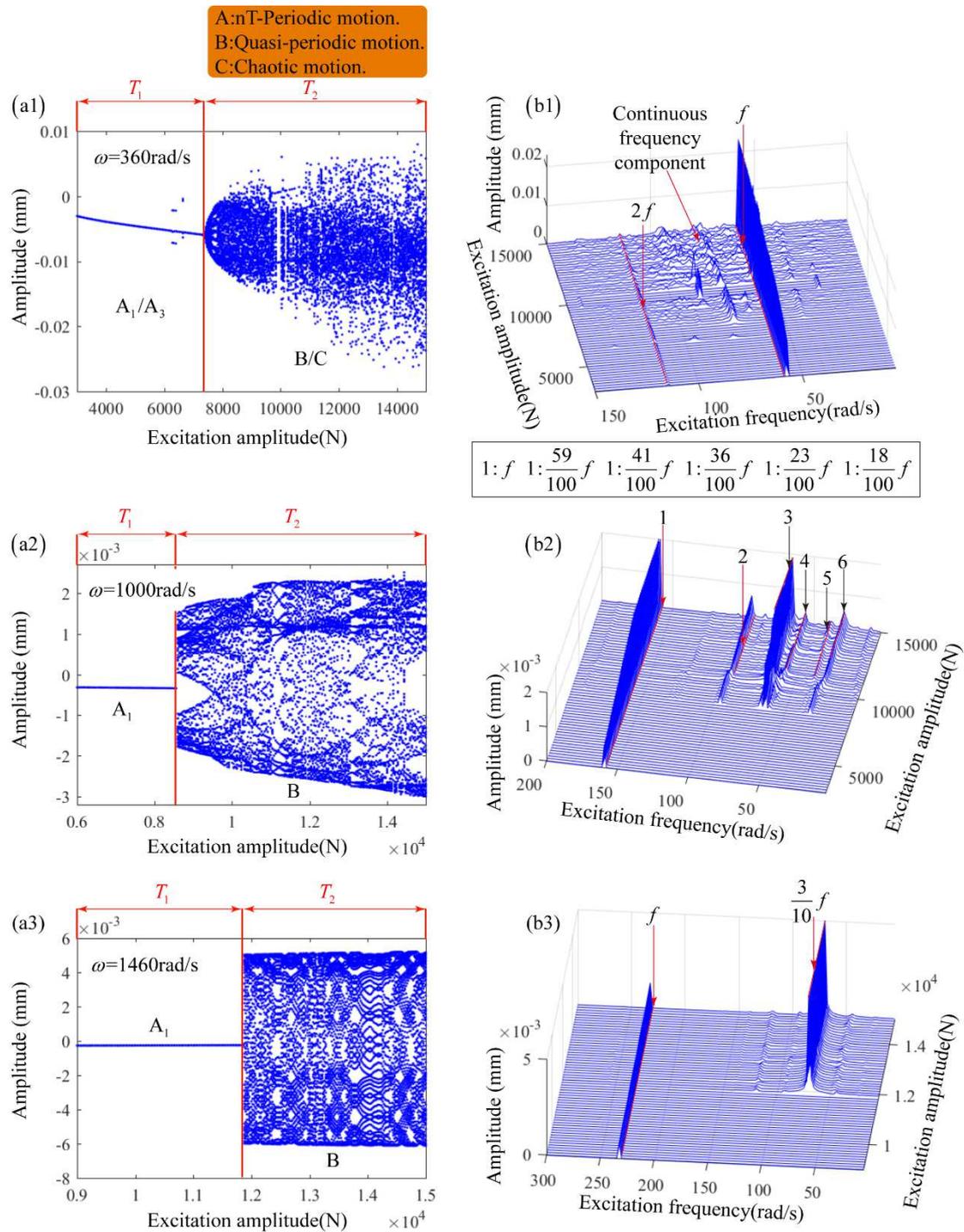


Fig.6 Dynamic responses of the system with excitation amplitude  $F_0$  as control parameter at  $\varphi=0^\circ$ ,  $N=23$  and Preload= $F_p$  for  $\omega=360\text{rad/s}$ ,  $\omega=1000\text{rad/s}$  and  $\omega=1460\text{rad/s}$  (a) Bifurcation diagram; (b) 3-D frequency spectrum.

Fig. 7 shows the corresponding vibration response at  $F_0=3000\text{N}$  and  $F_0=15000\text{N}$  when  $\omega=1000\text{rad/s}$ . In Fig. 7(c1), the Poincaré section appears only one point and the phase diagram shows a closed circle, which further proves the system exhibits a simple periodic-1 motion. In Fig. 7(c2), there exist finite points in the Poincaré section and the phase diagram is regular, where the system exhibits quasi-periodic motion. The comparison of vibration response between  $F_0=9000\text{N}$  and  $F_0=15000\text{N}$  at  $\omega=1460\text{rad/s}$  is shown in Fig. 8, where the system presents quasi-periodic motion at  $F_0=15000\text{N}$ . Fig. 9 shows the comparison between  $F_0=12000\text{N}$  and  $F_0=15000\text{N}$  at  $\omega=1590\text{rad/s}$ . When

$F_0=15000\text{N}$ , the waveform loses the periodicity in Fig. 9(a2), the continuous spectral appears in Fig. 9(b2). As can be seen in Fig. 9(c2), there exist infinite point in the Poincaré section, the phase diagram is irregular and loses periodicity, which illustrates that the system presents chaotic motion at this point.

Therefore, at relative low excitation amplitude, the vibration system exhibit less nonlinear dynamic properties. In addition, except the region near resonance frequency, the motion of the system exhibit simple periodic motion. When the excitation amplitude become higher, the system presents wide variety of dynamic behaviors. These further proved that the excitation amplitude has considerable impact on the dynamic characteristic of the system.

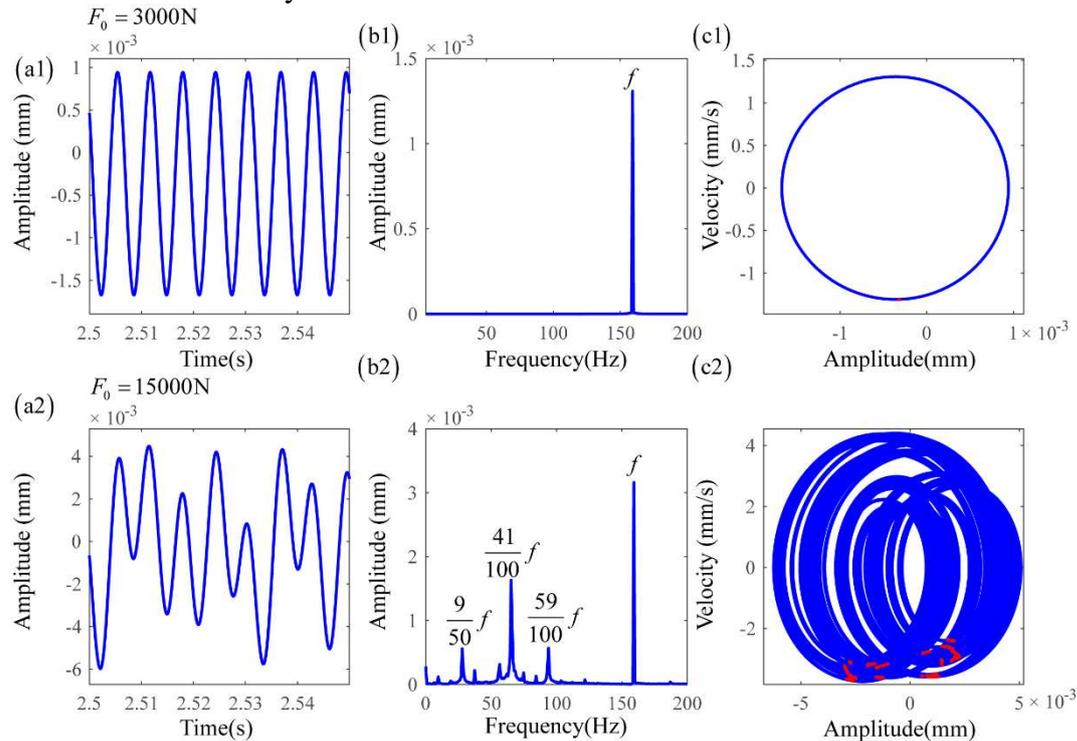


Fig.7 Vibration responses at  $\varphi=0^\circ$ ,  $N=23$ ,  $\omega=1000\text{rad/s}$  and Preload= $F_p$  for  $F_0=3000\text{N}$  and  $F_0=15000\text{N}$ . (a) Waveform, (b) Frequency, (c) Phase diagram and Poincaré section.

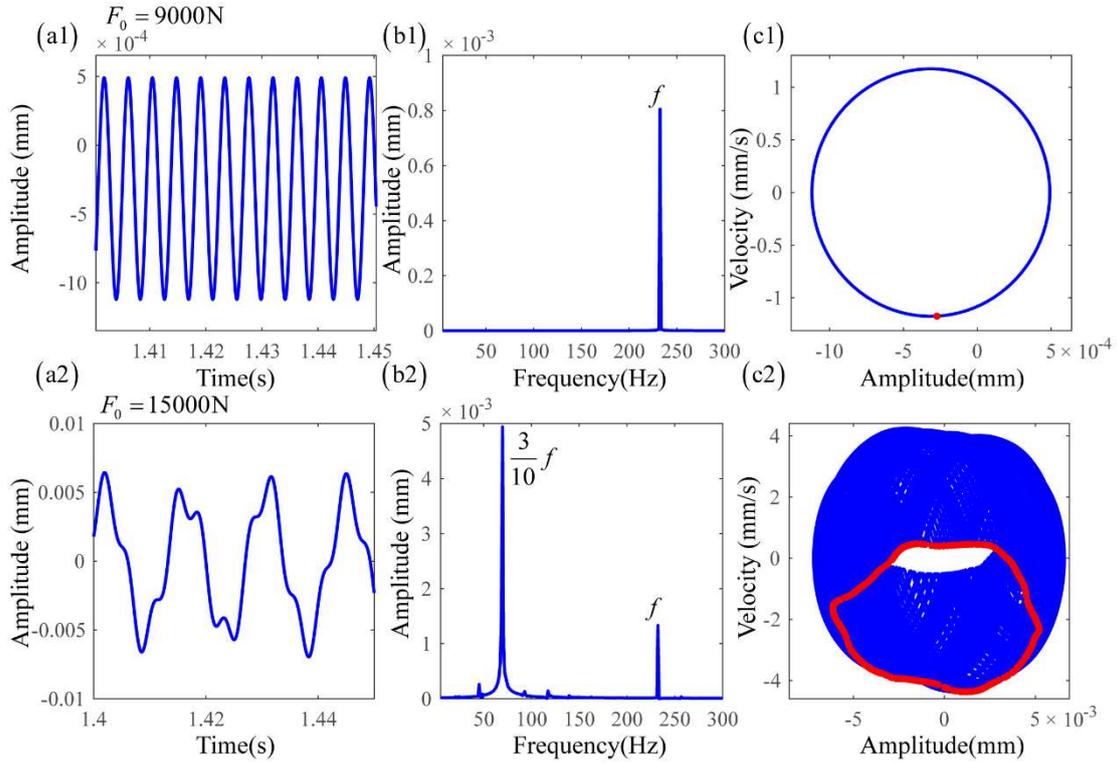


Fig.8 Vibration responses at  $\varphi=0^\circ$ ,  $N=23$ , and  $\omega=1460\text{rad/s}$  and Preload= $F_p$  for  $F_0=9000\text{N}$  and  $F_0=15000\text{N}$ . (a) Waveform, (b) Frequency, (c) Phase diagram and Poincaré section.

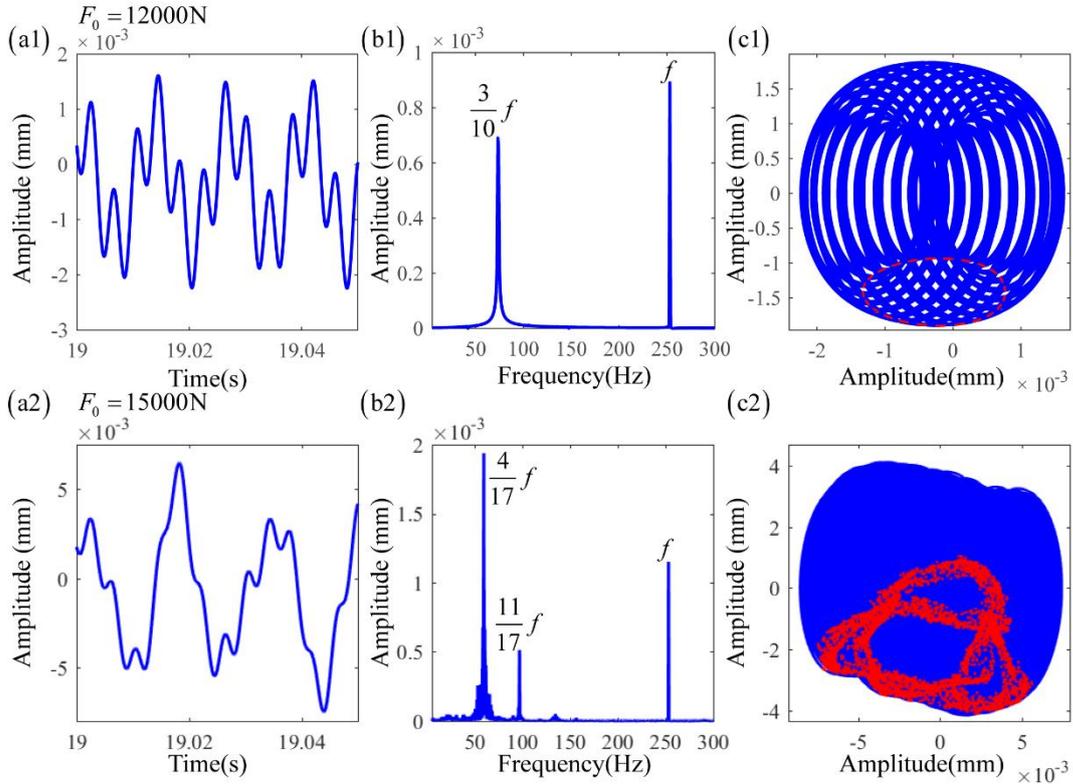


Fig.9 Vibration responses at  $\varphi=0^\circ$ ,  $N=23$ ,  $\omega=1590\text{rad/s}$  and Preload= $F_p$  for  $F_0=12000\text{N}$  and  $F_0=15000\text{N}$ . (a) Waveform, (b) Frequency, (c) Phase diagram and Poincaré section.

### 3.2 Effect of deflection angle

In the previous studies, the deflection angle of ball screw is often ignored.

However, due to human error and assembly error the deflection angle is inevitable to some extent. Considering the complex working conditions, the investigation of deflection angle on nonlinear dynamic characteristics is crucial. In this subsection, the nonlinear dynamics of the system with deflection angle and excitation frequency as control parameter are examined.

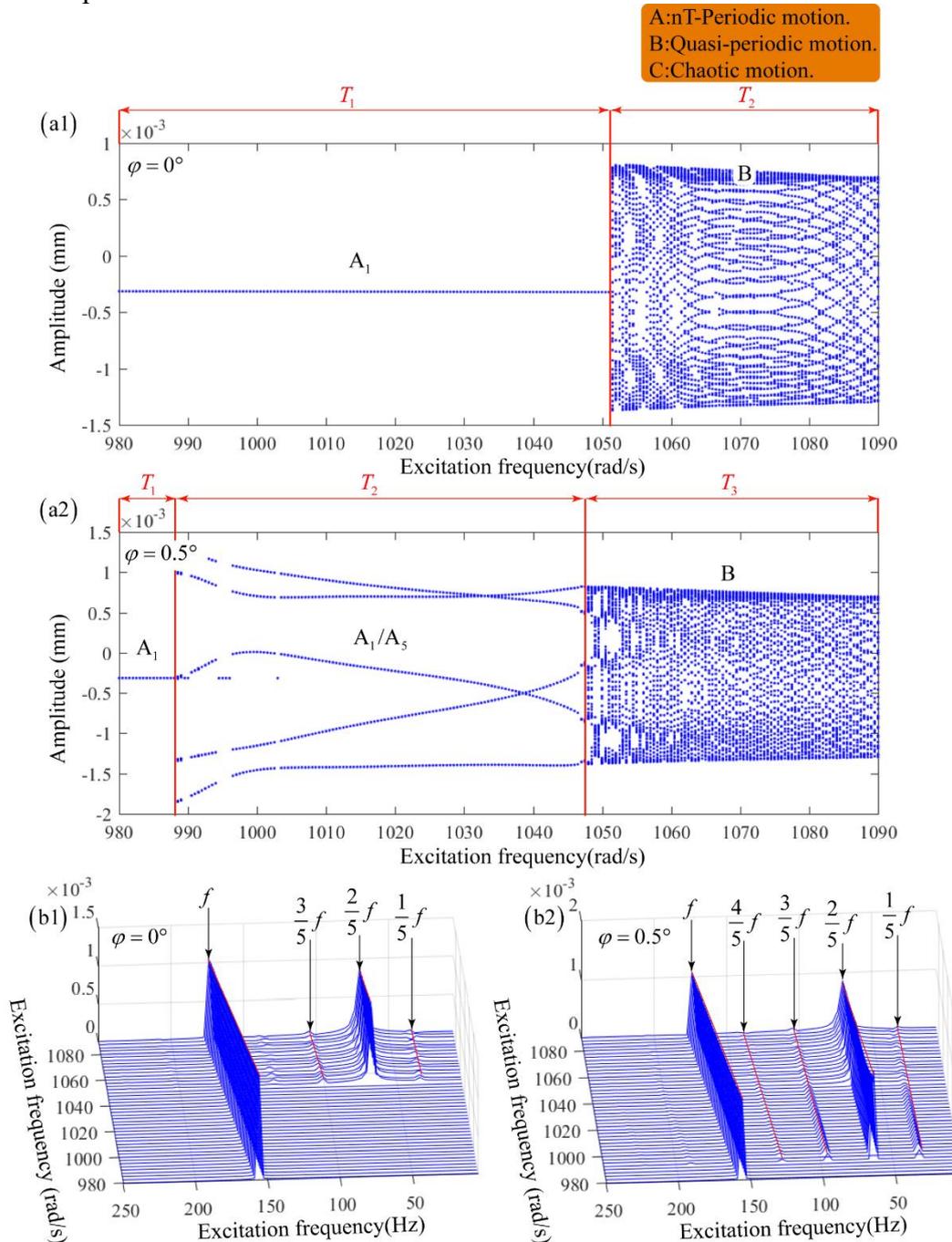


Fig.10 Dynamic responses of the system with excitation frequency  $\omega$  as control parameter at  $F_0=6000\text{N}$ ,  $N=23$  and Preload= $F_p$  for  $\varphi=0^\circ$  and  $\varphi=0.5^\circ$ .(a) Bifurcation diagram; (b) 3-D frequency spectrum.

The comparison of dynamic response between  $\varphi=0^\circ$  and  $\varphi=0.5^\circ$  can be shown in Fig. 10. As shown in Fig. 10(a1), the vibration system exhibits simple periodic-1 motion in the interval  $T_1$ , and in Fig. 10(b1) the main frequency is the fundamental frequency  $f$ . As the excitation frequency is increased, the system exhibits quasi-periodic motion in the range of  $T_2$ , and frequency demultiplication ( $f/5$ ,  $2f/5$ ,  $3f/5$ ) appears in Fig. 10(b1).

Increasing the deflection angle from  $0^\circ$  to  $0.5^\circ$ , the nonlinear characteristics strengthen. As seen in Fig.10(a2), the system shows periodic-1 motion in the range of  $T_1$ , experiencing periodic-1 and periodic-5 motion in the range of  $T_2$ , to quasi-periodic motion in the interval  $T_3$ . It can be seen in Fig. 10(b2), the fundamental frequency is the dominant frequency component, and the frequency demultiplication ( $f/5, 2f/5, 3f/5, 4f/5$ ) is lower than  $f$ . With the increase of excitation amplitude, the motion of system enters to quasi-periodic motion in the interval  $[1047,1090]$ rad/s. Furthermore, no complicated chaotic motion appears.

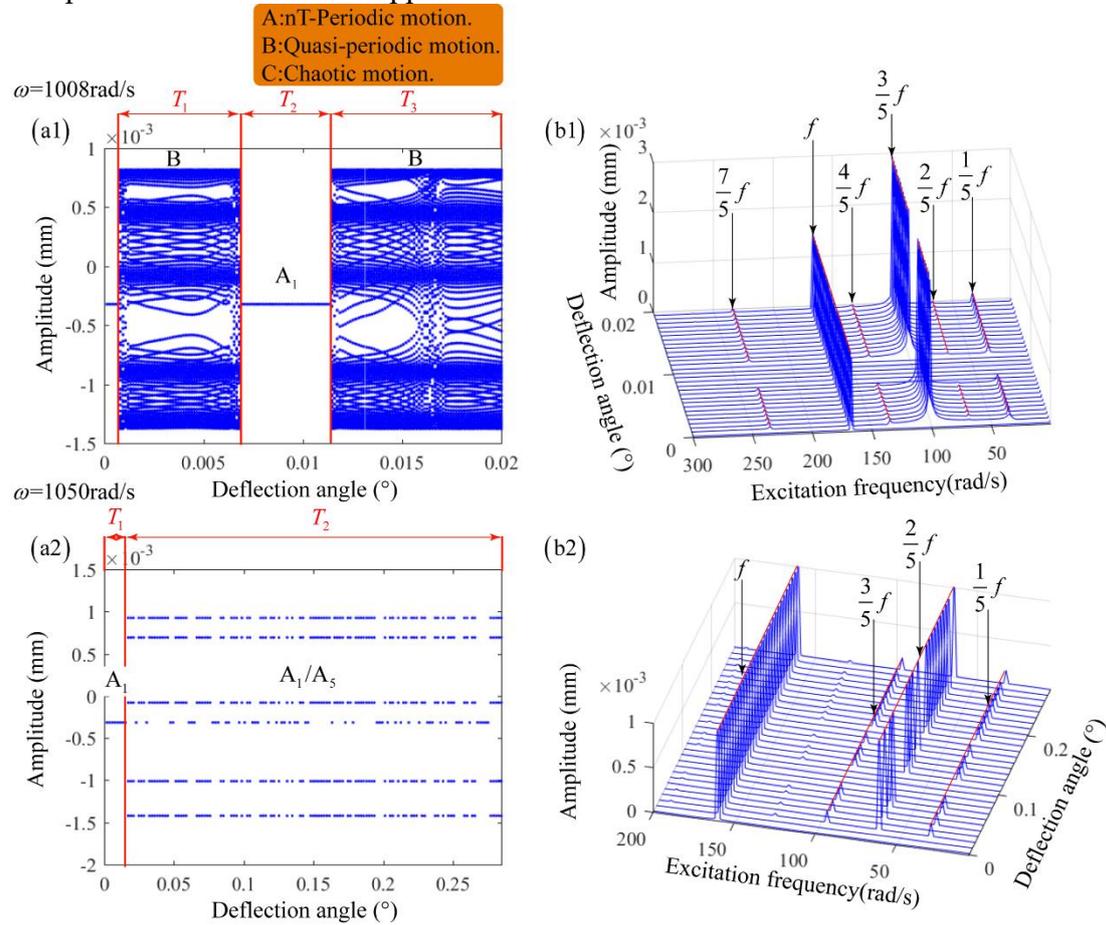


Fig.11 Dynamic responses of the system with deflection angle  $\varphi$  as control parameter at  $F_0=6000$ N,  $N=23$  and Preload= $F_p$  for  $\omega =1008$ rad/s and  $\omega =1050$ rad/s. (a) Bifurcation diagram; (b) 3-D frequency spectrum.

In order to further study the effect of deflection angle on dynamic response. Fig. 11 shows the corresponding bifurcation diagram and 3-D frequency spectrum with deflection angle as control parameter at  $\omega =1008$ rad/s and  $\omega =1050$ rad/s. It can be seen in Fig. 11(a1) and (b1), after experiencing a short periodic-1 motion at relative small deflection, the system enters into quasi-periodic motion in the interval  $T_1$ , the main frequency is demultiplication frequency  $3f/5$ , and the fundamental frequency  $f$  is the second largest. Increasing the deflection angle, the motion of the system presents periodic-1 motion again in the range of  $T_2$ , and the fundamental frequency  $f$  becomes the largest frequency component. By increasing the deflection angle further, the system enters to quasi-periodic motion, and demultiplication frequency  $3f/5$  becomes the dominant frequency in the interval  $T_3$ . Similarly, as shown in Fig. 11(a2), the system exhibits simple periodic-1 motion in the interval  $T_1$ , and the system presents two types of motion state at  $T_2$ , i.e. periodic-1 motion and periodic-5 motion, namely. Moreover, the fundamental frequency  $f$  is the dominant frequency component and the

demultiplication frequency component  $2f/5$  is the second largest which can be shown in Fig. 11(b2). The comparison of vibration response between  $\varphi=0^\circ$  and  $\varphi=0.02^\circ$  at  $\omega=1050\text{rad/s}$  is shown in Fig. 12. As seen in the figure, with the increase of deflection angle, the motion of the system changed from simple periodic-1 motion to quasi-periodic motion. Moreover, the comparison of vibration response between  $\varphi=0^\circ$  and  $\varphi=0.5^\circ$  at  $\omega=1008\text{rad/s}$  is shown in Fig. 13 which indicates that the increase of deflection leads the motion of the system to change from periodic-1 motion to quasi-periodic motion.

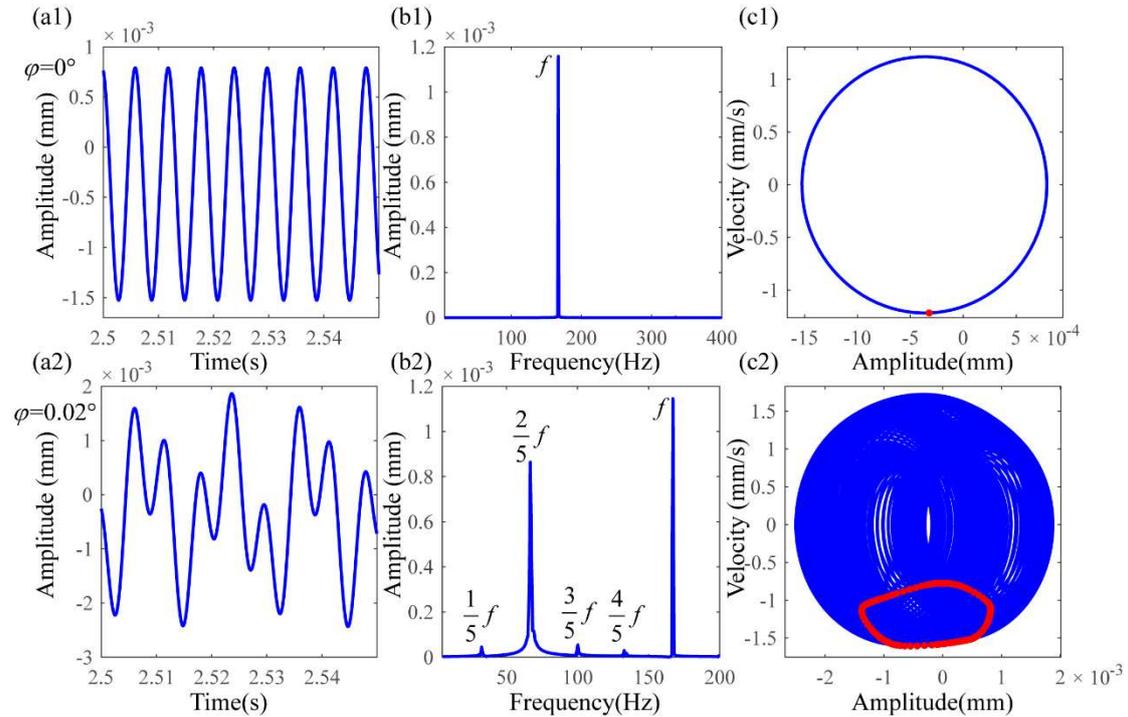


Fig.12 Vibration responses at,  $N=23$ ,  $F_0=6000\text{N}$   $\omega=1050\text{rad/s}$  and Preload= $F_p$  for  $\varphi=0^\circ$  and  $\varphi=0.02^\circ$ . (a) Waveform, (b) Frequency, (c) Phase diagram and Poincaré section.

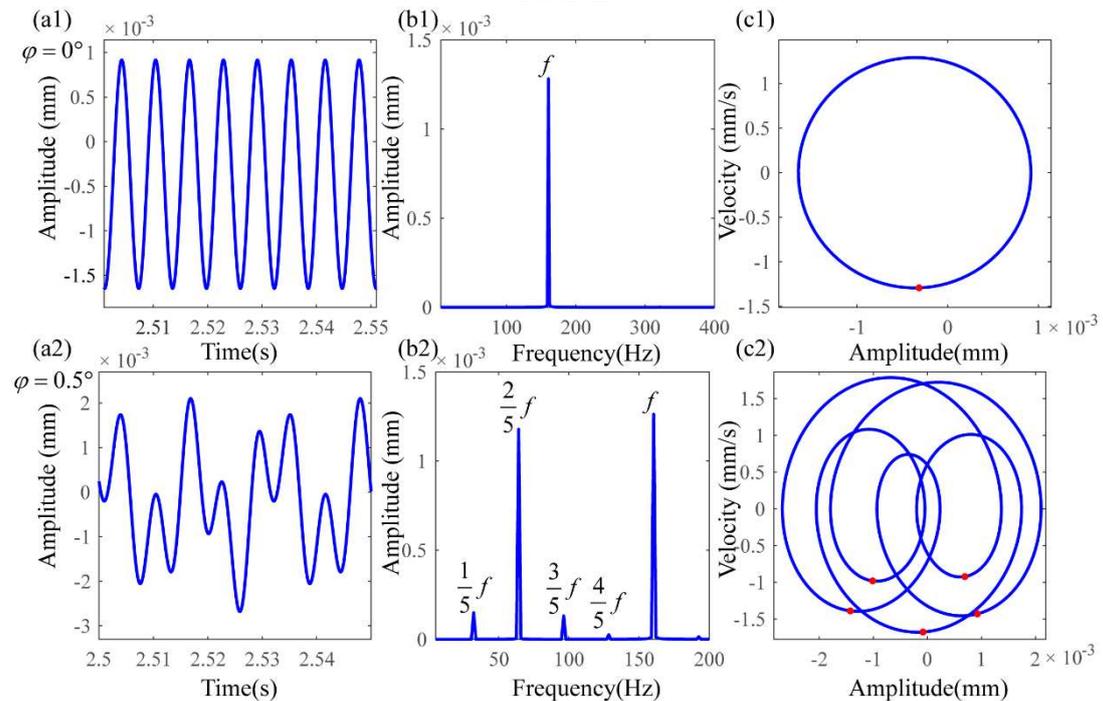


Fig.13 Vibration responses at  $N=23$ ,  $F_0=6000\text{N}$   $\omega=1008\text{rad/s}$  and Preload= $F_p$  for  $\varphi=0^\circ$

and  $\varphi=0.5^\circ$ . (a) Waveform, (b) Frequency, (c) Phase diagram and Poincaré section.

### 3.3 Effect of the number of balls

The number of balls is closely corresponding to the required rigid and the lead of screw shaft, then significantly changes the dynamic characteristics of ball screw. Due to the range of working frequency is fixed, it is important to determine a suitable number to guarantee the dynamic performance of the system. In this subsection, the dynamic responses are exhibited with respect to excitation frequency. To further study the effect of the number of balls on the dynamic characteristic, the bifurcation diagrams with different number of balls ( $N=10$ ,  $N=23$ ,  $N=36$ ) are shown in Fig. 14. It can be seen in the figure, chaotic motion, quasi-periodic motion, and periodic- $n$  motion can make a distinction between each other, and the critical transformation frequency among different kinds of motions can be identified. As the decrease of the number of balls, all three cases present chaotic motion near the resonance frequency, and the corresponding resonance frequency is reduced. Moreover, at the excitation frequency which is lower than resonance frequency, jump discontinuous phenomena can be observed repeatedly.

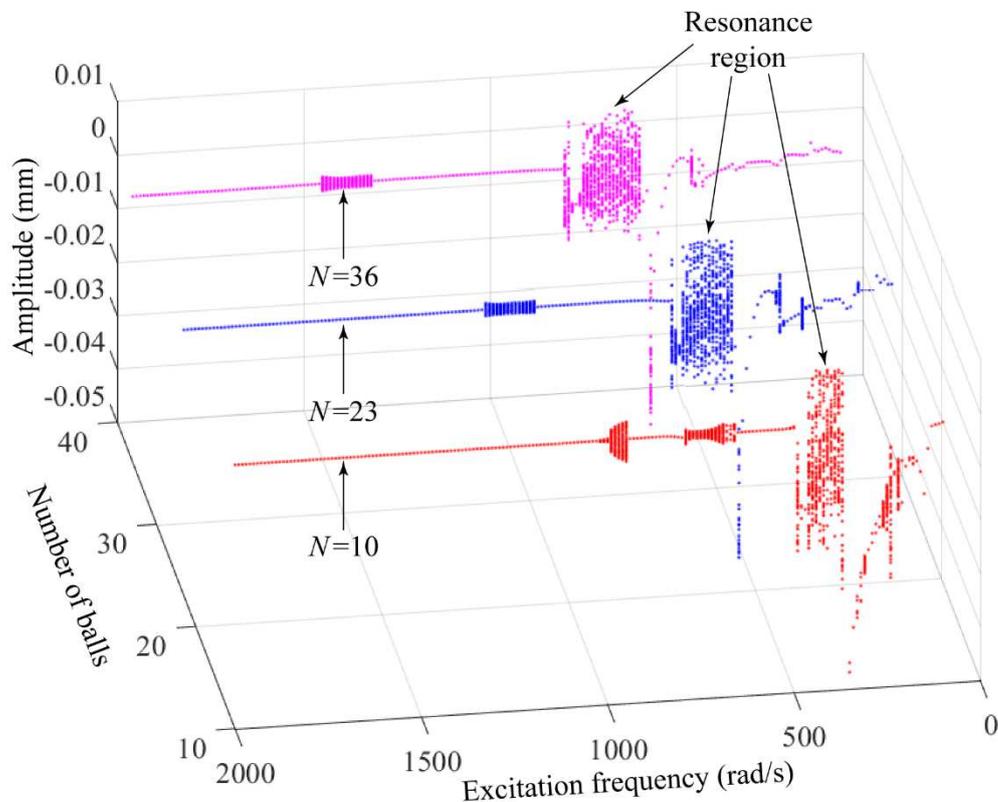


Fig.14 Bifurcation diagrams with excitation frequency  $\omega$  as the control parameter at ,  $F_0=6000\text{N}$ ,  $\varphi=0^\circ$ , and Preload= $F_p$  for  $N=10$ ,  $N=23$  and  $N=36$ .

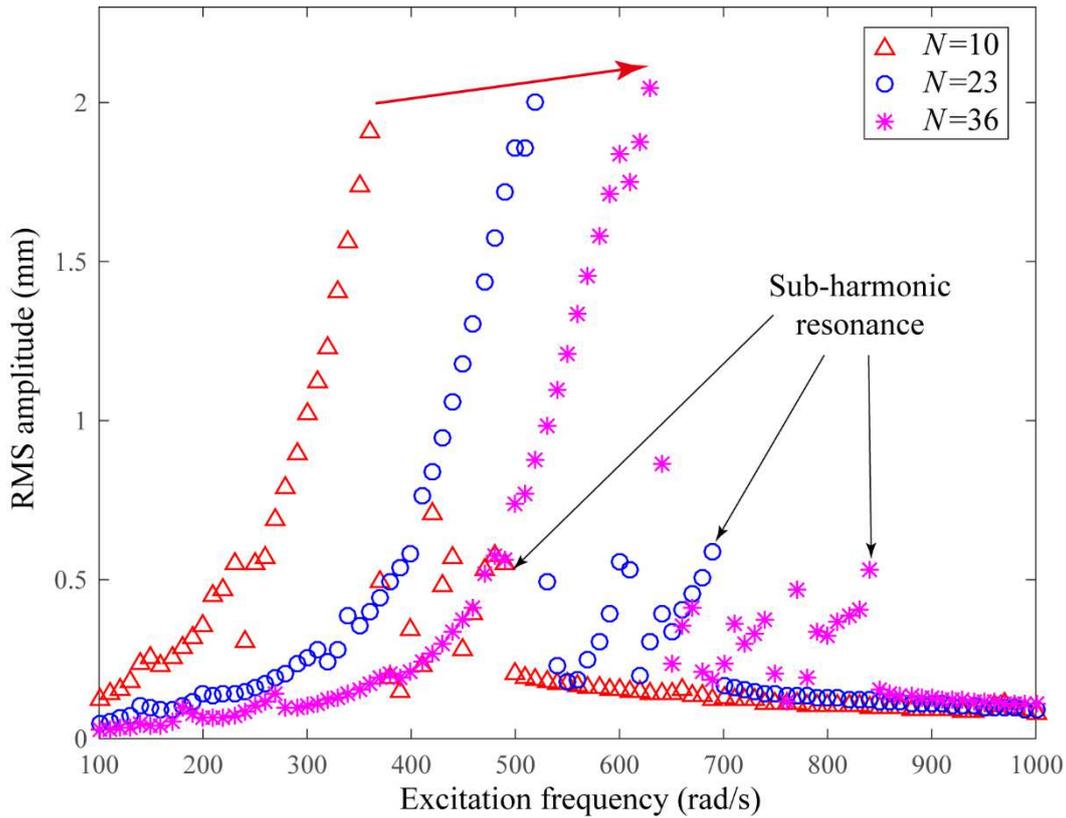


Fig.15 Amplitude frequency curve at  $N=23$ ,  $F_0=6000\text{N}$ ,  $\varphi=0^\circ$ , and Preload= $F_p$  for the number of balls  $N=10$ ,  $N=23$  and  $N=36$ .

Depending on the variation of stiffness caused by different number of balls, the resonance frequency increase with the number of balls. As shown in Fig. 15, the resonance frequency and sub-harmonic resonance frequency increase, the peak value of amplitude frequency curves enhance. In order to further illustrate the influence of the number of balls, Fig. 16 presents the corresponding bifurcation diagrams at relative low excitation frequency  $\omega \in [100, 500]\text{rad/s}$ . In Fig. 16(a), in the interval  $T_1$ , the response repeatedly appears jump discontinuous phenomenon, and the system exhibits periodic-1 and periodic-2 motion. Increasing the excitation frequency further, the system exhibits periodic- $n$  and chaotic motion in the range of  $T_2$ . As seen in Fig. 16(b), when the number of balls decrease from  $N=36$  to  $N=10$ , the system exhibits a wide variety of dynamic behaviors in the range of  $\omega \in [100, 500]\text{rad/s}$ . In the interval  $T_1$ , the motion of the system is periodic-1. As the increase of excitation frequency, different motion state can be observed such as: periodic- $n$ , quasi-periodic, and chaotic motion in the range of  $T_2$ . After experiencing chaotic motion in interval  $T_3$ , the motion of the system enters to periodic-1 and chaotic motion in the range of  $T_4$ . The corresponding 3-D frequency spectrums are shown in Fig.17, and the continuous frequency components can be distinguished from other discrete frequency components in Fig.17(b).

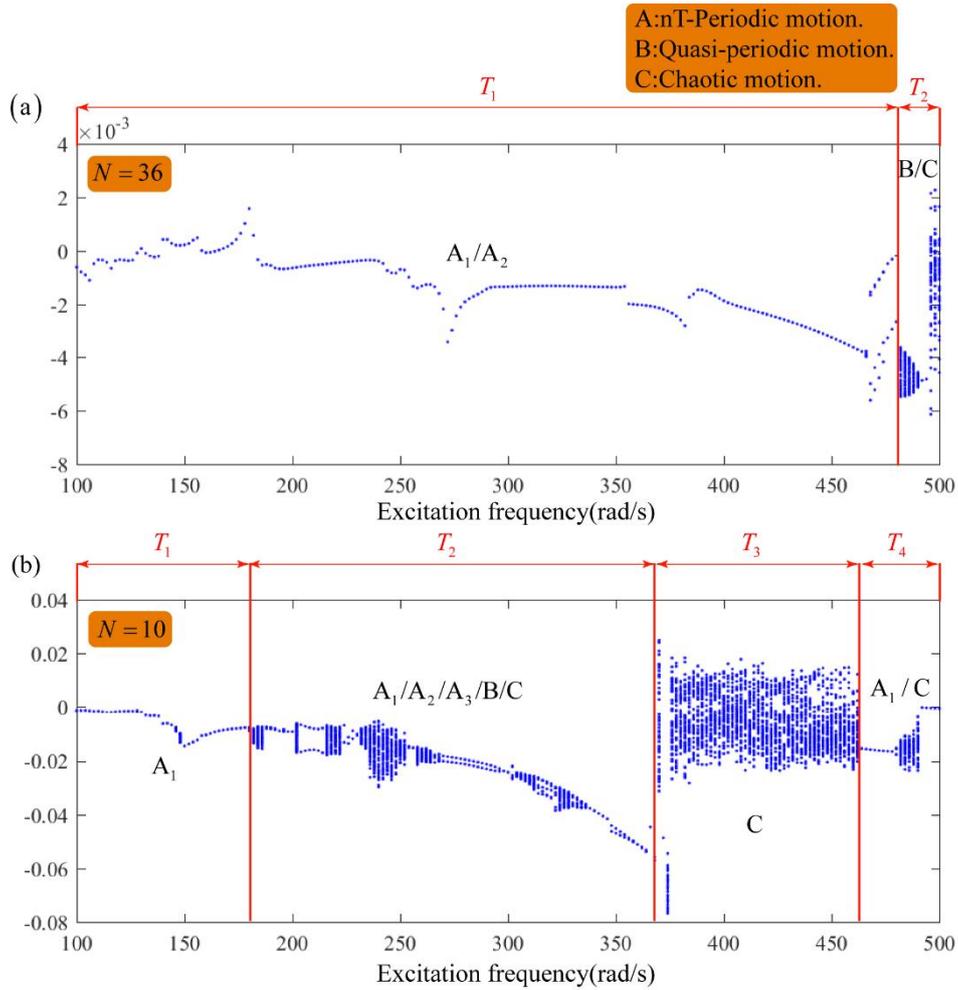


Fig.16 Bifurcation diagrams with excitation frequency  $\omega$  as the control parameter at  $\varphi=0^\circ$ ,  $F_0=6000\text{N}$ , and Preload= $F_p$  for different number of balls. (a)  $N=36$ , (b)  $N=10$ .

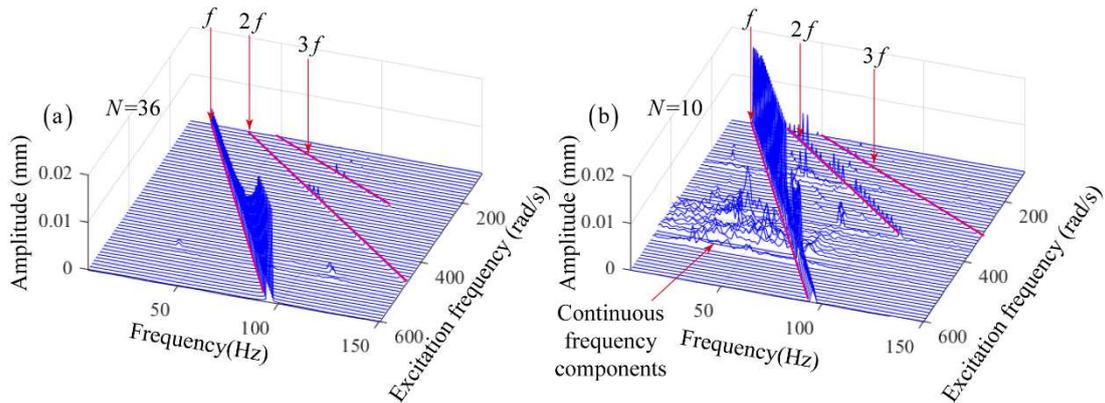


Fig.17 3-D frequency spectrums with excitation frequency  $\omega$  as the control parameter at  $\varphi=0^\circ$ ,  $F_0=6000\text{N}$ , and Preload= $F_p$  for different number of balls. (a)  $N=36$ , (b)  $N=10$ .

### 3.4 Effect of preload

By creating a load between screw shaft and ball nut, a preload can bring greater stiffness to ball screw. Therefore, the displacement of structure can be reduced when the system under external force. In this subsection, to investigate the influence of preload on the dynamic behaviors, the preload  $F_p$  and excitation frequency  $\omega$  are chosen as control parameter. The corresponding 3-D bifurcation diagrams with respect to excitation frequency are shown in Fig.18, five different values of the preload, i.e. Preload= $F_p$ , Preload= $1.2F_p$ , Preload= $1.4F_p$ , Preload= $1.6F_p$ , and Preload= $1.8F_p$ , are

studied. With the decrease of preload, the system exhibits different kinds of motion states. The system only exhibits simple periodic-1 motion and quasi-periodic motion when the preload  $\geq 1.6F_p$ , and the main frequency in Fig. 19(d) and Fig. 19(e) is the fundamental frequency  $f$ . Decrease the preload from  $1.4F_p$  to  $1.2F_p$ , as shown in Fig. 18, the interval of quasi-periodic and chaotic motion relative to the whole excitation frequency broaden, and frequency demultiplication ( $4f/13$ ) appears in Fig.19(b) and Fig.19(c). By further decrease the preload from  $1.2F_p$  to  $F_p$ , the region of periodic- $n$ , quasi-periodic and chaotic motion continuously broaden, and continuous spectrum appears in Fig.19(a). This indicates that lower preload can degrade the dynamic performance of ball screw in some specific excitation frequency range..

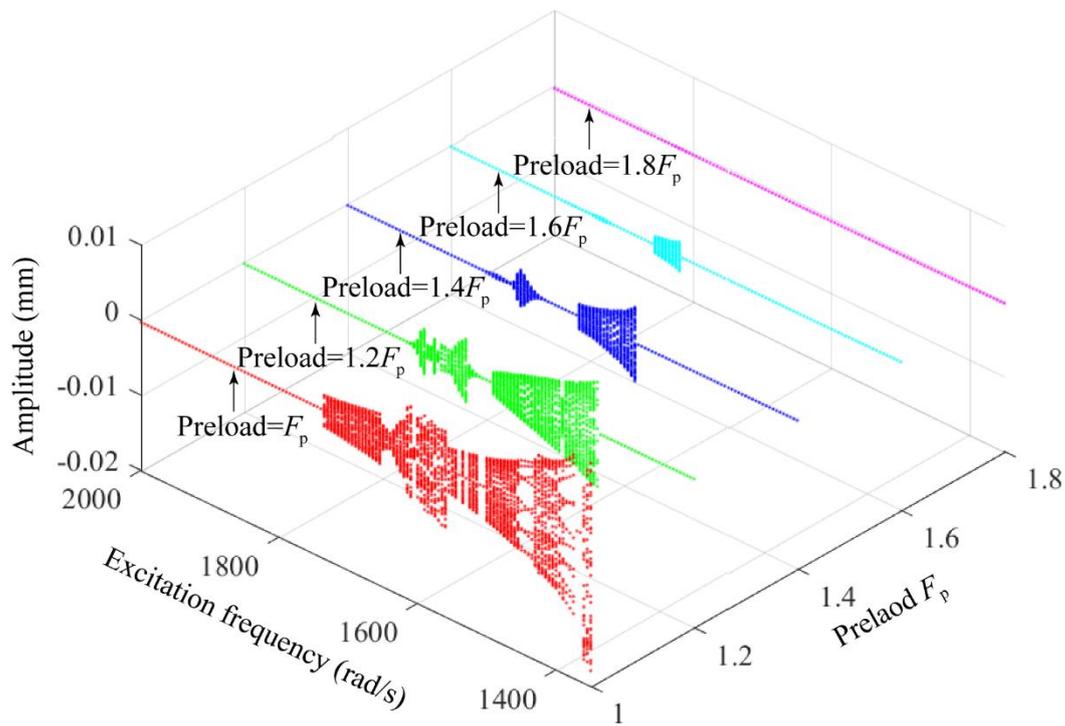


Fig.18 Bifurcation diagrams with excitation frequency  $\omega$  as the control parameter at  $N=23$ ,  $F_0=15000\text{N}$  and  $\varphi=0^\circ$  for different preload. (a) Preload= $F_p$ , (b) Preload= $1.2F_p$ , (c) Preload= $1.4F_p$ , (d) Preload= $1.6F_p$ , (e) Preload= $1.8F_p$ .

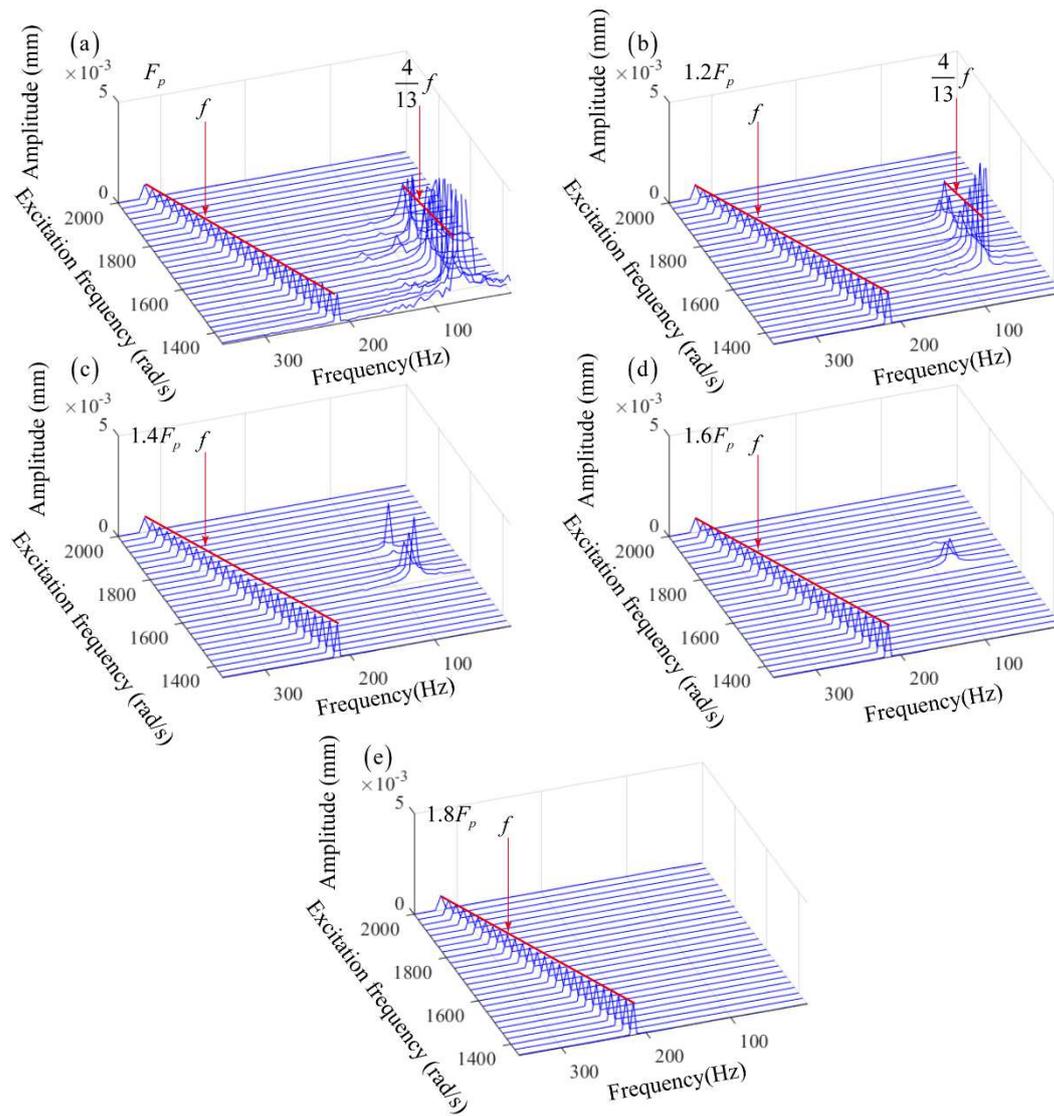


Fig.19 3-D frequency spectrum with excitation frequency  $\omega$  as control parameter at  $N=23$ ,  $F_0=15000\text{N}$  and  $\varphi=0^\circ$  for different preload. (a) Preload= $F_p$ , (b) Preload= $1.2F_p$ , (c) Preload= $1.4F_p$ , (d) Preload= $1.6F_p$ , (e) Preload= $1.8F_p$ .

As shown in Fig.20, in order to further investigate the influence of preload, the dynamic response at  $\omega=1585\text{rad/s}$  and  $\omega=1585\text{rad/s}$  are exhibited with preload as control parameter. It can be seen from Fig. 20(a1) and Fig.20(b1), the system exhibit chaotic motion in the range of  $T_1$ . The corresponding 3-D frequency spectrum shows the fundamental frequency component  $f$  is the second largest, and continuous frequency component appears in the interval  $T_1$ . As the increase of preload, the system enters to quasi-periodic motion in the range of  $T_2$ , except at a some points where the motion of the system is periodic-1, the continuous frequency component disappeared, and  $3f/10$  becomes the dominant frequency component. Increasing the preload further, the system presents simple periodic-1 motion in the range of  $T_3$ , and the main frequency in 3-D spectrum is the fundamental frequency  $f$ . Similarly in Fig. 20(a2) and Fig. 20(b2), after experiencing quasi-periodic motion in the range of  $T_1$ , the system enters to periodic-1 motion in the interval  $T_2$ , and the main frequency changed from frequency demultiplication ( $6f/25$  and  $7f/25$ ) in the range of  $T_1$  to fundamental frequency  $f$  in the range of  $T_2$ . The dynamic behavior of quasi-periodic motion at  $\omega=1680\text{rad/s}$ , preload =  $F_p$  is exhibited in Fig. 21, which presents that the system exhibits quasi-periodic motion

at this point.

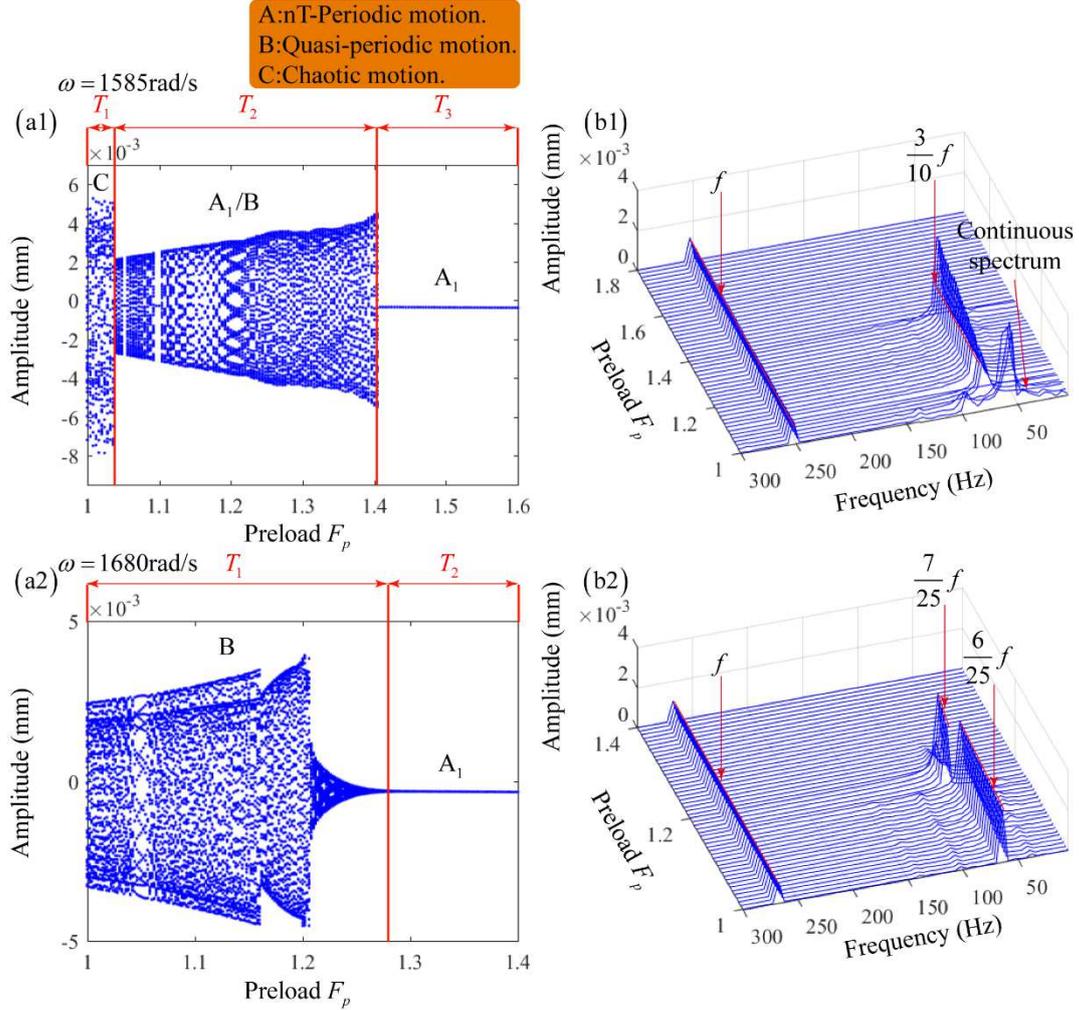


Fig.20 Dynamic responses of the system with deflection angle preload as control parameter at  $F_0=15000\text{N}$ , and  $N=23$  for  $\omega =1585\text{rad/s}$  and  $\omega=1680\text{rad/s}$ . (a) Bifurcation diagram; (b) 3-D frequency spectrum.

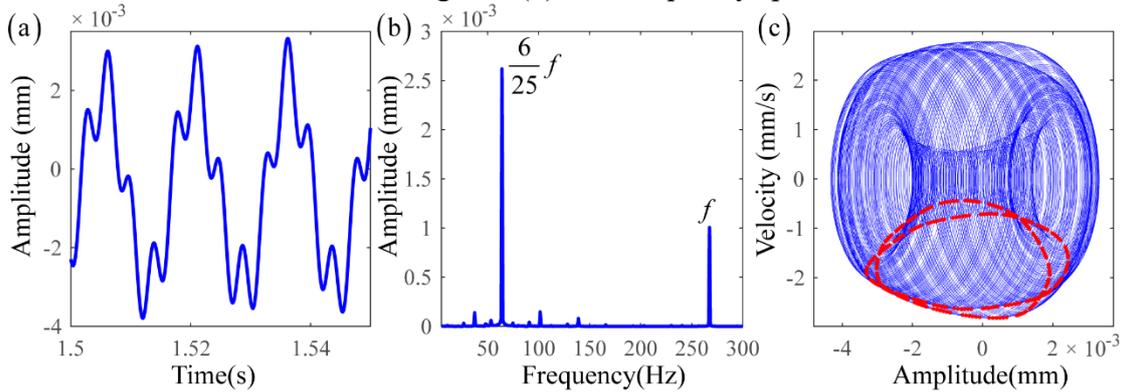


Fig.21 Vibration responses at Preload= $F_p$ ,  $N=23$ ,  $F_0=15000\text{N}$  and  $\varphi=0^\circ$  for  $\omega=1680\text{rad/s}$ . (a) Waveform, (b) Frequency, (c) Phase diagram and Poincaré section.

#### 4 Experimental verification

To validate the proposed dynamic model and obtain the dynamic parameters used in previous section, an experiment is setup in this section. The top view schematic of experiment setup can be shown in Fig. 22. As seen in the figure, the ball screw (THK SBN4016) is fixed on a test rig. The parameters of ball screw are listed in Table. 1. In the experiment, the deflection angle of screw shaft is generated by a hydraulic jack

acting on the right end of screw shaft, and the dial gauge is used to measure a the static displacement in circumferential direction. The vibration response is measured by the accelerometer (Sinocera CA-YD-189) mounted on the right end of screw shaft. Using an electromagnetic shaker (Sinocera JZK- 50) to generate harmonic excitation, and a piezoelectric force sensor (Sinocera CL-YD-331A) is used to measure the excitation amplitude acting on screw shaft along  $z$  axis. As shown in the figure, the signal collection and generating system consist of a power amplifier (YE5874A), a signal generator (Sinocera YE1311), a charge amplifier (Sinocera YE5874A), a data collection system (DH5956) and a PC.

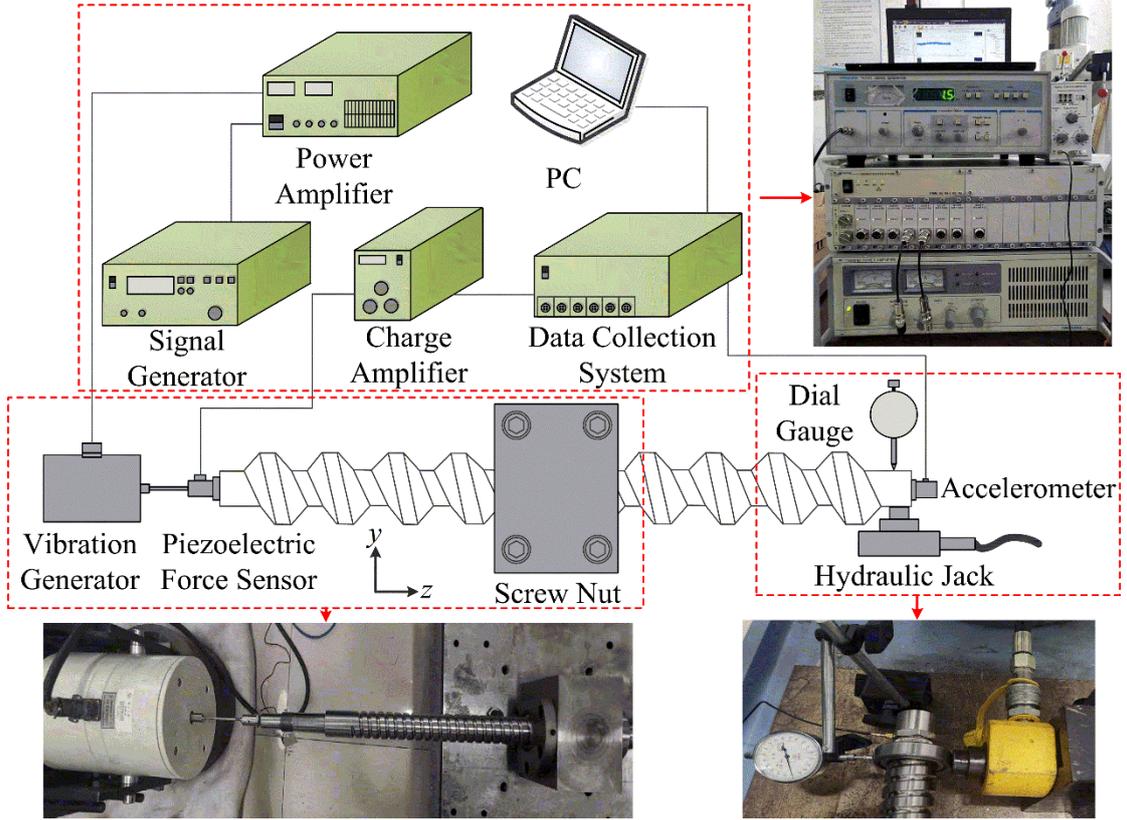


Fig. 22 Experimental setup.

In this study, the half-power bandwidth method is employed to extract damping ratios from the frequency response function. The frequency response function along  $y$  and  $z$  direction can be obtain by impact test, and the result can be shown in Fig. 23. According to the equation in Ref. [42], the expression of damping ratio  $\zeta$  can be given by:

$$\zeta = \frac{\omega_2 - \omega_1}{2\omega_{dm}} \quad 26$$

where  $\omega_{dm}$  represents the nature frequency of screw nut,  $\omega_1$  and  $\omega_2$  are the half-power frequency when the corresponding amplitude  $A=0.707A_{max}$ , where  $A_{max}$  represents the corresponding amplitude of  $\omega_{dm}$ . Therefore, the results can be shown in Fig. 23,  $\omega_{y1}=65.3\text{Hz}$ ,  $\omega_{y2}=66.935\text{Hz}$ ,  $\omega_{ydm}=65.9\text{Hz}$ ,  $\omega_{z1}=96.85\text{Hz}$ ,  $\omega_{z2}=101.2\text{Hz}$ ,  $\omega_{zdm}=102\text{Hz}$ . Based on the measurement results above, the dimensionless damping ratio along  $y$  axis  $\zeta_y=0.0124$  and  $z$  axis  $\zeta_z=0.0213$  can be estimated. The viscous damping coefficient  $c$  used in the previous section can be calculated by  $c=4\pi\zeta\omega_{dm}$ , hence the value of  $c_y$  and  $c_z$  can be evaluated. Furthermore, the amplitude frequency curve of screw nut from simulation and experiment are compared in Fig. 24. The differences of the nature frequency from experiment and simulation is 1.6%, which validate the proposed model.

In order to verify the proposed method with the consideration of deflection angle, two values of excitation frequency, i.e.  $\omega/2\pi=74\text{Hz}$  and  $\omega/2\pi=126\text{Hz}$  are selected to compare the vibration response between experiment and simulation, which can be seen in Fig.25. According to the measuring result of dial gauge, the deflection angle equals to  $0.7162^\circ$ . The comparison of vibration response between the selected two points can be shown in Fig. 25.

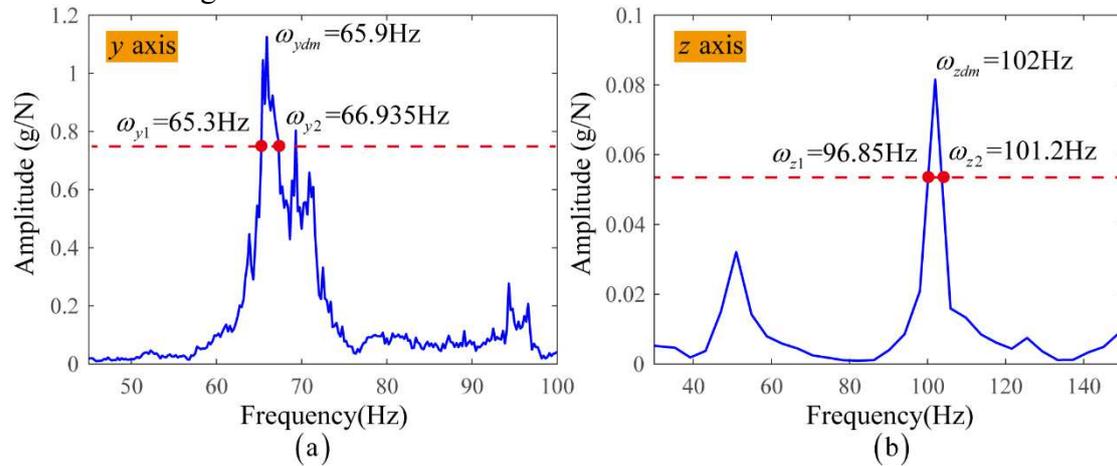


Fig. 23 Frequency response function obtained from impact test in two directions; (a) x axis. (b) z axis.

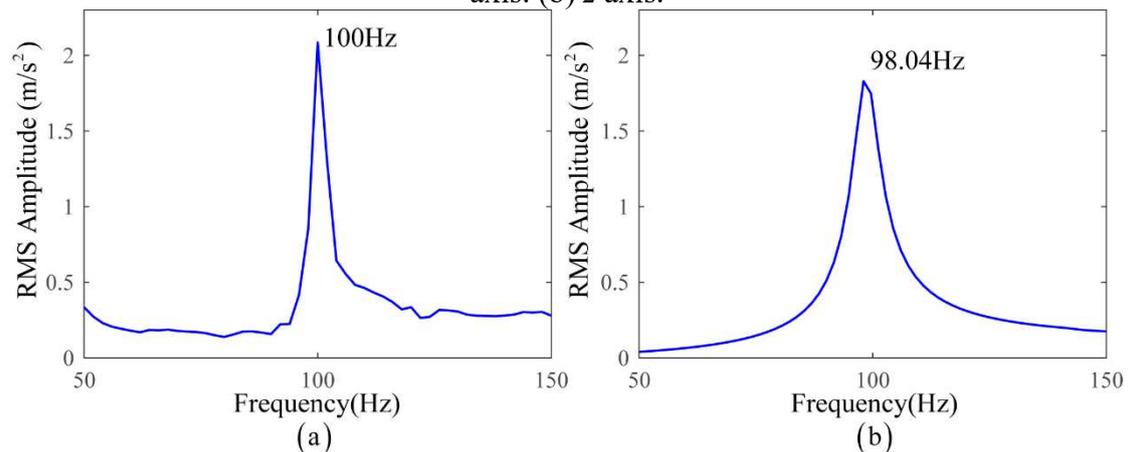


Fig. 24 Amplitude frequency curve obtained from impact test at  $F_0=50\text{N}$ . (a) Experimental result. (b) Simulation result

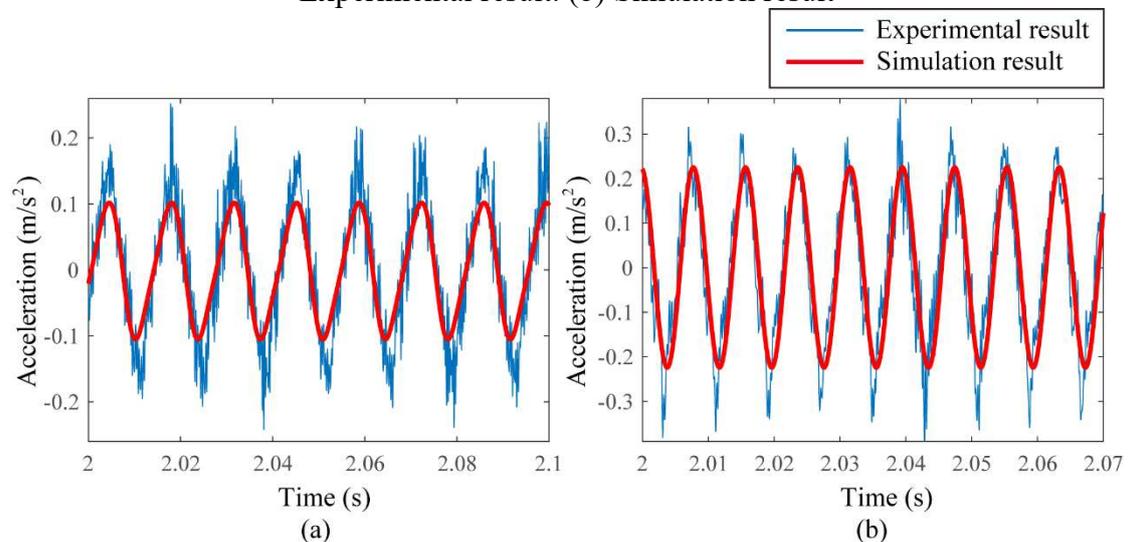


Fig. 25 Vibration response between simulation and experiment. (a)  $\omega/2\pi=74\text{Hz}$ ,

$$F_0=50\text{N. (b) } \omega/2\pi=126\text{Hz, } F_0=50\text{N.}$$

## 5. Conclusion

In this paper, a 5 degrees-of-freedom dynamic model of variable lead preloaded single nut ball screw considering the effect of deflection angle is proposed, the proposed model is investigated by numerical method. The relationship between the deflection angle and the contact deformation of each ball is given. The effects of different system parameters, i.e. excitation amplitude, deflection angle, the number of balls, and preload are quantitatively studied. Bifurcation diagram, 3-D frequency spectrum, and RMS value amplitude frequency curves with respect to different parameters are used to study the influence on dynamic behaviors. To validate the proposed dynamic model and estimate the dynamic parameters, a series of experiments are conducted. The main conclusions of this study are listed as follows:

(1) Based on different system parameters, the system presents various nonlinear dynamics, namely periodic- $n$  motion, quasi-periodic motion, chaotic motion and jump discontinuity phenomenon.

(2) In the whole range of excitation frequency, the excitation amplitude has impacts on system dynamic response. In contrast, for other system parameters, i.e. preload, the number of balls, the influenced range of excitation frequency is relatively limited.

(3) The system parameters can affect the dynamic characteristic by different ways. Excitation amplitude, preload, deflection angle can enhance the nonlinear properties of the system, and the number of balls can change the resonance frequency of the system.

## Declarations

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### Conflicts of interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

### Availability of data and material

The data sets supporting the results of this article are included within the article and its additional files.

### Authors' contributions

Zhendong Liu: Methodology; Investigation; Experimental; Writing-Original Draft; Writing- Review & Editing

Mengtao Xu: Resources and supervision

Hongzhuang Zhang: Resources, Writing- reviewing and editing, Supervision, Writing- Review & Editing

Zhenyuan Li: Carried out the experiment

Changyou Li: Conceived the presented idea

Guo Yao: Conceived the presented idea

Yimin Zhang: Resources and supervision

### Ethics approval

This chapter does not contain any studies with human participants or animals performed by any of the authors.

### Consent to participate

Not applicable. The article involves no studies on humans.

## Consent for publication

All authors have read and agreed to the published version of the manuscript.

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