

2 μm Passively Mode-Locked Thulium-doped Fiber Lasers with Ta₂AlC-deposited Tapered, Side-polished and D-shaped Fibers

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Abstract

In this work, mode-locked thulium-doped fiber lasers operating in the 2 μm wavelength region were demonstrated using tantalum aluminum carbide (Ta_2AlC)-based saturable absorbers (SAs) utilizing the evanescent wave interaction. The Ta_2AlC MAX Phase was prepared by dissolving the Ta_2AlC powder in isopropyl alcohol (IPA) and then deposited onto three different evanescent field-based devices, which were the tapered fiber, side-polished fiber (SPF), and D-shaped fiber. Flame-brushing and wheel-polishing techniques were used to fabricate the tapered and D-shaped fibers, respectively, while the side-polished fiber was purchased commercially. All three SA devices generated stable mode-locked pulses at center wavelengths of 1937, 1931, and 1929 nm for the tapered, side-polished, and D-shaped fibers. The frequency of the mode-locked pulses was 10.73 MHz for the tapered fiber, 9.58 MHz for the side-polished fiber, and 10.16 MHz for the D-shaped fiber. The measured pulse widths were 1.678, 1.734, and 1.817 ps for each of the three SA devices. The long-term stability of the mode-locked lasers was tested for each configuration over a 2-hour duration. The lasers also showed little to no fluctuations in the center wavelengths and the peak optical intensities, demonstrating a reliable, ultrafast laser system.

Introduction

The discovery of optical fibers by Charles Kao and George A. Hockam in 1966 gave the inroads for optical amplifier development with a wavelength range of 1.46 μm to 1.53 μm ^{1,2}, 1.53 μm to 1.565 μm ^{3,4} and 1.565 μm to 1.625 μm ⁵⁻¹⁰. Consequently, the development of laser configuration such as pulsed^{11,12}, dual and multiwavelength¹³⁻²², and optical sensors²³⁻²⁵ has been the focus of many research laboratories. There have been numerous activities on pulsed laser operating at wavelengths 1 μm and 1.5 μm , of interest will be to generate short pulses in the 2 μm wavelength region for several applications such as in spectroscopy²⁶, gas detection²⁷, laser ablation²⁸, long-range light detection and ranging (LIDAR)²⁹, plastic and glass processing³⁰ as well as in the medical field³¹. Lasing in the 2 μm is commonly achieved using thulium-doped fibers (TDFs) as the gain medium, as TDFs have a broad amplification range of 400 nm, ranging from 1700 to 2100 nm³². The 2 μm wavelength region is also of interest as it coincides with the absorption lines of water (H_2O) and several leading greenhouse gases such as carbon dioxide (CO_2) and nitrogen dioxide (NO_2)³³. Although lasing in the 2 μm region has traditionally been demonstrated with continuous wave (CW) outputs, recent advances in fiber laser technologies have increased the development of 2 μm pulsed fiber lasers that can generate short pulses with pulse durations in the pico- or femtosecond range.

Pulsed laser generation can be achieved either by Q-switching or mode-locking. In the former, short pulses with high output energies can be produced using an optical component incorporated in the laser cavity to modulate the Q-factor. In the latter, the oscillating longitudinal modes present in the laser cavity are phase locked when an optical component is introduced in the optical cavity. Pulse generation in fiber lasers could be obtained using two main techniques: active and passive techniques³⁴. Active techniques require

the use of external modulators such as acousto-optic and electro-optic modulators³⁵. It, however, causes the system to be bulky and inflexible due to the extra electronic components needed to be used.

In comparison, a passive technique allows for the development of a more compact and versatile system. Saturable absorbers (SAs) are used to saturate the molecules or atoms, whereby the optical absorption decreases as the light intensity increases. Due to this nonlinear optical response of SAs together with a narrow optical bandgap, a high damage threshold, and a wide bandwidth, SAs are suitable devices to generate short pulses using the Q-switching and mode-locking techniques. SAs can be divided into two groups, namely artificial SAs and real SAs. Nonlinear optical loop mirrors (NOLMs), nonlinear amplification loop mirrors (NALMs), or nonlinear polarization evolution (NPE) are examples of artificial SAs. Artificial SAs are not suitable for commercialization due to its sensitivity to environmental changes and large size despite their positive attributes of near-instantaneous response time and high modulation depth³⁶⁻³⁹. Semiconductors saturable absorbers mirrors (SESAM), a real SA, were chosen as the SA of choice for nonlinear absorption property that depends on light intensity^{40,41}. However, the disadvantages of SESAM are the operating bandwidth is narrow, complex design, costly, and has a low-damage threshold⁴².

In view of those limitations mentioned above, new SA nanomaterials are now the main focus of research in ultrafast laser worldwide. To date, various kinds of materials that exhibit intensity-dependent transmission have been used as SAs, namely graphene⁴³, carbon nanotubes (CNTs)⁴⁴, black phosphorus (BP)⁴⁵, transition metal dichalcogenides (TMDs)⁴⁶, and topological insulators (TIs)⁴⁷. Recently, a new type of material named MXenes has been widely explored for various optoelectronic applications due to their unique optical properties^{48,49}. It also makes them a great candidate to be used as SAs in generating ultrafast lasers. MXenes are typically obtained from their precursor, the MAX phases⁵⁰. Compared to its counterpart, MAX phases are favorable as it does not require the use of strong etching solutions that contain fluoride ions (F^-) such as hydrofluoric acid (HF), ammonium bifluoride (NH_4HF_2), or a mixture of hydrochloric acid (HCl) and lithium fluoride (LiF), thus minimizing the fabrication process and cost⁵¹. The MAX phases are also useful for high-temperature applications as they comprise of ternary transition-metal carbides that have metal and ceramic properties. This makes them to have a good thermal and electrical conductivity, as well as having a high damage threshold⁵². Several works demonstrate the use of MAX phases as SAs in generating pulses in fiber lasers. For instance, Lee et al. demonstrated the use of a titanium aluminum carbide (Ti_2AlC) SA to generate Q-switched pulses with a maximum pulse energy of 22.58 nJ in an erbium-doped fiber laser (EDFL) cavity⁵³. Jafry et al. also demonstrated an ultrashort pulse generation in an EDFL cavity using a Ti_3AlC_2 -PVA SA. The generated mode-locked pulses had a pulse width of 3.68 ps wavelength of 1577 nm⁵⁴. These demonstrations show the tremendous potential of MAX phases in generating short pulses in fiber lasers, which would allow further exploration of various MAX phases with other combination of the early transition metals.

In addition to these aforementioned SA materials, the structure of the SA devices has a significant impact on the SA performance. SA materials could be prepared and integrated into fiber laser cavities by several

arrangements, which are typically done using the optical deposition method onto fiber ferrules⁵⁵ or substrates and polymer hosts^{56,57}. Although these methods allow for a more direct integration, they limit the operation of the lasers at low power due to their poor heat dissipation and low optical damage thresholds. Another method that is attracting a great interest of late is the evanescent field-based devices, which utilize the nonlinear interaction of the propagating light in the optical fiber with the SA materials. Since the light-matter interaction is realized via the evanescent field, this approach is more efficient. It can provide better functionality by eliminating the issue of heat accumulation. Various types of evanescent field-based devices have been used as to generate pulses in fiber lasers. For example, Mouchel et al. utilized a graphene-coated tapered fiber to generate mode-locked pulses in an Er:Yb doped double-clad fiber laser with a high average output power of 520 mW when being pumped to a maximum pump power of 5 W⁵⁸. Zhou et al. obtained a mode-locked laser using a lead sulfide (PbS) being deposited onto a D-shaped fiber, in which the SA had a damage threshold of higher than 3.5 kW⁵⁹. Other reports have also shown the use of side-polished fibers (SPFs) and nonlinear materials to generate mode-locked pulses. Li et al. reported a harmonic mode-locking at 1879 nm using a graphene-deposited SPF as the SA, where up to the 21st harmonics were obtained at a pump power of 500 mW⁶⁰. Khazaeinezhad et al. integrated a molybdenum disulfide (MoS₂) SPF in an EDFL cavity to get mode-locked fiber lasers at both anomalous and normal-dispersion regimes, with the mode-locking operating being sustained at a pump power of more than 600 mW⁶¹. These demonstrations have shown that the use of evanescent field-based devices as saturable absorbers would be suitable in developing robust ultrafast laser systems to suit the demand for future high-power laser systems.

In this work, we explored three different types of evanescent field-based devices: the tapered fiber, the side-polished fiber, and the D-shaped fiber to generate mode-locked pulses in the 2 μm wavelength region. A Ta₂AlC MAX phase was first prepared in the solution form by ultrasonication Ta₂AlC powder in isopropyl alcohol (IPA), deposited onto the three fibers. The MAX phase was composed of tantalum (Ta) as the early transition metal instead of the common titanium (Ti). Each of the SA devices was then individually inserted into a TDFL cavity to generate mode-locked pulses with frequencies between 9 and 11 MHz and pulse widths between 1.678 and 1.817 ps. The pump power of the laser cavity could be increased up to 1 W with all three devices, maintaining the mode-locking operation in the 2 μm without any damage to the SA devices. The results show the role of the evanescent field-based devices as promising and robust SA devices for the development of high-power fiber lasers.

Characterization Of The TaalC Max Phase Solution

The crystalline phase of the Ta₂AlC MAX Phase was first assessed using X-ray powder diffraction (XRD). The measurement was recorded using a Malvern Panalytical Empyrean XRD by utilizing Cu-Kα radiation (λ = 0.1541 nm) as the X-ray source and by scanning in the 2θ range of 10° to 80°. The obtained XRD spectrum is shown in Fig. 1.

It was observed from the XRD pattern that a distinct prominent peak at $2\theta = 39.07^\circ$ together with a few sharp peaks was detected at $2\theta = 12.92, 25.84, 33.72, 34.34, 35.25, 39.07, 40.81, 42.84, 51.63, 52.46, 58.08, 60.06, 61.72, 66.48, 69.06, 70.61, \text{ and } 73.96^\circ$. These peaks were assigned to the (1 0 3), (0 0 2), (0 0 4), (1 0 0), (1 0 1), (1 0 2), (1 0 4), (1 0 5), (1 0 6), (0 0 8), (1 0 7), (1 1 0), (1 1 2), (1 1 4), (1 0 9), (2 0 1) and (2 0 3) planes of the Ta_2AlC MAX Phase, respectively. The result was in accordance with the ICDD card No. 29-0092 for Ta_2AlC ⁶² and similarly reported in previous works^{63,64}. The absence of other impurity peaks indicates the purity of the sample, and the high crystallinity of the sample was proven by the presence of sharp diffraction peaks in the XRD pattern of the Ta_2AlC MAX Phase.

The morphology of the Ta_2AlC MAX Phase was then examined using a JEOL JSM 7600-F Field Emission Electron Microscope (FESEM) at an accelerating voltage of 5.0 kV. The FESEM images captured at various magnifications are shown in Fig. 2. The Ta_2AlC MAX Phase had a flake-like structure with irregular shapes and had different lateral flake sizes varying from about 0.5 μm to 18 μm , as shown in Fig. 2 (a). The single grain of Ta_2AlC in Fig. 2 (b) indicates that the Ta_2AlC consisted of assembled Ta-C layers and Al layers that are alternately stacked to form the Ta_2AlC . At the highest magnification shown in Fig. 2 (c), it is seen that the layers are densely packed and that a single layer of the Ta_2AlC had an average thickness of less than 100 nm.

Fabrication and Characterization of the Ta_2AlC -deposited Tapered, D-shaped and Side Polished Fibers

Ta_2AlC -deposited Tapered Fiber

The tapered fiber was fabricated using the flame brushing method, as shown in Fig. 3 (a). An approximately 0.6 m-long bare fiber was first stripped and then fixed onto Newport FCL100 translation stages using fiber holders. The stages were used to pull and stretch the fiber after it has been softened using oxy-LPG flames so that the width of the optical fiber was reduced from a diameter of 125 μm to 6 μm , with a tapered length of about 2.5 cm. Upon completion, the insertion loss of the tapered fiber was measured using a light source (LS) and an optical power meter (OPM), giving a value of approximately 4 dB at 2000 nm.

The Ta_2AlC in the solution form was then deposited onto the tapered fiber using the drop-cast method and was left until dry at room temperature. The microscopic image of the fabricated tapered fiber taken at 50 times magnification is shown in Fig. 3 (b), where the tapered waist was reduced to only 6 μm . In Fig. 3 (c), the microscopic image of the tapered fiber after the deposition of the Ta_2AlC was also taken at 100 times magnification. It shows that the Ta_2AlC was successfully coated around the waist of the fabricated tapered fiber. The graphical illustration of the MAX phase Ta_2AlC -coated tapered fiber is also given in Fig. 3 (d).

The twin detector measurement technique⁶⁵ was used to evaluate the nonlinear optical absorption properties of the Ta₂AlC-deposited tapered fiber. The pulsed laser for the measurement was a FemtoFERb 1950 nm femtosecond fiber laser, from Toptica Photonics, with a pulse width of 90 fs and a repetition rate of 30 MHz. To calculate the modulation depth, the non-saturable absorption and the saturation intensity of the Ta₂AlC-deposited tapered fiber, the experimental data obtained were fitted using the typical two-level saturation absorption model, which can be expressed as follows⁶⁶:

$$\alpha(I) = \frac{\alpha_s}{1+I/I_{sat}} + \alpha_{ns} \quad (1)$$

where I is the input intensity, I_{sat} is the saturation intensity and α_{ns} is the non-saturable absorption. For the Ta₂AlC-deposited tapered fiber, the computed modulation depth and the saturation intensity as measured in Fig. 4 were 6.02% and 0.36 MW/cm², respectively.

Ta₂AlC-deposited Side-Polished Fiber

The side-polished fiber (SPF) was commercially obtained from Phoenix Photonics. The polishing depth of the SPF was approximately 1 μm from the edge of the core, and the polished length was 1.7 cm, as given by the manufacturer. The Ta₂AlC was drop-casted onto the polished region of the SPF, similarly as it was with the tapered fiber. Fig. 5 (a) shows the image of the Ta₂AlC-deposited SPF when being injected with a red-light source. It is observed that scattered light was seen along the polished region of the SPF, confirming that the polished length was approximately 1.7 cm. The deposited Ta₂AlC at the polished part of the SPF is illustrated in Fig. 5 (b), while the microscopic image of the deposited polished region of the SPF at 100 times magnification is shown in Fig. 5 (c). The measured insertion loss for the side polished fiber Ta₂AlC was 4.8 dB at 2000 nm. The nonlinear absorption test for the Ta₂AlC-deposited side polished fiber was conducted using the twin detection measurement method as described previously. From Fig. 6, the modulation depth and the saturation intensity were calculated to be 1.09 % and 1.63 MW/cm² for the Ta₂AlC-deposited SPF.

Ta₂AlC-deposited D-shaped Fiber

The D-shaped fiber was fabricated using the polishing wheel method⁶⁷, schematically shown in Fig. 7 (a). Two fiber holders and two mechanical alignment stages were used to hold a single-mode fiber (SMF-28) above a polishing wheel with a diameter of approximately 1.5 cm. The polishing wheel was wrapped with P800 grit sandpaper and then fixed on a motor shaft. The wheel was rotated and gradually raised onto the bottom side of the SMF. The polishing process was carried out while the insertion loss of the D-shaped fiber was simultaneously monitored by connecting one end of the fiber to a light source (LS) and

the other to an optical power meter (OPM). The polishing process was stopped when the insertion loss of the D-shaped fiber was measured to be approximately 4.2 dB. After the process was completed, the D-shaped fiber was carefully transferred onto a glass slide. The deposition of the Ta₂AlC solution onto the D-shaped fiber was carried out in the same manner as described in the previous section. When injected with a red-light source, scattered light was observed in the polished region of the D-shaped fiber, as shown in Fig. 7 (b). The microscopic image of the D-shaped fiber's side profile at 50 times magnification, as shown in Fig. 7 (c) and Fig. 7 (d), illustrates the deposited Ta₂AlC onto the D-shaped fiber. Fig. 8 shows the nonlinear absorption plot of the Ta₂AlC-deposited D-shaped fiber. From the graph, the modulation depth and the saturation intensity were calculated to be 0.82 % and 1.03 MW/cm², respectively.

Table 1

Summary of the nonlinear properties of each of the SA devices

SA Device	Modulation Depth (%)	Saturation Intensity (MW/cm ²)
Ta ₂ AlC-deposited tapered fiber	6.02	0.36
Ta ₂ AlC-deposited SPF	1.06	1.63
Ta ₂ AlC-deposited D-shaped fiber	0.82	1.03

The modulation depths and the saturation intensities of the Ta₂AlC-based SA devices are summarized in Table 1. The Ta₂AlC-deposited tapered fiber had the lowest saturation intensity of 0.36 MW/cm², with a modulation depth of 6.02%. Although the Ta₂AlC-deposited SPF and D-shaped fiber had lower modulation depth at 1.06% and 0.82%, the modulation depth of 6.02% was still capable of generating mode-locked pulses. Previous works have been reported using SAs with modulation depths of above 5%⁶⁸⁻⁷⁰. It was expected that Ta₂AlC-deposited tapered fiber would have the lowest mode-locking threshold, as it had the lowest saturation intensity compared to the 1.63% of the Ta₂AlC-deposited SPF and 1.03% of the Ta₂AlC-deposited D-shaped fiber. A low saturation intensity would allow the SA to be saturated at a lower pump power to induce the mode-locking operation⁷¹.

Experimental Setup

The experimental setup for the 2 μm mode-locked thulium-doped fiber laser (TDFL) using the Ta₂AlC-based SAs is presented in Fig. 9. A 1565 nm laser source (LS) with a maximum output power of 1 W was used to pump the gain medium through the 1550 nm port of the 1550/2000 nm wavelength division

multiplexer (WDM). The gain medium was a 4-meter long TmDF200 thulium-doped fiber (TDF) from OFS Inc, connected to the common port of WDM₁. The absorption of the TDF was 22 dB/m at 1565 nm, and the core and cladding diameters of the TDF were 9 and 125 μm, respectively. The TDF was then connected to the common port of WDM₂, where the 1550 nm port was used to remove any excess pump. The 2000 nm port of WDM₂ was connected to the input port of a 2000 nm isolator to obtain a unidirectional operation in the laser cavity. A polarization controller (PC) was connected to the isolator's output to adjust the propagating signal's polarization state. The Ta₂AlC-based SA, either in tapered fiber, a side-polished fiber, or a D-shaped fiber, was connected to the other end of the PC. The cavity was completed by connecting the other end of the SA device to a 90:10 coupler, where 90% of the signal was circulated back into the cavity, and another 10% was taken as the laser output.

Since there were three different types of SA devices used in this experiment, the total length for each fiber was slightly different, resulting in a different net cavity dispersion for each setup. Table 2 gives the calculated group velocity dispersion (GVD) values for each SA device at their respective operating wavelengths. The manufacturer provided the GVDs of the TmDF200, while the GVDs of SMF-28 were estimated from Corning dispersion equation.

Table 2
The computed dispersion values for the TDFL cavities with different types of SA devices.

Device	Central wavelength (nm)	GVD TmDF 200 (ps ² /m)	TmDF200 length (m)	GVD SMF-28 (ps ² /m)	SMF-28 length (m)	Total cavity length (m)	Net dispersion (ps ²)
Tapered	1937	-0.214	4	-0.0649	15.35	19.35	-1.85
SPF	1931	-0.211	4	-0.0645	17.81	21.81	-1.99
D-shaped	1929	-0.208	4	-0.0643	16.43	20.43	-1.88

The net cavity dispersion for each of the mode-locked laser was computed by

$$Dispersion_{net} = GVD_{TmDF200} Length_{TmDF200} + GVD_{SMF} Length_{SMF}.$$

From the values given in Table 2, the calculated net cavity dispersions for the mode-locked lasers were -1.85, -1.99, and -1.88 ps² for the cavity with the tapered fiber, SPF, and D-shaped, respectively. It indicates that all the mode-locked lasers operated in the anomalous dispersion regime.

Results And Discussions

The experiment was first conducted without any of the SA devices to observe the operation of the fiber laser. As expected, the TDFL without any SA only operated in the continuous wave (CW) regime as no

pulses were observed in the oscilloscope as the pump power was increased to a maximum of 1 W. The experiment was then continued by integrating the Ta₂AlC-deposited tapered fiber into the laser cavity. With fine-tuning of the PC, the fundamental mode-locking of the TDFL was observed at a threshold pump power of 245 mW. The fundamental mode-locking could be obtained until the pump power reached 480 mW. As the pump power was further increased until the maximum, the mode-locked operation was sustained but the laser operated at higher harmonics. As our interest was mainly on the fundamental operation of the mode-locked laser, the laser characteristics were only recorded when the pulsed laser operated at the fundamental frequency.

The characteristics of the mode-locked laser using the Ta₂AlC-deposited tapered fiber are shown in Fig. 10. The optical spectrum recorded in Fig. 10 (a) shows a broad laser spectrum at a center wavelength of 1937 nm, having a 3-dB bandwidth of 2.8 nm. It was apparent that Kelly sidebands were observed in the soliton spectrum when the mode-locking operation was initiated, which was expected as the mode-locked TDFL was operating in the anomalous dispersion regime. From the values obtained from the optical spectrum, the transform-limited pulse width could be estimated by the equation;

$$\tau_p = \frac{TBP \times (\lambda_c)^2}{c \times \Delta\lambda} \quad (2)$$

where c is the speed of light, TBP is the time-bandwidth product, λ_c is the center wavelength, and $\Delta\lambda$ is the 3-dB bandwidth. Taking the TBP to be 0.315 for a sech^2 pulse profile, the center wavelength to be 1937 nm, and the 3-dB bandwidth to be 2.8 nm, the transform-limited pulse width was calculated to be 1.407 ps. Figure 10 (b) shows that the oscilloscope trace of the mode-locked pulses had a repetition rate of 10.73 MHz, which correlates with the cavity round trip time estimated from the cavity length of approximately 19.3 m. A comparable pulse train was achieved by Zhou et al.⁵⁹, where the slight fluctuation of the pulse peak intensities was affected by the relatively low sampling rate of the photodetector bandwidth and the low sampling points per pulse. Nonetheless, it was found that the pulse train's amplitude jitter was within reasonable limits. The radio frequency (RF) spectrum of the mode-locked pulse is shown in Fig. 10 (c), whereby a sharp peak was observed at about 10.73 MHz with a measured signal-to-noise ratio (SNR) of ~ 55 dB. Figure 10 (d) shows the autocorrelation trace of the modelocked pulse, where the pulse width was measured to be 1.678 ps when fitted with a sech^2 fitting. It was only about 19% longer than the calculated transform-limited pulse width. The corresponding TBP was 0.375, indicating that the pulse width is only slightly chirped.

The experiment was then further conducted by replacing the Ta₂AlC-deposited tapered fiber with the Ta₂AlC-based side-polished fiber that has been described in Sect. 3.3. The fundamental mode-locking operation was achieved at a threshold pump power of 351 mW. The laser spectrum is given in Fig. 11 (a) also shows a typical soliton spectrum with distinct Kelly sidebands. However, it was observed that the sidebands were uneven, with the longer being higher than the shorter-wavelength sidebands. It is highly

likely due to the optical fiber birefringence filtering effect in the cavity, as was theoretically and experimentally proven by Man et al.⁷². The center wavelength and the 3-dB bandwidth of the TDFL were recorded to be 1931 nm and 3.1 nm, respectively. As for the frequency of the mode-locked pulses, the oscilloscope trace shown in Fig. 11 (b) shows a repetition rate of 9.52 MHz. The frequency was lower compared to the mode-locked laser with the Ta₂AlC-deposited tapered fiber, which was due to the slightly longer length of the SPF. The RF spectrum of the mode-locked laser in Fig. 11 (c) shows a sharp peak at around 9.52 MHz, having an SNR of ~ 50.5 dB. From the autocorrelation trace of the mode-locked pulse shown in Fig. 11 (d), the pulse width was measured to be 1.743 ps, fitted with a sech² profile. The corresponding TBP was then calculated to be 0.434, also indicating that the pulse was chirped.

The TDFL cavity was further tested by inserting the fabricated D-shaped fiber with the Ta₂AlC solution. The mode-locked pulses were generated at a pump power of 380 mW. The output characteristics of the mode-locked laser are shown in Fig. 12. From the optical spectrum plotted in Fig. 12 (a), the mode-locked laser had a center wavelength of 1929 nm with a 3-dB bandwidth of 2.2 nm. The presence of minor dips in the optical spectrum could be attributed to water absorption lines in 2 μm⁷³. The mode-locked pulse had a repetition rate of 10.16 MHz with a 9.84 ns interval between peaks, measured from the oscilloscope trace in Fig. 12 (b). It tallies well with the fundamental frequency of the mode-locked laser, which was estimated by $f = c/nL$ where c is the speed of light, n is the refractive index of an optical fiber, and L is the length of the cavity. By taking n to be approximately 1.44 at 2000 nm and L to be 20.43 meters, the fundamental frequency was about 10.2 MHz. Figure 12 (c) shows the RF spectrum with a peak that corresponds to the repetition rate of the mode-locked laser, in which it has an SNR value of ~ 47 dB. The autocorrelation trace of the pulse recorded using the autocorrelator is shown in Fig. 12 (d). When fitted with a sech² fitting curve, the pulse width was measured to be 1.817 ps.

The stability of the mode-locking operation with all three SA devices was evaluated by conducting a long-term stability test over two hours. The laser output for each of the three SA devices was monitored at every 10-minute interval, in which their optical spectrum was recorded and plotted in Fig. 13. As seen from Fig. 13 (a), the mode-locked TDFL with the Ta₂AlC-deposited tapered fiber exhibit a steady output intensity, and the contour plot shows no changes in the central wavelength of the output spectrum. For Ta₂AlC-deposited SPF, the output spectrum displayed in Fig. 13 (b) also shows a very stable output as the central wavelength and the Kelly sideband exhibit no changes throughout the stability test. Meanwhile, the D-shaped fiber demonstrates a steady output as depicted in Fig. 13 (c).

The pump power against the output power of the mode-locked lasers is plotted in Fig. 14. As seen from the graph, the average output power of the laser with each type of SA device increased almost linearly after the mode-locking operation was initiated. The mode-locking threshold was 245 mW for the TDFL with the Ta₂AlC-deposited tapered fiber, 351 mW for the Ta₂AlC-deposited SPF, and 380 mW the Ta₂AlC-deposited D-shaped fiber. At this pump power, the fundamental mode-locking (FML) was observed. The FML operation was sustained until a pump power of 480 mW for the Ta₂AlC-deposited tapered fiber and 479 mW and 550 mW for the Ta₂AlC-deposited SPF and D-shaped fiber, respectively. At the maximum

pump power in which the FML was sustained, the average output power recorded for the mode-locked TDFL with the tapered fiber, SPF, side polished, and D-shaped fiber was 1.91 mW, 0.8 mW 1.37 mW, respectively. When the pump power was further increased beyond these pump power, harmonic mode-locking was observed and could be maintained up until the maximum pump power of 1 W. At the maximum pump power of 1 W, the maximum average output power obtained were 3.48 mW, 2.27 mW and 2.71 mW for the TDFL with the tapered fiber, SPF and D-shaped fiber, respectively. The TDFL with the Ta₂AlC-deposited tapered fiber had the highest output power as the structure of the tapered fiber was maintained, only its dimension was reduced⁷⁴.

In contrast, the mode-locked TDFL with the Ta₂AlC-deposited SPF and D-shaped fiber had a lower output power since the fiber structure had been modified during the grinding or polishing process. It caused a higher amount of light to escape easily by scattering light due to imperfection of the surfaces. Nevertheless, all the three SA devices could operate even when the pump power was increased until a maximum pump power of 1 W, without being optically damaged. Compared with materials embedded in polymer hosts, SAs in the film form typically had a lower damage threshold and can be easily burnt when the power was high⁷⁵. It limits the application of polymer-based SAs in ultrafast fiber laser systems, eliminating the possibility of power-scaling of fiber lasers.

The performance of the mode-locked TDFLs with each of the three SA devices is summarized in Table 3.

Table 3
Summary of the optical characteristics of the TDFL with each mode-locking device in the fundamental operation.

Type of device	Threshold pump power (mW)	Repetition rate (MHz)	Signal to noise ratio (SNR)	Pulse width (ps)	Max. Output power (mW)	Max. Peak Power (W)
Tapered	245	10.73	55	1.678	1.91	106
SPF	351	9.52	50.5	1.743	0.8	43
D-shaped	380	10.16	47	1.817	1.37	82

From Table 3, the lowest pump power needed to initiate the mode-locking operation was that of the Ta₂AlC-deposited tapered fiber at 245 mW. This was followed by the Ta₂AlC-deposited SPF and the D-shaped fiber at the pump power of 351 mW and 380 mW, respectively. A low pump power threshold to induce the mode-locking operation was favorable as it could reduce the energy consumption. The mode-locked TDFL using the Ta₂AlC-deposited tapered fiber generated the highest maximum average output power, which was as high as 3.48 mW. Compared to the output power of the TDFL using the Ta₂AlC-deposited SPF and D-shaped fiber, the generated output power was only 2.71 mW for the former and 2.27 mW for the latter. It was lower by about 22% from the output power generated by the TDFL with the Ta₂AlC-deposited tapered fiber. The corresponding peak power could be calculated by dividing the

average output power with the repetition rate and then dividing the value with the pulse width. It gives peak power values of 106 W, 43 W, and 82 W for the TDFL with the Ta₂AlC-deposited tapered fiber, SPF, and D-shaped fiber. The lowest and highest repetition rate recorded was 9.52 MHz with the Ti₂AlC-deposited SPF and 10.73 MHz with the Ta₂AlC-deposited tapered fiber. The difference in the repetition rate was only due to the length of the fiber used for the evanescent field-based fibers. In this regard, a higher repetition rate of the mode-locked TDFL could be obtained by having a shorter cavity length. The SNR values recorded from the RF spectrum for all three cases were more than 47 dB, which indicates that all three evanescent field-based SAs could generate stable mode-locked pulses. The shortest pulse width was recorded to be 1.678 ps, obtained using the Ta₂AlC-deposited tapered fiber. The pulse widths obtained using the Ta₂AlC-deposited SPF and D-shaped fiber were slightly longer, with values of 1.734 ps and 1.817 ps being measured. Overall, it is seen that all the evanescent field-based SA devices could generate mode-locked pulses in the 2 μm region, with the Ta₂AlC-deposited tapered fiber providing the best performance in terms of the low mode-locking threshold, the highest output, and peak power, as well as having the shortest pulse width.

Methods

Preparation of Ta₂AlC MAX Phase.

The Ta₂AlC was purchased from Forsman Scientific (Beijing) Co., Ltd in powder form with a purity of ≥ 98 %. It was prepared directly by dissolving 30 mg of Ta₂AlC powder in 3 ml isopropyl alcohol (IPA) and then sonicated for 1 hour.

Laser Characterization.

The output characteristics of the laser were obtained using a Rohde & Schwarz RTM3002 oscilloscope (OSC), Yokogawa AQ6375 optical spectrum analyzer (OSA), Rohde & Schwarz FPC1000 radio frequency spectrum analyzer (RFSA), and an A.P.E. PulseCheck150 optical autocorrelator.

Conclusion

Mode-locked thulium-doped fiber lasers (TDFLs) operating in the 2 μm were successfully demonstrated using three different evanescent field-based devices: tapered fiber, side-polished fiber (SPF), and also the D-shaped fiber. The tapered fiber was fabricated using the flame-brushing method, while the D-shaped fiber was polished using the wheel-polishing technique. Meanwhile, the SPF was obtained commercially. A new type of MAX phase, the Ta₂AlC, was drop-casted onto the three fibers and then inserted separately into the TDFL cavity. The MAX phase was composed of tantalum (Ta) as the early transition metal instead of the common titanium (Ti). Stable mode-locked pulses were obtained for all cases using the three SA devices. The stability measurements showed little to no fluctuations in the center wavelength

and the peak optical power of the mode-locked TDFLs. The SNR values were also recorded to be more than 47 dB for all three cases, further proving the stability of the generated mode-locked pulses. It was observed from the results that the Ta₂AlC-deposited fiber had the best performance as it can generate the highest output power of 1.91 mW, the highest peak power of 106 W, and had the shortest pulse width of 1.678 ps. These demonstrations show the potential of evanescent field-based devices to be used as SAs for the development of high-power ultrafast fiber lasers sensing.

Declarations

Acknowledgement

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Author contributions statements

Harith Ahmad proposed the study and designed the experiment. Muhammad Fadhli Mohd Azri and Rizal Ramli performed the experiments. Muhamad Zharif Samion contributed to data analysis and writing of the manuscript. Norazriena Yusoff prepared the Nb₂C MXene, performed the material characterization and contributed to writing of the manuscript. Kok Sing Lim contributed to data analysis and writing of the manuscript.

Competing interests

The authors declare no competing financial interests.

Additional information

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References

1. Harun, S. W., Dimyati, K., Jayapalan, K. K. & Ahmad, H. An overview on S-band erbium-doped fiber amplifiers. *Laser Phys. Lett*, **4**, 10–15 (2007).
2. Supe, A. *et al.* Raman Assisted Fiber Optical Parametric Amplifier for S-Band Multichannel Transmission System. *Fibers*, **9**, 9 (2021).
3. Mears, R. J., Reekie, L., Jauncey, I. M. & Payne, D. N. Low-noise erbium-doped fibre amplifier operating at 1.54 μm . *Electron. Lett*, **23**, 1026–1028 (1987).

4. Pedersen, B., Bjarklev, A., Povlsen, J. H., Dybdal, K. & Larsen, C. C. The design of erbium-doped fiber amplifiers. *J. Light. Technol*, **9**, 1105–1112 (1991).
5. Harun, S. W., Tamchek, N., Poopalan, P. & Ahmad, H. Double-pass L-band EDFA with enhanced noise figure characteristics. *IEEE Photonics Technol. Lett*, **15**, 1055–1057 (2003).
6. Singh, S. & Kaler, R. S. Flat-gain L-band raman-EDFA hybrid optical amplifier for dense wavelength division multiplexed system. *IEEE Photonics Technol. Lett*, **25**, 250–252 (2013).
7. Harun, S. W., Low, S. K., Poopalan, P. & Ahmad, H. Gain clamping in L-band erbium-doped fiber amplifier using a fiber Bragg grating. *IEEE Photonics Technol. Lett*, **14**, 293–295 (2002).
8. Yang, J., Meng, X. & Liu, C. Accurately control and flatten gain spectrum of L-band erbium doped fiber amplifier based on suitable gain-clamping. *Opt. Laser Technol*, **78**, 74–78 (2016).
9. Ahmad, H., Shahi, S. & Harun, S. W. Bismuth-based erbium-doped fiber as a gain medium for L-band amplification and Brillouin fiber laser. *Laser Phys*, **20**, 716–719 (2010).
10. Lu, Y., Bin, Chu, P. L., Alphones, A. & Shum, P. A 105-nm ultrawide-band gain-flattened amplifier combining C-and L-band dual-core EDFAs in a parallel configuration. *IEEE Photonics Technol. Lett*, **16**, 1640–1642 (2004).
11. Ismail, E. I., Kadir, N. A., Latiff, A. A., Ahmad, H. & Harun, S. W. Black phosphorus crystal as a saturable absorber for both a Q-switched and mode-locked erbium-doped fiber laser. *RSC Adv*, **6**, 72692–72697 (2016).
12. Tuo, M. *et al.* Ultrathin 2D transition metal carbides for ultrafast pulsed fiber lasers. *ACS Photonics*, **5**, 1808–1816 (2018).
13. Soltanian, M. R. K. *et al.* A stable dual-wavelength Thulium-doped fiber laser at 1.9 μm using photonic crystal fiber. *Sci. Rep*, **5**, 1–8 (2015).
14. Zhou, Y., Lou, S., Tang, Z., Zhao, T. & Zhang, W. Tunable and switchable C-band and L-band multi-wavelength erbium-doped fiber laser employing a large-core fiber filter. *Opt. Laser Technol*, **111**, 262–270 (2019).
15. Ahmad, H., Zulkifli, M. Z., Thambiratnam, K., Latif, S. F. & Harun, S. W. High power and compact switchable bismuth based multiwavelength fiber laser. *Laser Phys. Lett*, **6**, 380 (2009).
16. Husshini, N. F. H. *et al.* Multiwavelength fiber laser employing semiconductor optical amplifier in nonlinear optical loop mirror with polarization controller and polarization maintaining fiber. in *AIP Conference Proceedings* vol. 2203 20030 (AIP Publishing LLC, 2020).
17. Harun, S. W., Shirazi, M. R. & Ahmad, H. Multiple wavelength Brillouin fiber laser from injection of intense signal light. *Laser Phys. Lett*, **4**, 678 (2007).
18. Qureshi, K. K. Multiwavelength Fiber Laser Covering far L and U Bands in a Dual Cavity Configuration. *IEEE Photonics Technol. Lett*, **33**, 321–324 (2021).
19. Parvizi, R., Arof, H., Ali, N. M., Ahmad, H. & Harun, S. W. 0.16 nm spaced multi-wavelength Brillouin fiber laser in a figure-of-eight configuration. *Opt. Laser Technol*, **43**, 866–869 (2011).

20. Oehler, A. E. H., Zeller, S. C., Weingarten, K. J. & Keller, U. Broad multiwavelength source with 50 GHz channel spacing for wavelength division multiplexing applications in the telecom C band. *Opt. Lett*, **33**, 2158–2160 (2008).
21. Ahmad, H., Zulkifli, M. Z., Latif, A. A. & Harun, S. W. Tunable dual wavelength fiber laser incorporating AWG and optical channel selector by controlling the cavity loss. *Opt. Commun*, **282**, 4771–4775 (2009).
22. He, X., Fang, X., Liao, C., Wang, D. N. & Sun, J. A tunable and switchable single-longitudinal-mode dual-wavelength fiber laser with a simple linear cavity. *Opt. Express*, **17**, 21773–21781 (2009).
23. Shang, Y., Guo, R., Liu, Y. & Yi, X. Managing Brillouin frequency spacing for temperature measurement with Brillouin fiber laser sensor. *Opt. Quantum Electron*, **52**, 1–8 (2020).
24. Yasin, M., Harun, S. W., Abdul-Rashid, H. A. & Ahmad, H. The performance of a fiber optic displacement sensor for different types of probes and targets. *Laser Phys. Lett*, **5**, 55 (2007).
25. Guo, K., He, J., Xu, G. & Wang, Y. Dual-polarization distributed feedback fiber laser sensor based on femtosecond laser-inscribed in-fiber stressors for simultaneous strain and temperature measurements. *IEEE Access*, **8**, 97823–97829 (2020).
26. Whitenett, G., Stewart, G., Yu, H. & Culshaw, B. Investigation of a tuneable mode-locked fiber laser for application to multipoint gas spectroscopy. *J. Light. Technol*, **22**, 813 (2004).
27. McAleavey, F. J. *et al.* Narrow linewidth, tunable Tm³⁺-doped fluoride fiber laser for optical-based hydrocarbon gas sensing. *IEEE J. Sel. Top. quantum Electron*, **3**, 1103–1111 (1997).
28. Fried, N. M. & Murray, K. E. High-power thulium fiber laser ablation of urinary tissues at 1.94 μm . *J. Endourol*, **19**, 25–31 (2005).
29. Spiers, G. D. *et al.* Atmospheric CO₂ measurements with a 2 μm airborne laser absorption spectrometer employing coherent detection. *Appl. Opt*, **50**, 2098–2111 (2011).
30. Scholle, K., Lamrini, S., Koopmann, P. & Fuhrberg, P. 2 μm Laser Sources and Their Possible Applications(IntechOpen, 2010).
31. Shoji, T. *et al.* Glaucomatous changes in lamina pores shape within the lamina cribrosa using wide bandwidth, femtosecond mode-locked laser OCT. *PLoS One*, **12**, e0181675 (2017).
32. Li, Z. *et al.* Thulium-doped fiber amplifier for optical communications at 2 μm . *Opt. Express*, **21**, 9289–9297 (2013).
33. Sotor, J. *et al.* All-fiber Ho-doped mode-locked oscillator based on a graphene saturable absorber. *Opt. Lett*, **41**, 2592–2595 (2016).
34. Villanueva, G. E., Ferri, M. & Perez-Millan, P. Active and passive mode-locked fiber lasers for high-speed high-resolution photonic analog-to-digital conversion. *IEEE J. Quantum Electron*, **48**, 1443–1452 (2012).
35. Kim, J., Koo, J. & Lee, J. H. All-fiber acousto-optic modulator based on a cladding-etched optical fiber for active mode-locking. *Photonics Res*, **5**, 391–395 (2017).

36. Chong, A., Wright, L. G. & Wise, F. W. Ultrafast fiber lasers based on self-similar pulse evolution: a review of current progress. *Reports Prog. Phys*, **78**, 113901 (2015).
37. Oktem, B., Ülgüdür, C. & Ilday, F. Ö. Soliton–similariton fibre laser. *Nat. Photonics* **4**, 307–311(2010).
38. Tang, Y., Chong, A. & Wise, F. W. Generation of 8 nJ pulses from a normal-dispersion thulium fiber laser. *Opt. Lett*, **40**, 2361–2364 (2015).
39. Chong, A., Buckley, J., Renninger, W. & Wise, F. All-normal-dispersion femtosecond fiber laser. *Opt. Express*, **14**, 10095–10100 (2006).
40. Gomes, L. A., Orsila, L., Jouhti, T. & Okhotnikov, O. G. Picosecond SESAM-based ytterbium mode-locked fiber lasers. *IEEE J. Sel. Top. Quantum Electron*, **10**, 129–136 (2004).
41. Jung, I. D. *et al.* Semiconductor saturable absorber mirrors supporting sub-10-fs pulses. *Appl. Phys. B Lasers Opt.* **65**, (1997).
42. Ma, C. *et al.* Recent progress in ultrafast lasers based on 2D materials as a saturable absorber. *Appl. Phys. Rev*, **6**, 41304 (2019).
43. Sun, Z. *et al.* Graphene mode-locked ultrafast laser. *ACS Nano*, **4**, 803–810 (2010).
44. Solodyankin, M. A. *et al.* Mode-locked 1.93 μ m thulium fiber laser with a carbon nanotube absorber. *Opt. Lett*, **33**, 1336–1338 (2008).
45. Zhao, R. *et al.* Triwavelength synchronously mode-locked fiber laser based on few-layered black phosphorus. *Appl. Phys. Express*, **9**, 092701 (2016).
46. Chen, H. *et al.* Transition-metal dichalcogenides heterostructure saturable absorbers for ultrafast photonics. *Opt. Lett*, **42**, 4279–4282 (2017).
47. Yan, P. *et al.* A practical topological insulator saturable absorber for mode-locked fiber laser. *Sci. Rep*, **5**, 1–5 (2015).
48. Gogotsi, Y. & Anasori, B. The Rise of MXenes. *ACS Nano* **13**, 8491 – 8494(2019).
49. Li, J. *et al.* Highly stable femtosecond pulse generation from a MXene Ti₃C₂T_x (T = F, O, or OH) mode-locked fiber laser. *Photonics Res*, **7**, 260–264 (2019).
50. Barsoum, M. W. & El-Raghy, T. The MAX phases: unique new carbide and nitride materials: ternary ceramics turn out to be surprisingly soft and machinable, yet also heat-tolerant, strong and lightweight. *Am. Sci*, **89**, 334–343 (2001).
51. Ahmad, H., Kamely, A. A., Yusoff, N., Bayang, L. & Samion, M. Z. Generation of Q-switched Pulses in Thulium-doped and Thulium/Holmium-co-doped Fiber Lasers using MAX phase (Ti₃AlC₂). *Sci. Rep*, **10**, 1–10 (2020).
52. Mo, Y., Rulis, P. & Ching, W. Y. Electronic structure and optical conductivities of 20 MAX-phase compounds. *Phys. Rev. B*, **86**, 165122 (2012).
53. Lee, J., Kwon, S. & Lee, J. H. Ti₂AlC-based saturable absorber for passive Q-switching of a fiber laser. *Opt. Mater. Express*, **9**, 2057–2066 (2019).
54. Jafry, A. A. A. *et al.* MAX phase based saturable absorber for mode-locked erbium-doped fiber laser. *Opt. Laser Technol*, **127**, 106186 (2020).

55. Martinez, A., Fuse, K., Xu, B. & Yamashita, S. Optical deposition of graphene and carbon nanotubes in a fiber ferrule for passive mode-locked lasing. *Opt. Express*, **18**, 23054–23061 (2010).
56. Zhang, M. *et al.* Tm-doped fiber laser mode-locked by graphene-polymer composite. *Opt. Express*, **20**, 25077–25084 (2012).
57. Cheng, P. K. *et al.* Passively Q-switched Ytterbium-doped fiber laser based on broadband multilayer Platinum Ditelluride (PtTe₂) saturable absorber. *Sci. Rep.*, **9**, 1–7 (2019).
58. Mouchel, P. *et al.* High power passively mode-locked fiber laser based on graphene nanocoated optical taper. *Appl. Phys. Lett.*, **111**, 31106 (2017).
59. Zhou, Y. *et al.* A passively mode-locked thulium-doped fiber laser based on a D-shaped fiber deposited with PbS nanoparticles. *J. Mater. Chem. C*, **7**, 11215–11219 (2019).
60. Li, X. *et al.* High-power graphene mode-locked Tm/Ho co-doped fiber laser with evanescent field interaction. *Sci. Rep.*, **5**, 1–8 (2015).
61. Khazaeinezhad, R. *et al.* Mode-Locked All-Fiber Lasers at Both Anomalous and Normal Dispersion Regimes Based on Spin-Coated MoS₂ Sheets on a Side-Polished Fiber. *IEEE Photonics J*, **7**, 1–9 (2014).
62. Hu, C. *et al.* In-situ reaction synthesis and decomposition of Ta₂AlC. *Int. J. Mater. Res.*, **99**, 8–13 (2008).
63. Hu, C. *et al.* Microstructure and properties of bulk Ta₂AlC ceramic synthesized by an in situ reaction/hot pressing method. *J. Eur. Ceram. Soc.*, **28**, 1679–1685 (2008).
64. Yeh, C. L. & Shen, Y. G. Effects of Al content on formation of Ta₂AlC by self-propagating high-temperature synthesis. *J. Alloys Compd.*, **482**, 219–223 (2009).
65. Du, J. *et al.* Phosphorene quantum dot saturable absorbers for ultrafast fiber lasers. *Sci. Rep.*, **7**, 1–10 (2017).
66. Woodward, R. I., Hasan, T. & Kelleher, E. J. R. Optical nonlinearity of few-layer MoS₂ devices and applications in short-pulse laser technology. *Proc. Opto-Electron. Commun. Conf. (OECC)* 1–3(2015).
67. Ahmad, H., Hassan, H., Zulkifli, A. Z., Thambiratnam, K. & Amiri, I. S. Characterization of arc-shaped side-polished fiber. *Opt. Quantum Electron.*, **49**, 207 (2017).
68. Feng, J. *et al.* SnS nanosheets for 105th harmonic soliton molecule generation. *Ann. Phys.*, **531**, 1900273 (2019).
69. Liu, W. *et al.* CVD-grown MoSe₂ with high modulation depth for ultrafast mode-locked erbium-doped fiber laser. *Sci. Rep.*, **29**, 394002 (2018).
70. Cabasse, A., Martel, G. & Oudar, J. L. High power dissipative soliton in an Erbium-doped fiber laser mode-locked with a high modulation depth saturable absorber mirror. *Opt. Express*, **17**, 9537–9542 (2009).
71. Wang, Y. R. *et al.* Ultralow saturation intensity topological insulator saturable absorber for gigahertz mode-locked solid-state lasers. *IEEE Photonics J*, **10**, 1–10 (2018).

72. Man, W. S., Tam, H. Y., Demokan, M. S., Wai, P. K. A. & Tang, D. Y. Mechanism of intrinsic wavelength tuning and sideband asymmetry in a passively mode-locked soliton fiber ring laser. *JOSA B*, **17**, 28–33 (2000).
73. Sotor, J. *et al.* Ultrafast thulium-doped fiber laser mode locked with black phosphorus. *Opt. Lett*, **40**, 3885–3888 (2015).
74. Birks, T. A. & Li, Y. W. The shape of fiber tapers. *J. Light. Technol*, **10**, 432–438 (1992).
75. Ryu, S. Y., Kim, K. S., Kim, J. & Kim, S. Degradation of optical properties of a film-type single-wall carbon nanotubes saturable absorber (SWNT-SA) with an Er-doped all-fiber laser. *Opt. Express*, **20**, 12966–12974 (2012).

Figures

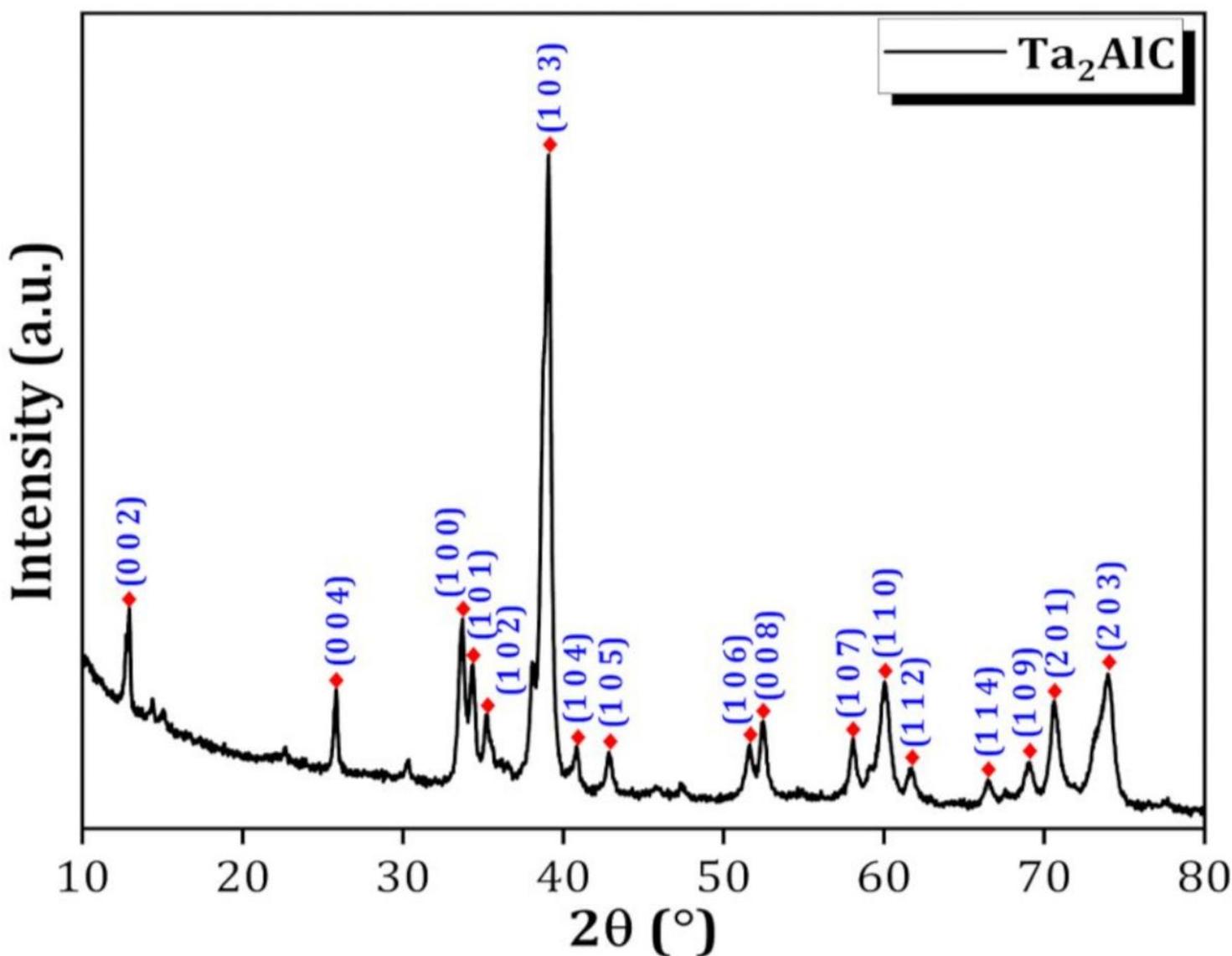


Figure 1

XRD pattern of the Ta₂AlC MAX Phase.

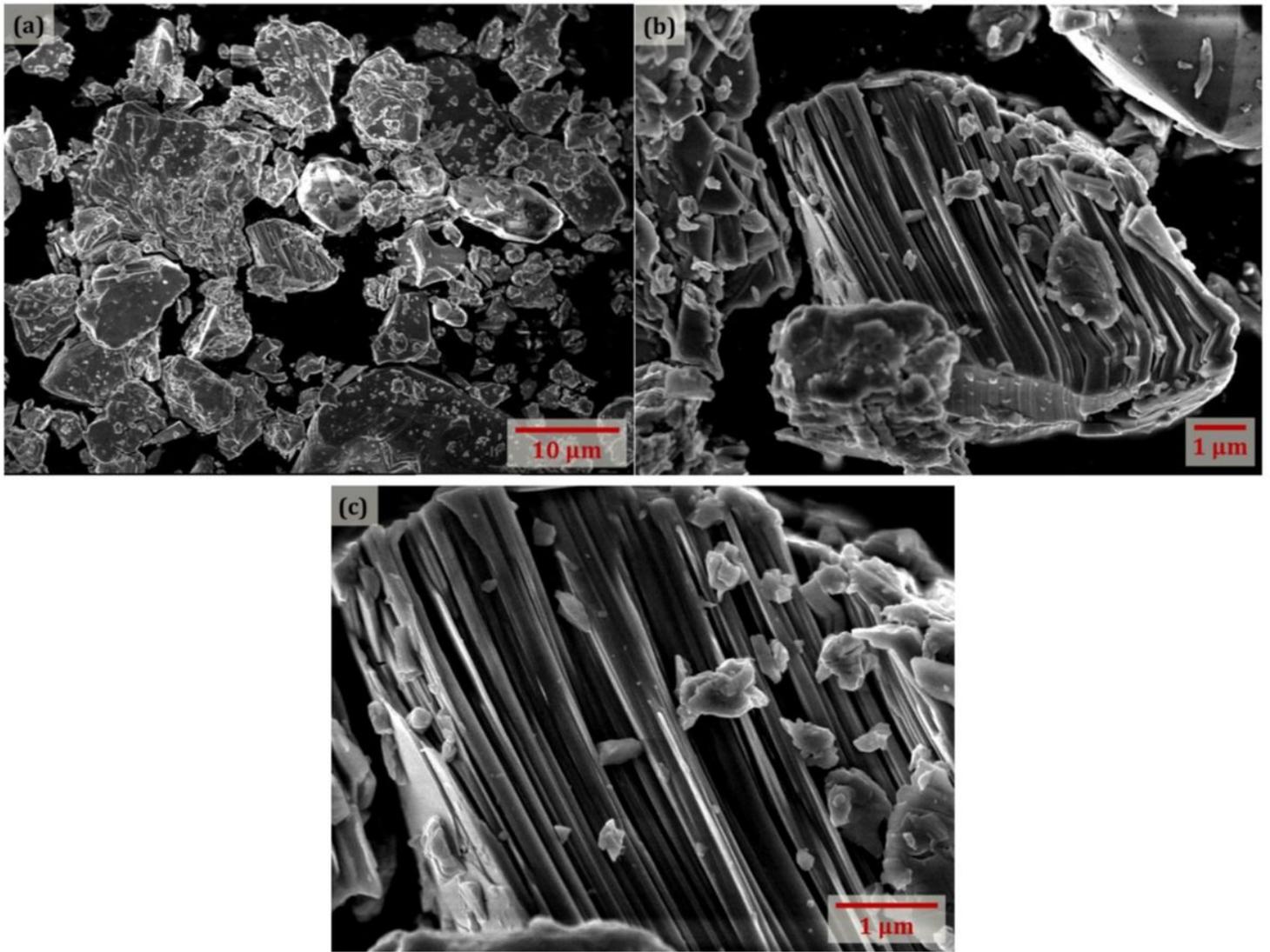


Figure 2

FESEM images of Ta₂AlC MAX Phase captured at (a) 5k, (b) 10k, and (c) 20k x magnification with 5.0 kV of operating voltage.

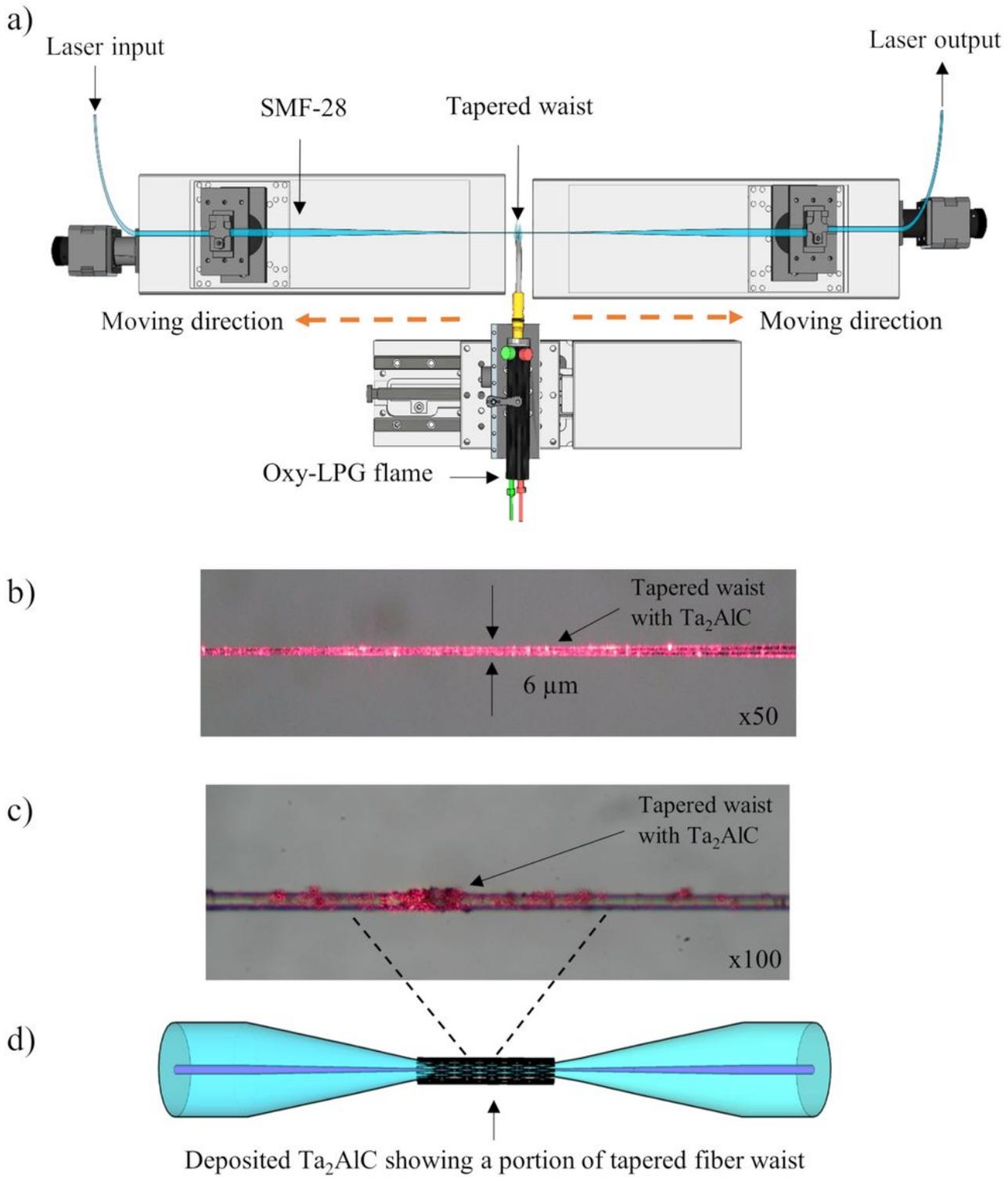


Figure 3

(a) The setup of flame brushing method to fabricate tapered fiber. (b) The microscopic image of the tapered fiber at x50 and (c) x100 magnification with the material. (d) The schematic drawing of the coated Ta_2AlC tapered fiber.

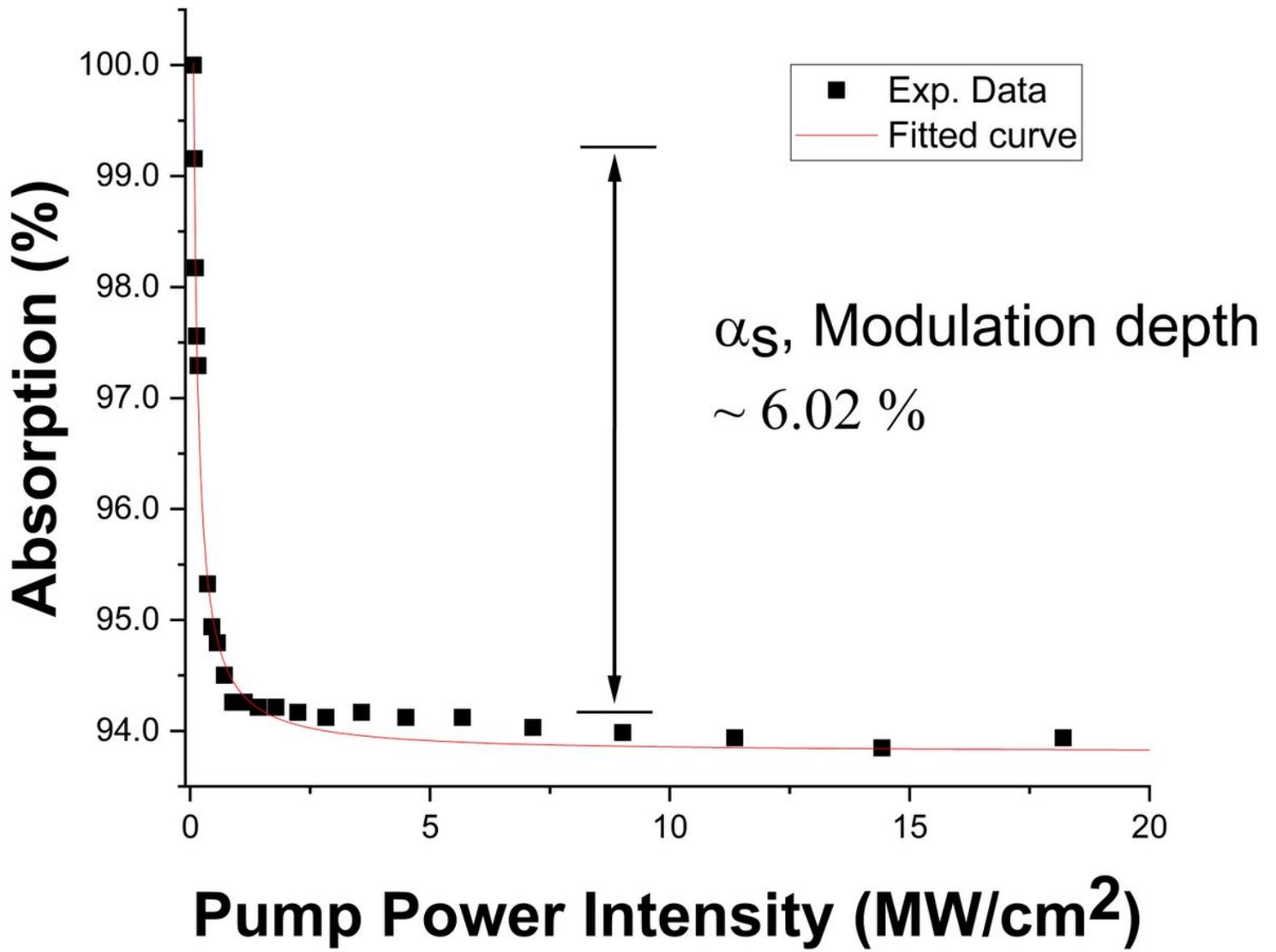


Figure 4

The nonlinear absorption measurement of the Ta₂AlC-based tapered fiber.

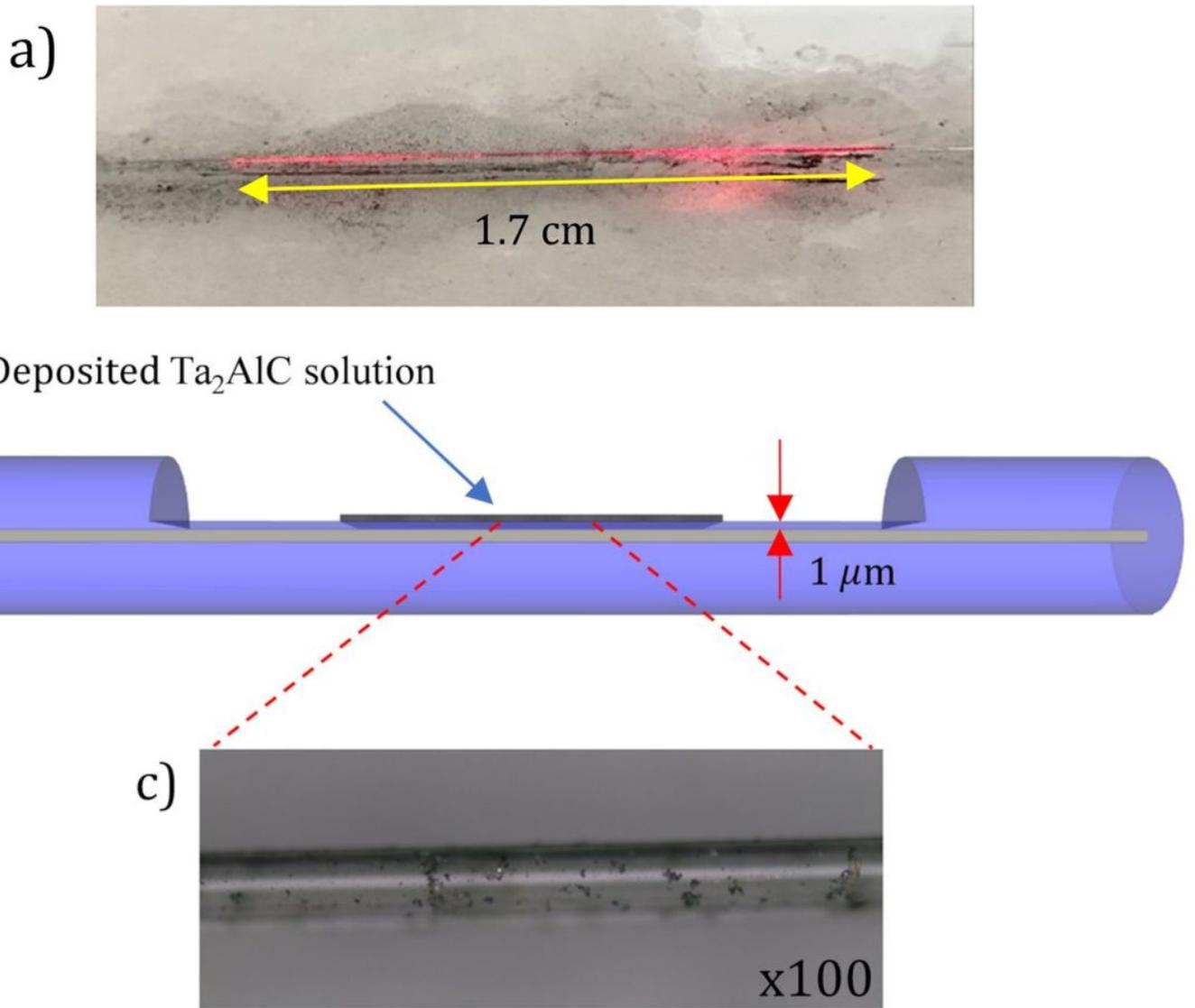


Figure 5

(a) The evanescent field at the polished region using a red-light source. (b) The deposited Ta_2AlC on the SPF (c) Microscopic image of the polished part of the side polished fiber at 100 times magnification.

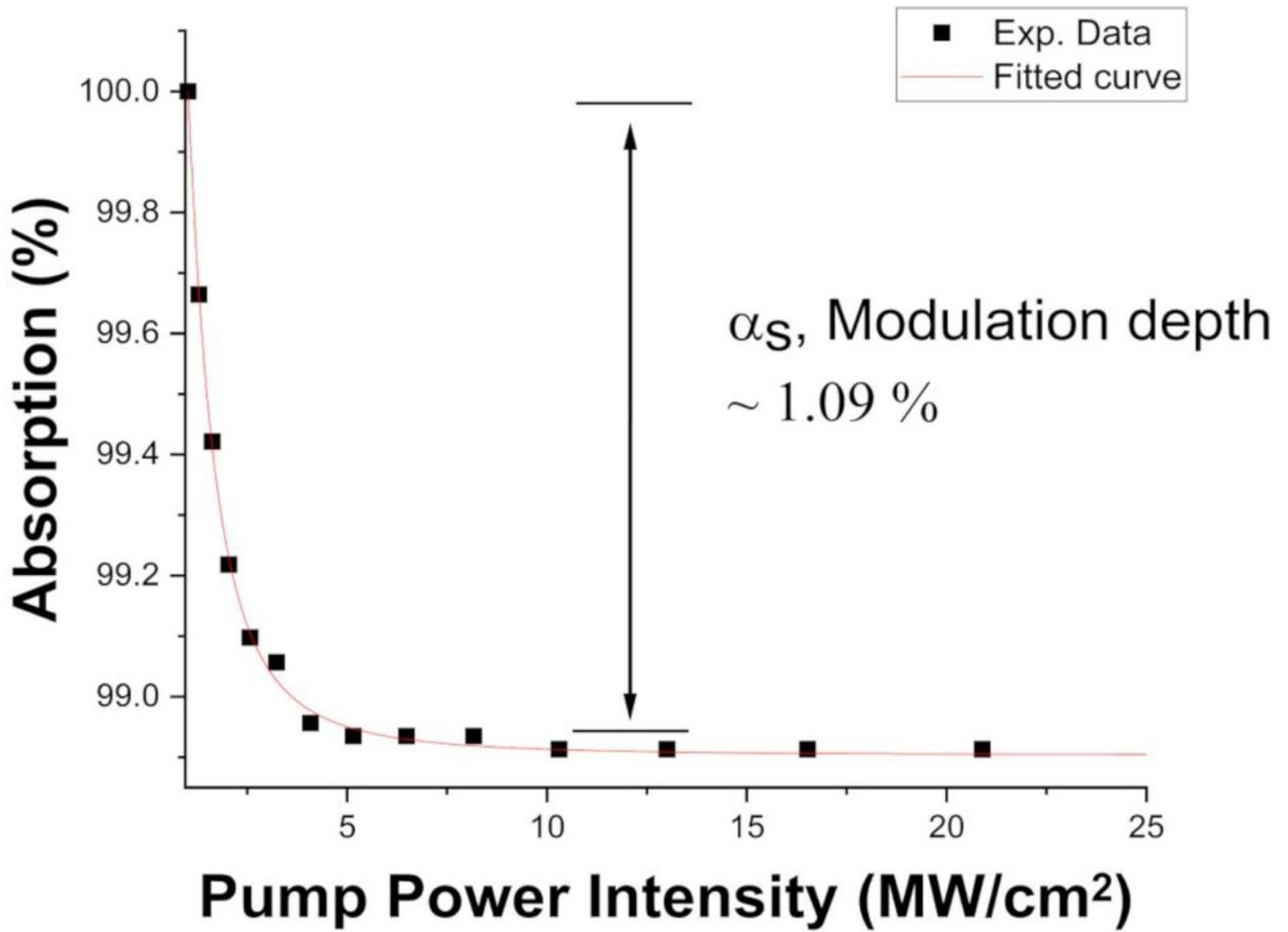


Figure 6

The nonlinear absorption test of the side polished fiber Ta2AlC.

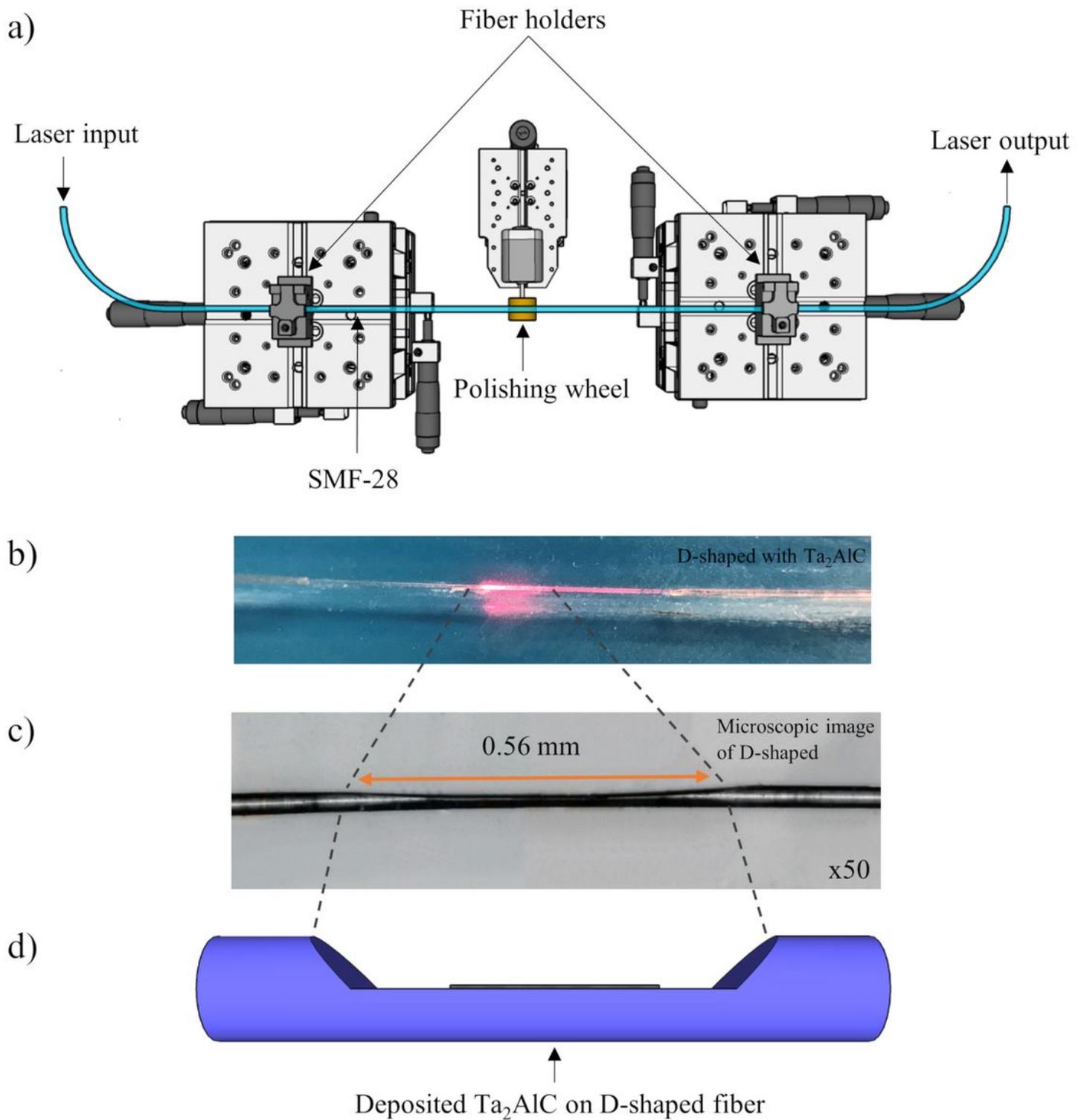


Figure 7

(a) Schematic illustration of wheel polishing method to fabricate the D-shaped fiber (b) The interaction of the D-shaped fiber embedded with the Ta_2AlC nanoparticles when a red-light source was injected (c) The microscopic image of the side profile of the D-shaped fiber at x50 magnification (d) The illustration of the deposited Ta_2AlC onto the D-shaped fiber.

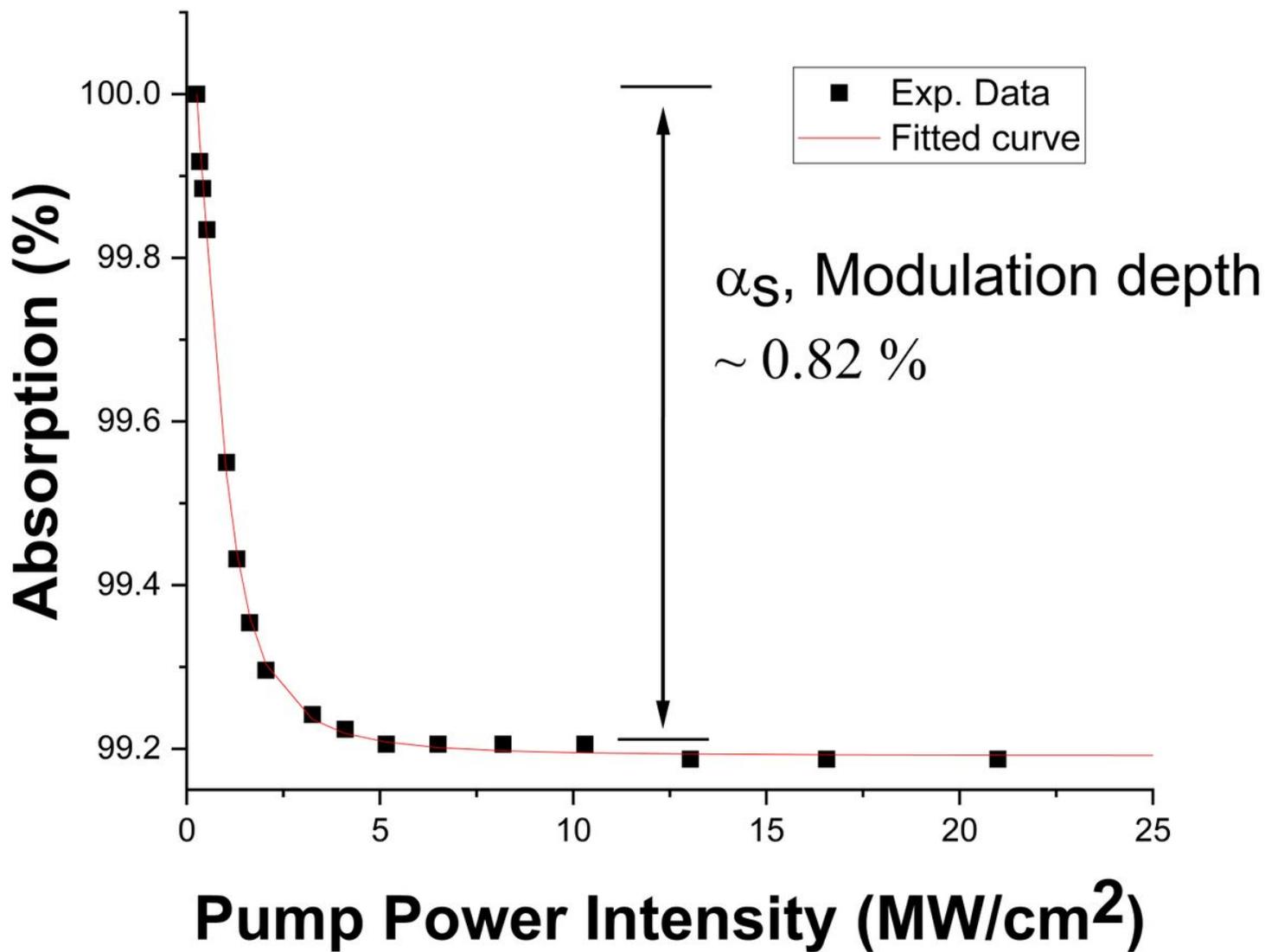


Figure 8

The nonlinear absorption test of the Ta₂AlC-deposited D-shaped fiber.

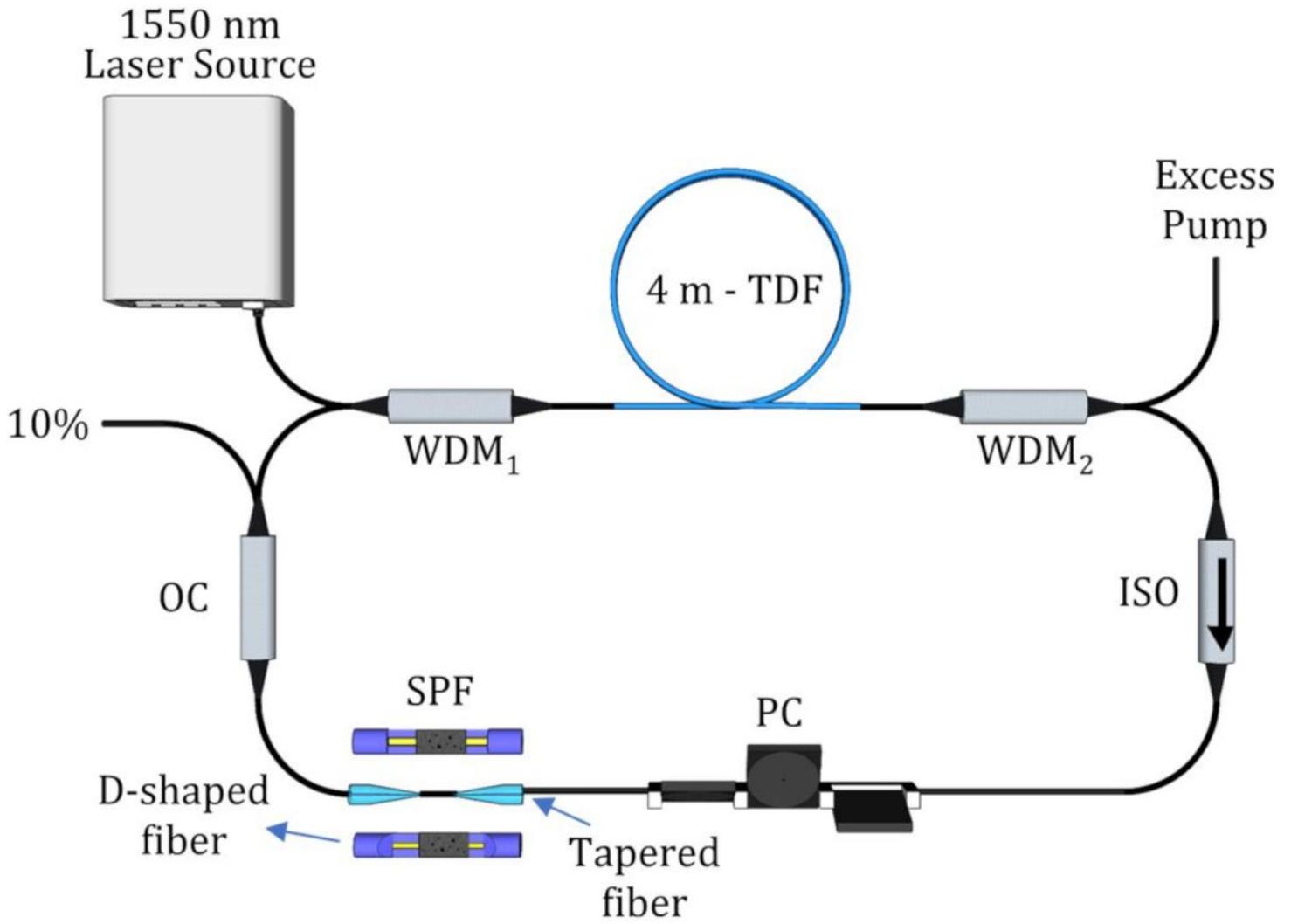


Figure 9

Illustration of the experimental setup for the 2 μm mode-locked Thulium-doped fiber laser (TDFL).

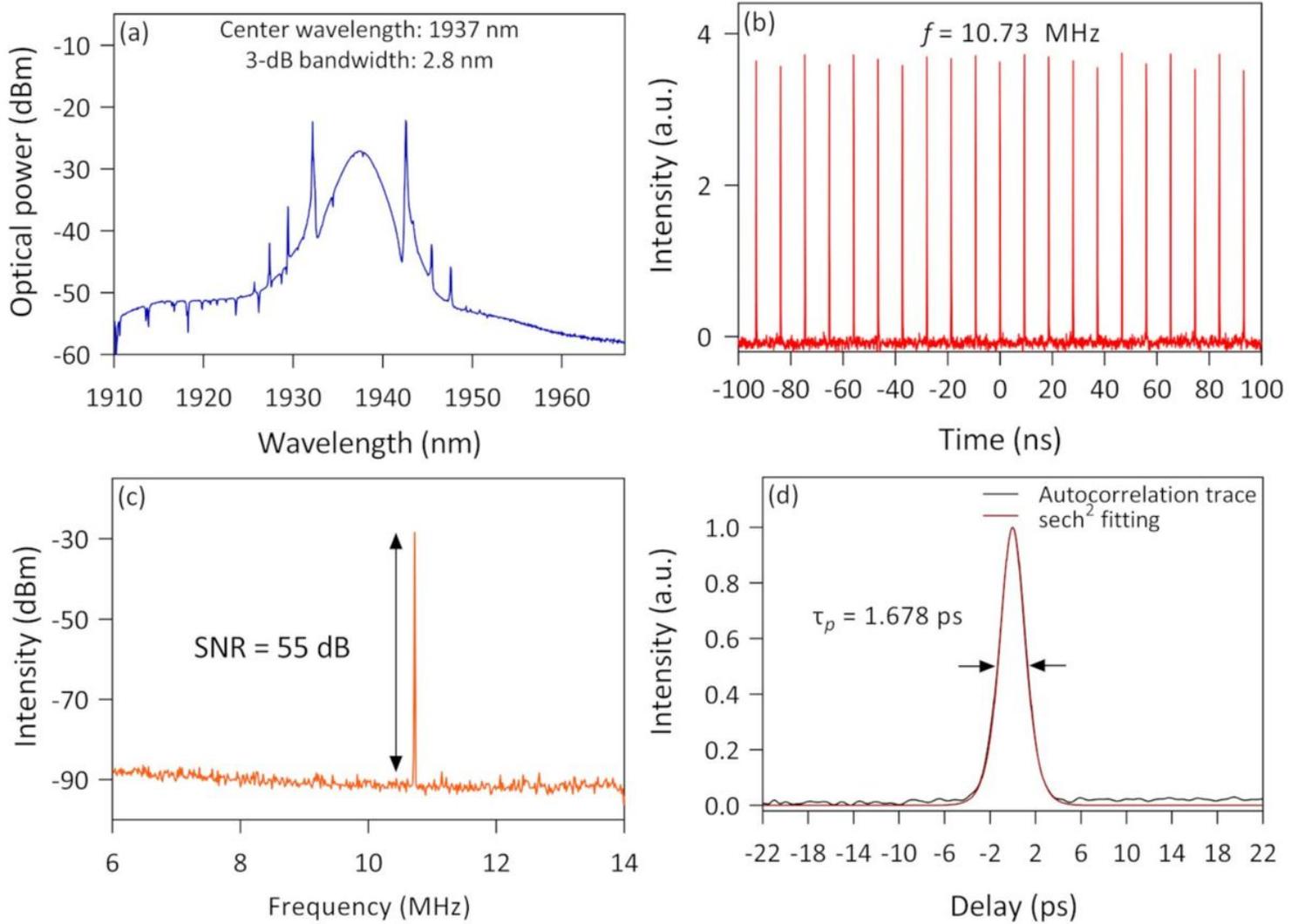


Figure 10

The output characteristic for mode-locked thulium-doped fiber laser using the Ta₂AlC tapered fiber saturable absorber (a) optical spectrum, (b) pulse train, (c) radio frequency spectrum, and (d) autocorrelation trace.

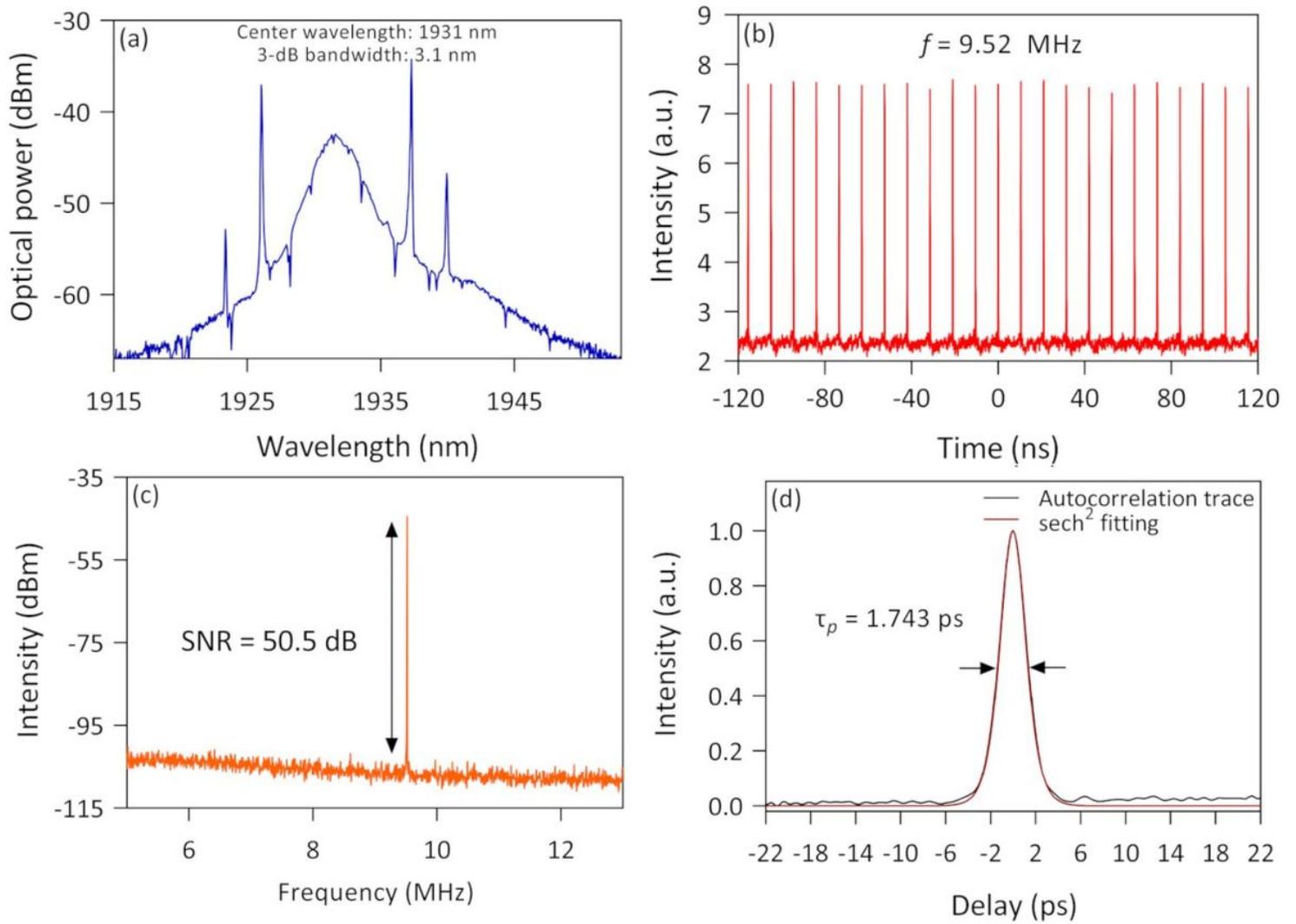


Figure 11

The output characteristic for mode-locked thulium-doped fiber laser using Ta2AlC side polished fiber saturable absorber (a) optical spectrum, (b) pulse train, (c) radio frequency spectrum, and (d) autocorrelation trace.

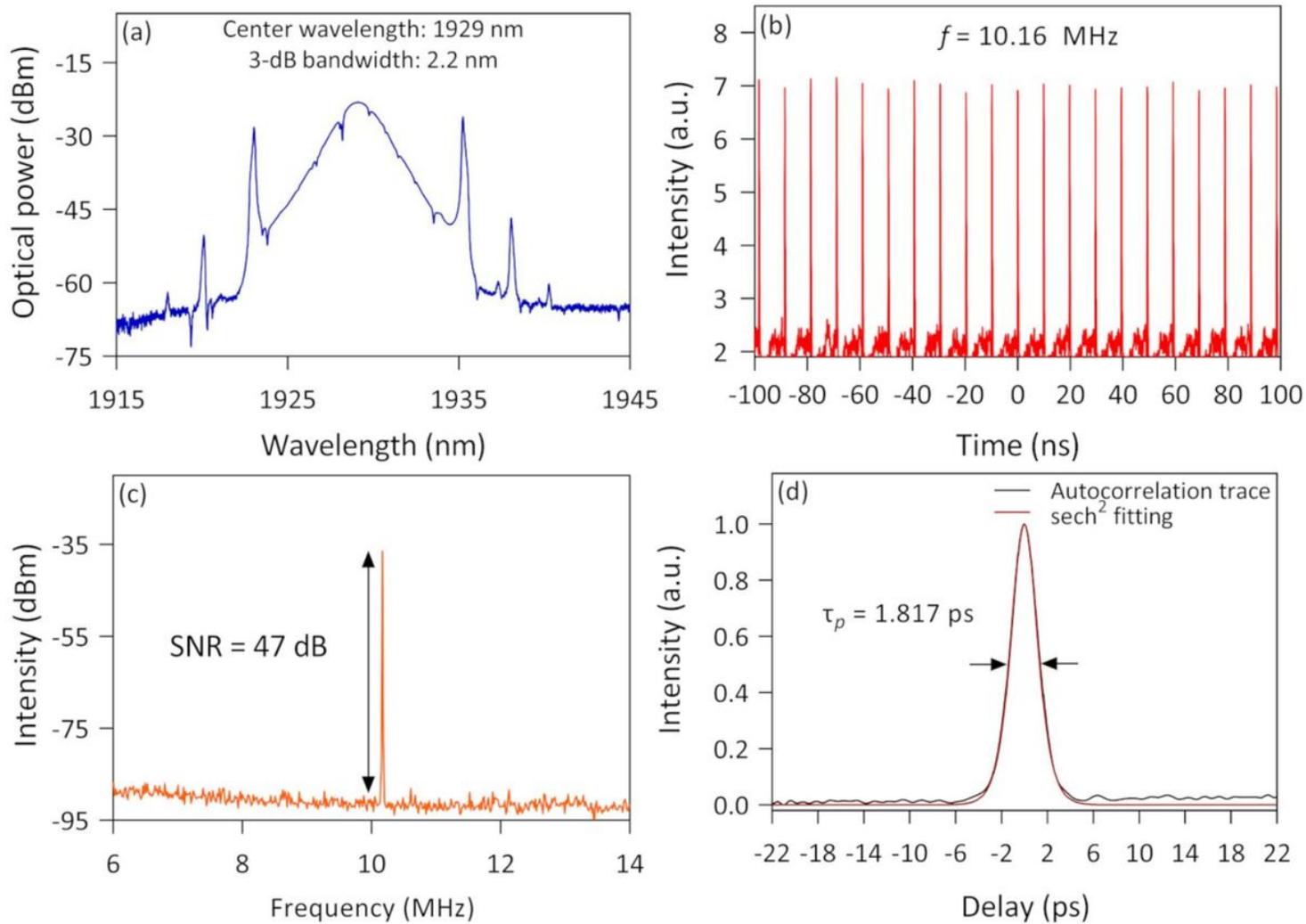


Figure 12

The output characteristic for mode-locked thulium-doped fiber laser using Ta₂AlC D-shaped fiber saturable absorber: (a) optical spectrum, (b) soliton train, (c) radio frequency spectrum, and (d) autocorrelation trace.

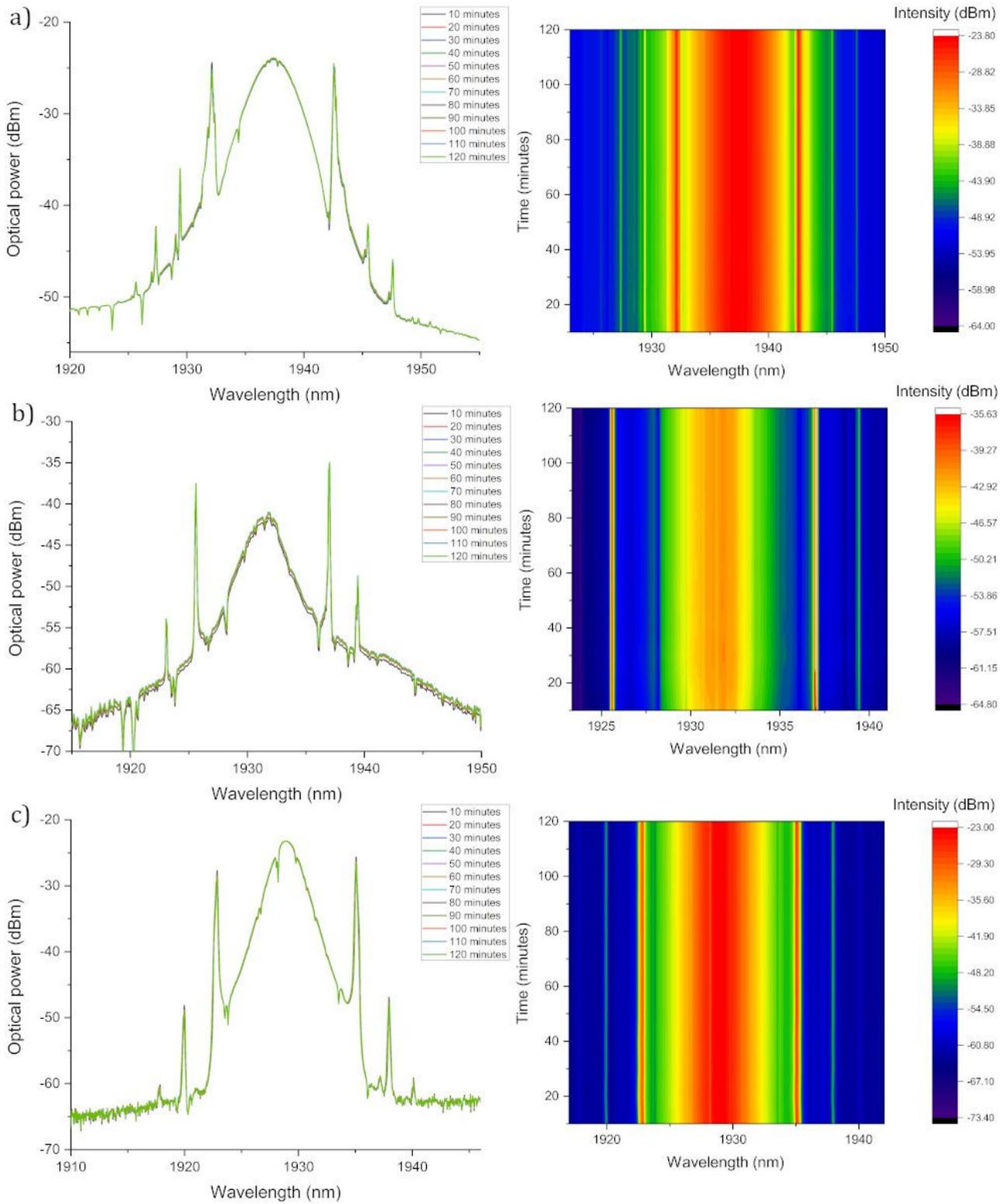


Figure 13

The stability of the output pulse spectrum for Ta₂AlC and the corresponding contour plot for (a) tapered fiber, (b) SPF, and (c) D-shaped fiber saturable absorber. The stability test was conducted for two hours at a 10-minute interval.

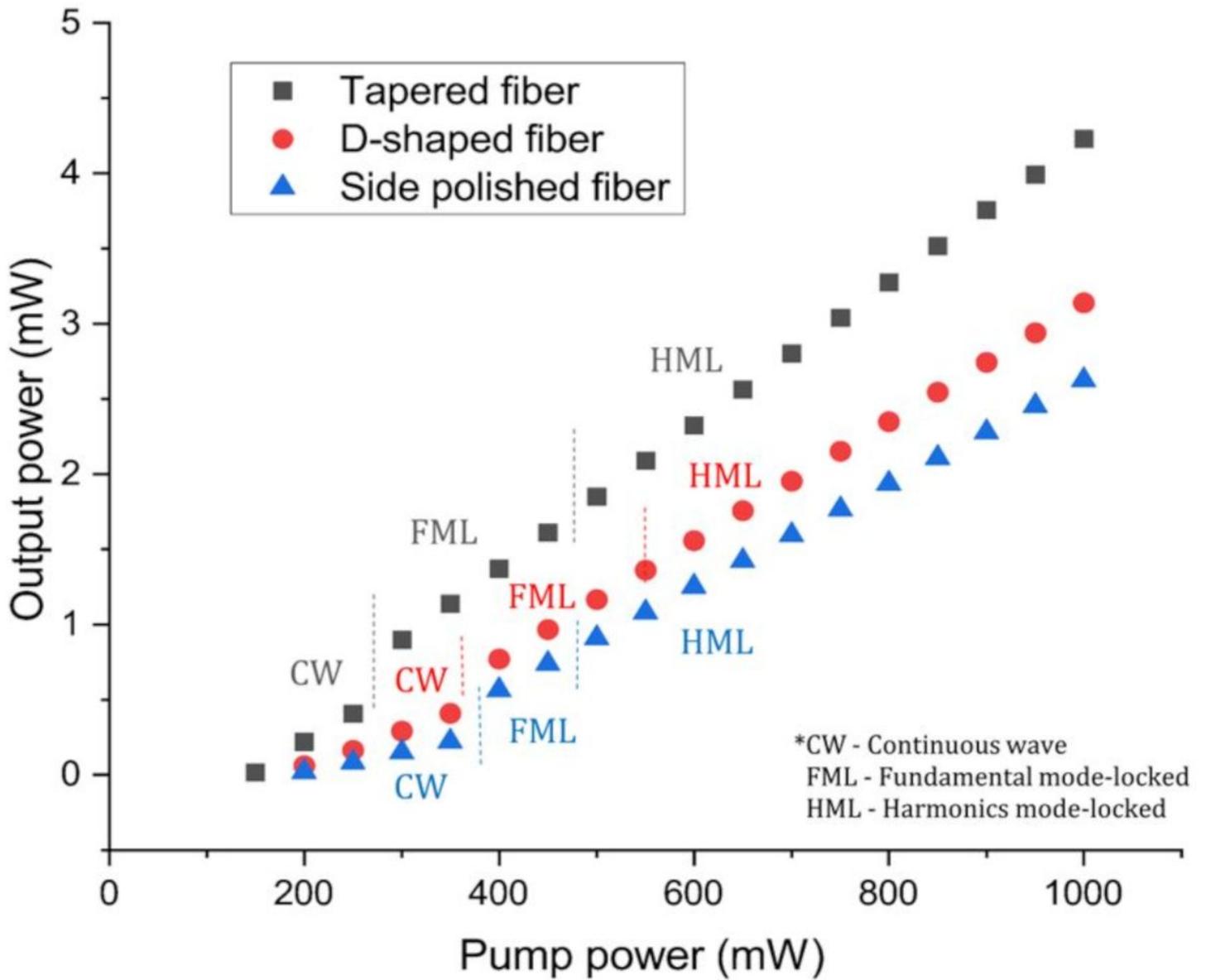


Figure 14

The pump power versus the output power of the mode-locked TDFLs.